

International Institute of Seismology  
and Earthquake Engineering (IISEE)  
Seismology Course Lecture Notes

TRAINING COURSE IN  
SEISMOLOGY AND EARTHQUAKE ENGINEERING

# **Introduction to Digital Data Processing**

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# 1. Introduction

In this lecture note, several basic and essential topics that are useful for understanding digital data processing techniques are explained. These topics are directly related to measurements using digital equipment.

## 1.1. Linear System

Almost all measuring equipment is “linear systems.” This is because “linear systems” make it easier to reproduce the measured value of the physical parameters from the results of the measurement.

### 1.1.1. Signal, Input, and Output

Since it could be difficult to define the terms “signal”, “input” and “output” exactly, we employ their practical definitions.

A signal is any quantity that we measure or input. For example, mechanical signals may be displacement, velocity, acceleration, or force, whereas electric signals may be charge, voltage, or current. Signals often depend on time. In seismology, we typically treat time-dependent signals. Such signals can have many relations. A system can be defined either as a relation between more than two signals itself, or an equipment (or algorithm) that yields such relations. The input can be defined as the conditions given to a system. The signal generated in the system that corresponds to the input is called the output (Fig. 1).

For example, consider ground motion. Ground motion will be the input to a seismograph (which is the “system”). The output is the recorded seismogram. If we consider a seismometer to be the system, the input will be the ground motion again while the voltage imbalance between the two terminals of the seismometer will be the output.



Fig. 1 Block diagram of a system.

### 1.1.2. Linear System

Almost all the systems used in the measurement of physical quantities have “linear” characteristics. This can imply the following. Suppose that the output of a system that corresponds to the input  $x_1(t)$  is  $y_1(t)$  while the output that corresponds to the input  $x_2(t)$  is  $y_2(t)$ . The first requirement of a linear system is that the output corresponding to the input  $x_1(t) + x_2(t)$  is  $y_1(t) + y_2(t)$ . The second is that the output corresponding to  $\alpha x_1(t)$  is  $\alpha y_1(t)$ , where  $\alpha$  is a constant. These requirements lead to a proportional relation between the input and the output. Such a relation is called “linearity.” Namely,

$$\text{If } \begin{cases} x_1(t) \rightarrow y_1(t) \\ x_2(t) \rightarrow y_2(t) \end{cases}, \text{ then } \begin{cases} x_1(t) + x_2(t) \rightarrow y_1(t) + y_2(t) \\ \alpha x_1(t) \rightarrow \alpha y_1(t) \end{cases}$$

Such “linear” characteristics of the relation between an input and an output ensure that the input signal can be reconstructed by using the output signal. This is why almost all measuring equipment is linear systems. However, in reality, linear characteristics can be usually obtained only for a limited range of the input signal. Clipping of a seismometer and saturation of an amplifier are simple examples. When “linearity” is lost, the system becomes “non-linear” (Fig. 2).

There are several important “linear” transformations in mathematics that can be applied to physics, for example, the Fourier transform.

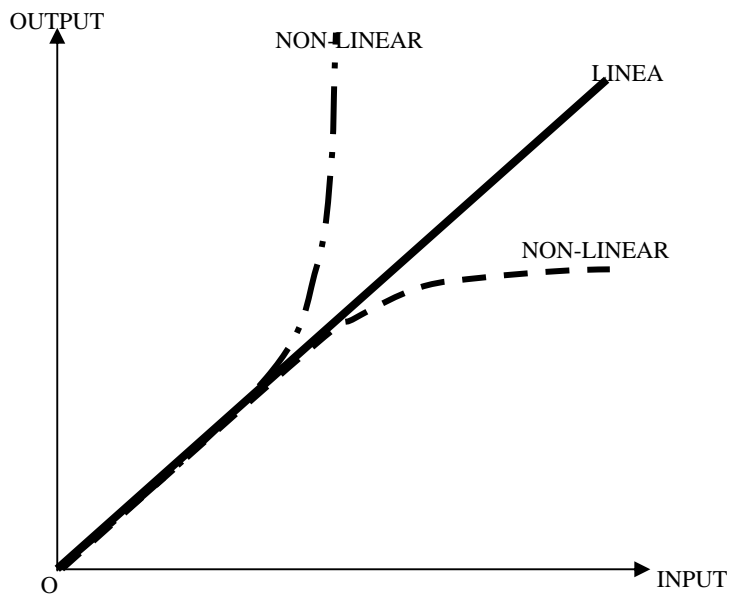


Fig. 2 Schematic figure showing both linearity and non-linearity.

### 1.1.3. Response Characteristics of Linear Systems

Suppose that  $g(t)$  is the output of a system when the input signal is a unit impulse  $\delta(t)$ . An arbitrary function  $x(t)$  can be expressed by using the technique of “convolution” as follows.

$$x(t) = \int_{-\infty}^{\infty} x(u)\delta(t-u)du = \int_{-\infty}^{\infty} x(t-u)\delta(u)du = \lim_{\Delta u \rightarrow 0} \sum_{m=-\infty}^{\infty} x(t-m\Delta u)\delta(m\Delta u).$$

The last term shows that  $x(t)$  is the weighted sum of the delta functions at an infinitively small  $\Delta u$ . Each delta function  $\delta(m\Delta u)$  gives the output  $g(m\Delta u)$ . The second requirement (mentioned previously in 1.1.2) suggests that the input  $x(t-m\Delta u)\delta(m\Delta u)$  gives  $x(t-m\Delta u)g(m\Delta u)$ , since  $x(t-m\Delta u)$  is a constant. The first requirement suggests that the input, that is, the sum of  $x(t-m\Delta u)\delta(m\Delta u)$ , gives the sum of the output for each weighted delta function  $x(t-m\Delta u)g(m\Delta u)$ .

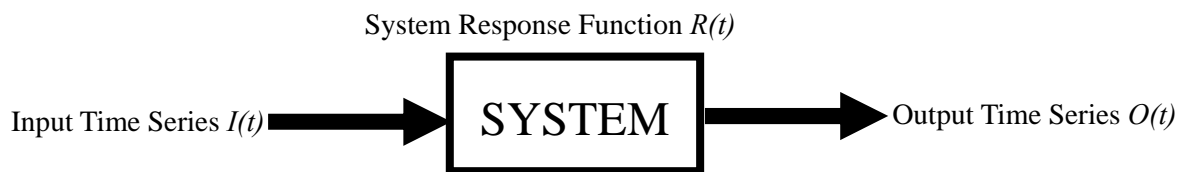
Thus, the output signal that corresponds to the input signal  $x(t)$  is given by the following.

$$\lim_{\Delta u \rightarrow 0} \sum_{m=-\infty}^{\infty} x(t-m\Delta u)g(m\Delta u) = \int_{-\infty}^{\infty} x(t-u)g(u)du = \int_{-\infty}^{\infty} x(u)g(t-u)du = y(t).$$

For a seismic signal, it is reasonable to suppose that  $g(t) = 0$  at  $t < 0$ ; then

$$y(t) = \int_0^{\infty} x(u)g(t-u)du .$$

These formulas show that the output that corresponds to an arbitrary input signal can be defined by  $g(t)$ . In other words, the response to the unit impulse  $g(t)$  contains all information on the characteristics of a linear system. We call such a response “impulse response” or “system characteristics.”



$$O(t) = I(t)*R(t)$$

Fig. 3 System response, input, and output signals.

## 1.2. Digitization of Time-Dependent Functions

Today, digital data acquisition and processing systems are so dominant in physical measurements that we can expect analog or continuous systems to become obsolete by the first half of the twenty-first century. We can convert every analog signal to a digital once they are converted to the form of a voltage imbalance. Some basic knowledge is required to prevent difficulties that might occur during such analog-to-digital conversions, which are referred to as “discretization.”

### 1.2.1. Dirac Comb Function

First, let us understand the discretization process. Consider a box car function  $b(t)$ :

$$b(t) = \begin{cases} B, & |t| \leq t_0 / 2, \\ 0, & |t| > t_0 / 2. \end{cases}$$

The Fourier expansion in  $[-T/2, T/2]$  ( $T > t_0$ ) gives

$$b(t) = \frac{Bt_0}{T} \left[ 1 + 2 \sum_{n=1}^{\infty} \frac{\sin(\omega_n t_0 / 2)}{\omega_n t_0 / 2} \cos \omega_n t \right], \quad \omega_n = 2n\pi / T.$$

Note that the time window of the length  $T$  is used implicitly.  $T$  has to be longer than  $t_0$  and no other constraint is placed on it.

Let  $t_0$  tend to zero under the condition  $Bt_0 = l$ ; in this manner, the Fourier expansion of  $\delta(t)$  is obtained.

$$\delta(t) = \frac{1}{T} \left[ 1 + 2 \sum_{n=1}^{\infty} \cos \omega_n t \right] = \frac{1}{T} \sum_{n=-\infty}^{\infty} e^{i\omega_n t}, \quad \omega_n = 2n\pi / T.$$

The Fourier expansion of the signal in a limited time window implicitly assumes the periodicity of the signal. This formula, when written in exact terms, represents an infinite series of the impulse and it is called as the Dirac comb function. Note that the interval between the delta functions is  $T$ .

Let us change the notation from  $T$  to  $\Delta t$  for the convenience of the following description (Fig. 4).

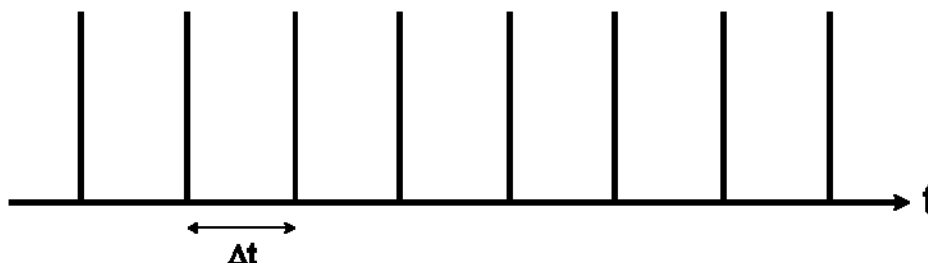


Fig. 4 Dirac comb function.

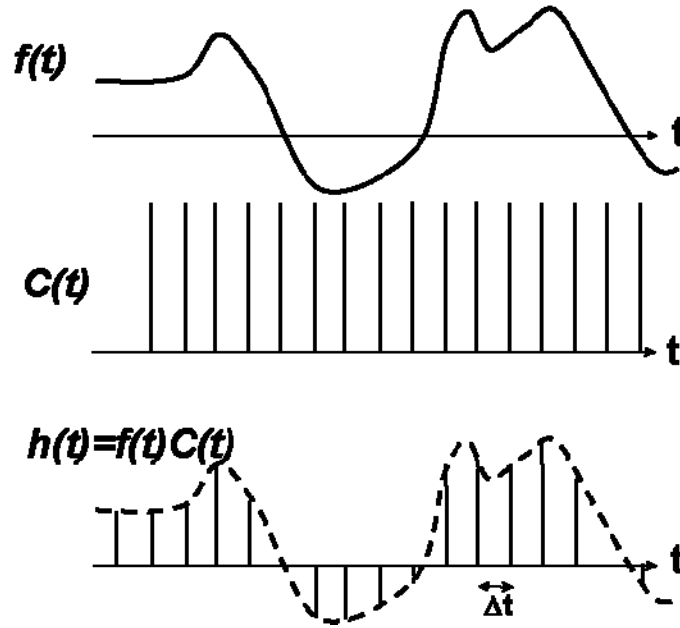


Fig. 5 Continuous function, Dirac comb function, and time series.

### 1.2.2. Time Series Discretization

Today, it is popular and convenient to handle time-dependent data by using computers. For doing this, it is necessary that recorded data is digital or discrete. The process for converting an continuous analog signal  $f(t)$  to its digital equivalent is called “analog-to-digital conversion,” “digitization,” or “discretization.” This process is expressed mathematically as the multiplication of a continuous function  $f(t)$  with the Dirac comb function  $C(t)$ , in which the interval of the neighboring delta function is  $\Delta t$  (Fig. 5).

$$h(t) = f(t) \cdot C(t).$$

### 1.2.3. Folding and Aliasing

From the above, we obtain the Fourier expansion as follows:

$$C(t) = \frac{1}{\Delta t} \sum_{n=-\infty}^{\infty} e^{i(2n\pi/\Delta t)t},$$

where  $\Delta t$  is the sampling interval.

The Fourier transform of the discrete function  $h(t)$  is given by

$$h(\omega) = \int_{-\infty}^{\infty} f(t) \cdot C(t) e^{-i\omega t} dt = \frac{1}{\Delta t} \sum_{n=-\infty}^{\infty} \int_{-\infty}^{\infty} f(t) e^{-i(\omega - 2n\pi/\Delta t)t} dt = \sum_{n=-\infty}^{\infty} F\left(\omega - \frac{2n\pi}{\Delta t}\right).$$

This shows that  $h(\omega)$  is a repetition of  $F(\omega)$  with an interval  $1/\Delta t$  and that a spectra outside the frequency range  $[-1/2\Delta t, 1/2\Delta t]$  does not make any sense (Fig. 6). This border frequency  $1/2\Delta t$  is called the folding frequency or the Nyquist frequency.

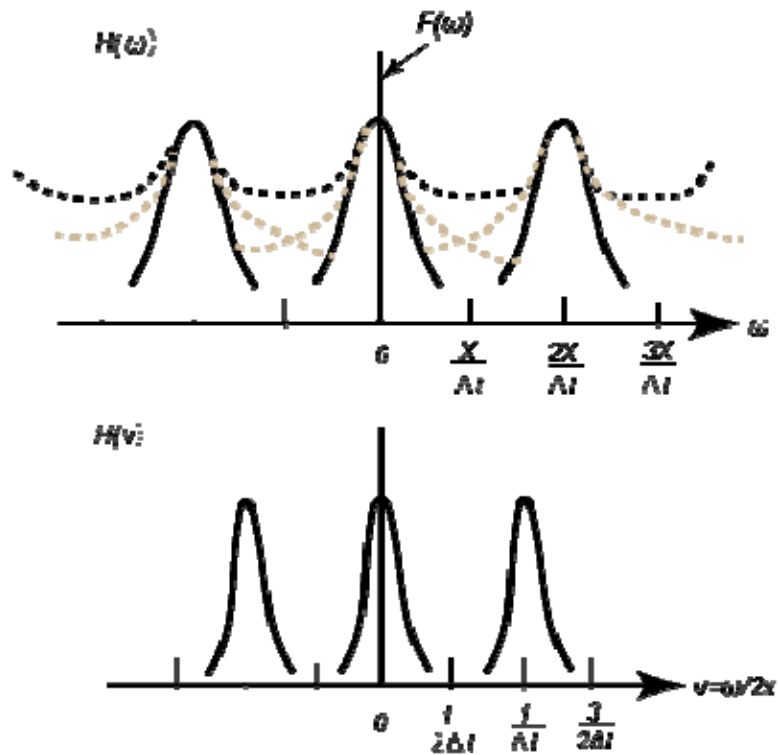


Fig. 6 Spectra of a time series.

The influence of digitization can be minimized easily when the Fourier spectra of the original continuous function have a negligible value at the Nyquist frequency. Otherwise, the foot of the neighboring spectral peak invades in the range  $[-1/2\Delta t, 1/2\Delta t]$  and contaminates the signal. This phenomenon in the frequency domain is referred to as folding. The above consideration suggests that the sampling interval  $\Delta t$  must be shorter than half of the shortest period that is included in the original continuous function. In other words, the frequency components of the period that is shorter than twice the sampling interval must be eliminated before digitization.

The disturbance is more clearly shown in the time domain. Fig. 7 clearly shows that coarse discretization cannot identify fine peaks of the original continuous function, and the result is completely different from the original one.

Folding in the frequency domain and aliasing in the time domain represent the same phenomenon. The relationship between aliasing and folding is schematically shown in Fig. 8.

An analog filter that is applied to the original analog signal in order to prevent the aliasing or folding is called an anti-alias filter. Re-sampling of the digital data also requires the application of an anti-alias filter.



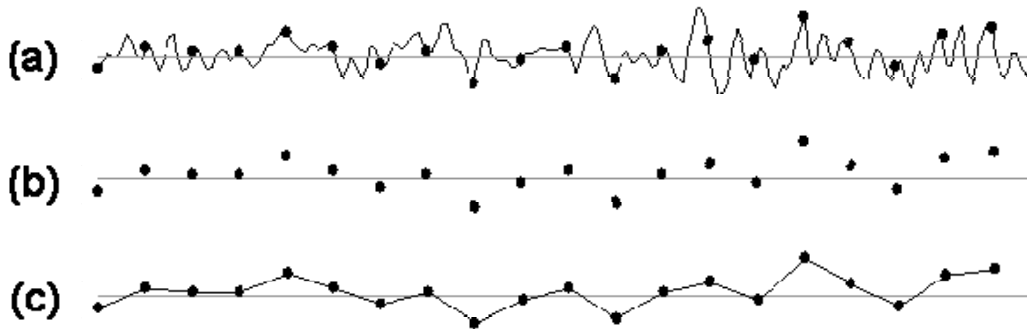


Fig. 7 Example of Aliasing in the time domain. (a) Original analog signal where dots denote the sampling, (b) digitized signal, and (c) Re-constructed analog signal by the linear interpolation of the digitized data. Produced newly based on the concept of Yilmaz(1994).

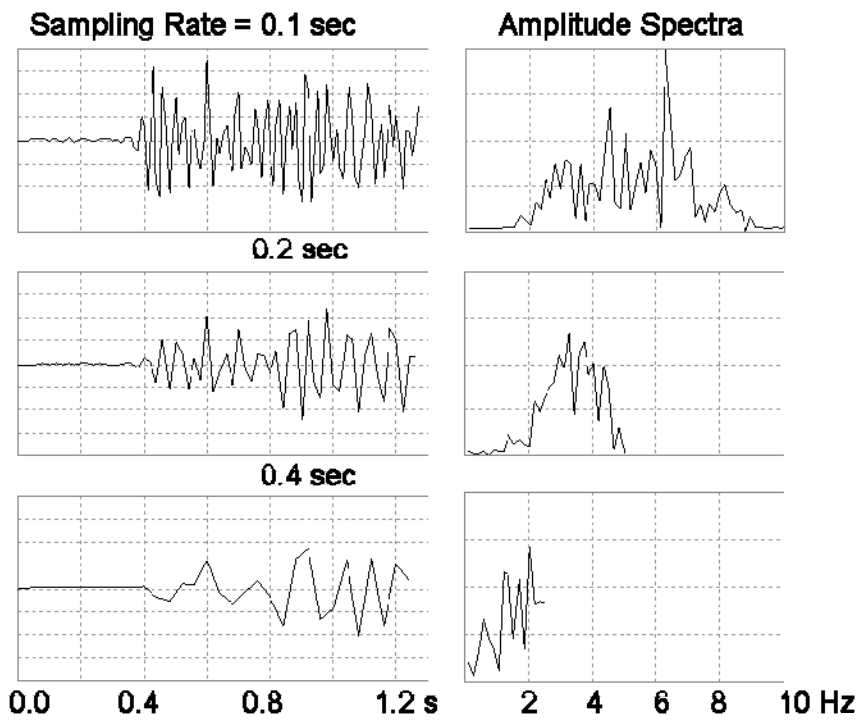


Fig. 8 Example of Folding in the frequency domain. A waveform sampled at 0.1sec has Nyquist frequency of 5.0 Hz. Resampling to 0.2 and 0.4 sec confines the frequency band to 2.5 and 1.25 Hz, respectively. Note the loss of high frequencies at larger sampling intervals. Produced newly based on the concept of Yilmaz(1994).

## 2. Fast Fourier Transform (FFT)

The Fourier transform is one of the basic mathematical tools used for data processing. A signal in the time domain can be converted to one in the frequency domain by applying the Fourier transform and in this manner, different features of the converted data can be obtained. The Fourier transform of digital data is defined in this chapter.

There are several published subroutines of FFT in BASIC, FORTRAN, and C, which are very useful and simplify our task. However, we must focus on the definitions of the Fourier transform and its inverse transform that are given in different books. There are several possible definitions and mathematically they are all equivalent. When we use a subroutine given in a textbook, it is important to carefully read the main text. This lecture note employs the definition of Papoulis (1962, 1984) and Ohsaki (1976).

$$\begin{cases} X(\omega) = \int_{-\infty}^{\infty} x(t) \cdot e^{-i\omega t} dt, \\ x(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} X(\omega) \cdot e^{i\omega t} d\omega. \end{cases}$$

### 2.1. Fourier Expansion of a Finite Time Series

The Fourier expansion of a time series  $x_m$  ( $m = -N/2 + 1, \dots, -1, 0, 1, \dots, N/2$ ) is given as follows.

The coefficient of expansion is given by

$$C_k = \frac{1}{N} \sum_{m=-N/2+1}^{N/2} x_m e^{-i(2\pi km/N)}, \quad k = -N/2+1, \dots, -1, 0, 1, \dots, N/2. \quad (1)$$

By using these coefficients, the original time series  $x_m$  is expanded as

$$x_m = \sum_{k=-N/2+1}^{N/2} C_k e^{i2\pi km/N}, \quad m = -N/2+1, \dots, N/2. \quad (2)$$

Naturally, these formulas imply that any limited time series can be expanded into a finite number of sinusoidal waves of which the frequency is discrete, as shown in Fig. 9. Note that the periodicity in the time domain is implicitly introduced.

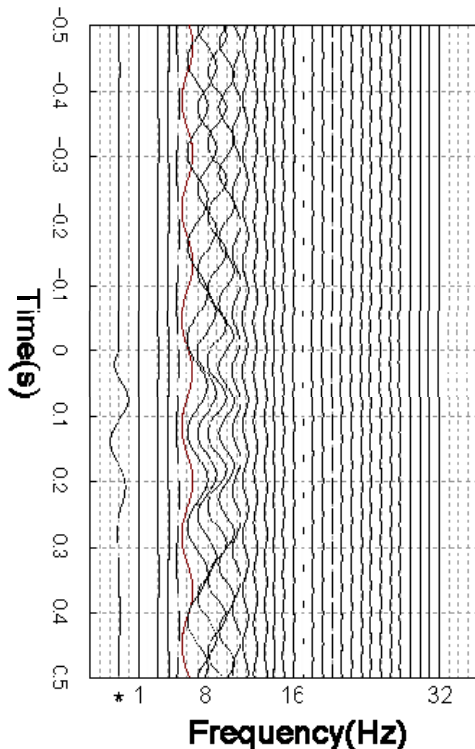


Fig. 9 A series of sinusoidal curves with different frequencies, peak amplitudes, and phase lags can be superimposed to synthesize a waveform on the left-most curve as indicated by the asterisk. The sampling frequency of this waveform is 512Hz. The sinusoidal curves of frequencies higher than 33 Hz are omitted because their amplitudes are negligible. Produced newly based on the concept of Yilmaz(1994).

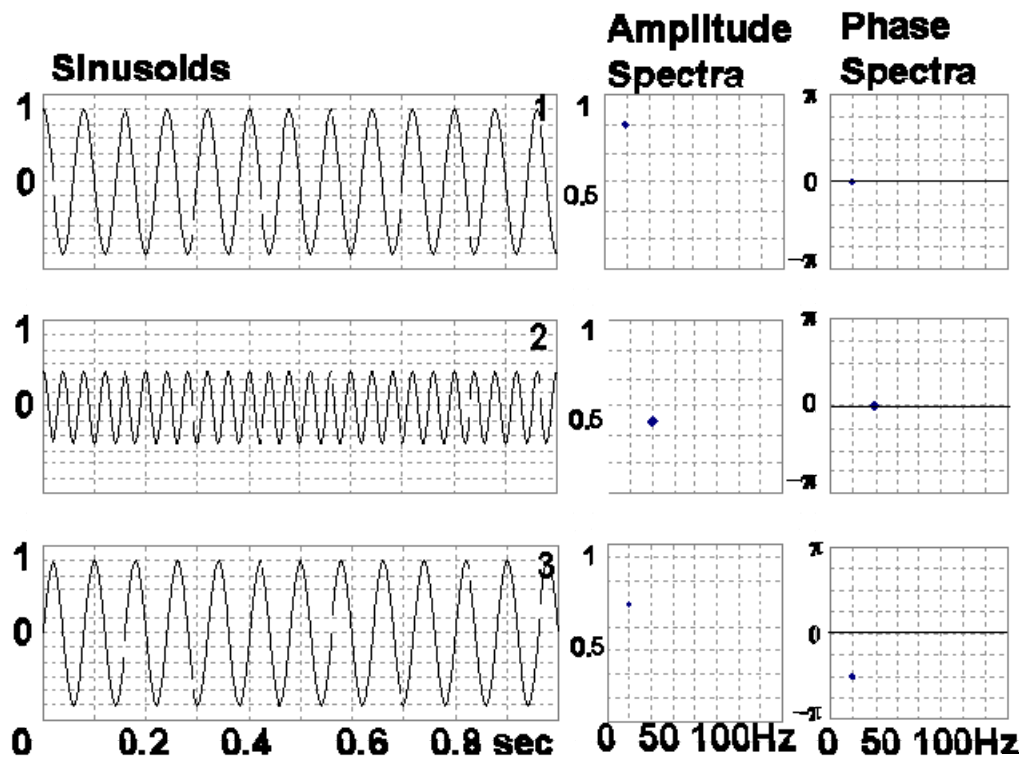


Fig. 10 Three sinusoids (*left*) and their amplitude (*center*) and phase spectra (*right*). The time between two consecutive peaks is the period of the sinusoid, the inverse of which is called frequency. Finally, the time delay of the onset is defined as phase lag. Produced newly based on the concept of Yilmaz(1994).

## 2.2. Sinusoidal Functions and Fourier Spectra of Amplitude and Phase

Yilmaz(1994) shows a persuasive way of explaining Fourier Spectra of Amplitude and Phase. A sinusoidal function is defined by its frequency, amplitude, and time shift, as shown by an example given in Fig. 10. A phase lag, that is, a time shift normalized by the period is usually used. Assume that the phase lag of the signals in the *top* panels is zero and its amplitude is unity. The frequency of signals in the *top* panels is  $12.5\text{ Hz}$ . The *middle* panels have a half amplitude, frequency of  $25.0\text{ Hz}$  and the phase lag is zero. The *bottom* panels have unit amplitude, a frequency of  $12.5\text{ Hz}$ , and phase lag of  $-\pi/2$ .

Every sinusoid drawn in Fig. 9 has a frequency, amplitude, and phase lag. The latter two variables can be plotted against the frequency (Fig. 11). Each point along the amplitude spectrum curve (Fig. 11 *Top*) corresponds to the peak amplitude of the sinusoid at that frequency, as shown in Fig. 9. Note the correspondence of the peak in the amplitude spectra with the high-amplitude frequency range in Fig. 9. Each point along the phase spectrum (Fig. 11 *Bottom right*) corresponds to the phase delay of a peak or trough along the sinusoid at that frequency with respect to the timing line at  $t = 0$  in Fig. 9. Note the correspondence of the phase curve with the trend of a positive peak from trace to trace (Fig. 12).

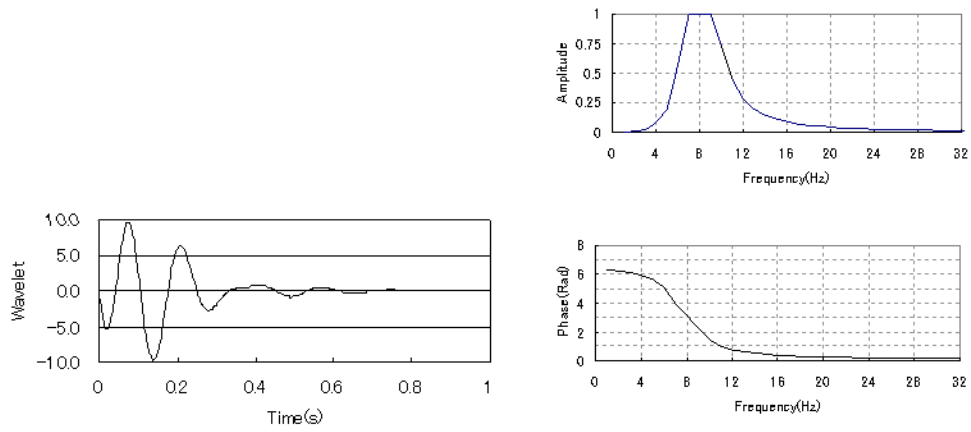


Fig. 11 *Bottom left*: The waveform in Fig.9, its amplitude and phase spectra (top and *bottom right* panels). Produced newly based on the concept of Yilmaz(1994).

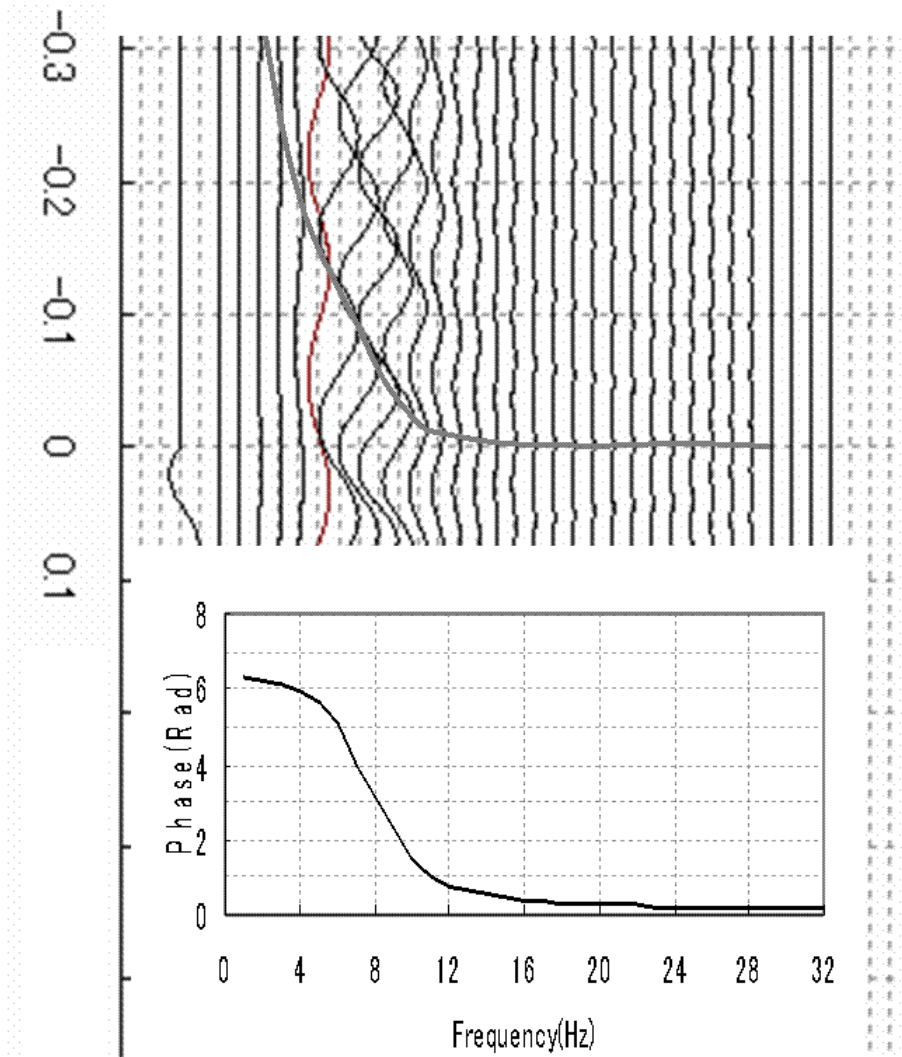


Fig. 12 An enlarged view of Fig. 9 that delineate the trend of the phase curve from curve to curve in comparison of the phase spectra in Fig. 11 *bottom right* panel. Produced newly based on the concept of Yilmaz(1994).

### 2.3. Discrete Fourier Transform

Define the time window length of the time series  $x_m$  as  $T = N\Delta t$ . Eq. (1) can be written as

$$C_k = \frac{1}{N\Delta t} \sum_{m=-N/2+1}^{N/2} x_m \Delta t e^{-i(2\pi km\Delta t / N\Delta t)}, \quad k = -N/2+1, \dots, -1, 0, 1, \dots, N/2.$$

Let  $\Delta t$  tend to zero while  $T$  is kept constant ( $m\Delta t \rightarrow t$ ). This gives a continuous function in a limited time window.

$$C_k = \frac{1}{T} \int_{-T/2}^{T/2} x(t) e^{-i(2\pi kt/T)} dt, \quad -\infty < k \leq \infty, k : \text{discrete}. \quad (3)$$

Similarly,

$$x(t) = \sum_{-\infty}^{\infty} C_k e^{i(2\pi kt/T)}, \quad -T/2 < t \leq T/2. \quad (4)$$

The frequency  $f$  in these formulas is given by  $f = k/T$ . Since  $k$  is an integer, the frequency  $f$  takes a discrete value with an interval  $\Delta f = 1/T$ . Eq. (3) and Eq. (4) show the Fourier expansion of a continuous time-windowed time function. Note again that Eq. (4) shows the repetitive nature of the re-constructed time function.

Change Eq. (4) by using  $\Delta f = 1/T$ .

$$x(t) = \sum_{-\infty}^{\infty} (TC_k) e^{i(2\pi k\Delta f t)} \Delta f, \quad -1/2\Delta f < t \leq 1/2\Delta f.$$

Let  $\Delta f$  tend to zero, i. e., let  $T$  tend to infinity ( $k\Delta f \rightarrow f$ ). This means that the periodicity in Eq. (3) and Eq. (4) becomes eliminated. Eq. (3) gives

$$TC_k = \int_{-\infty}^{\infty} x(t) e^{-i(2\pi ft)} dt, \quad -\infty < k \leq \infty, k : \text{continuous}. \quad (5)$$

$$x(t) = \int_{-\infty}^{\infty} (TC_k) e^{i(2\pi ft)} df = \frac{1}{2\pi} \int_{-\infty}^{\infty} (TC_k) e^{i(2\pi\omega t)} d\omega, \quad -\infty < t \leq \infty. \quad (6)$$

A comparison with the definition of Fourier transform shows that  $(TC_k)$  corresponds to the Fourier transform. Thus, the discrete Fourier transform is given by

$$F(f) = TC_k = \frac{T}{N} \sum_{m=-N/2+1}^{N/2} x_m e^{-i(2\pi km/N)}, \quad (7)$$

$$f = k/T, \quad k = -N/2+1, \dots, -1, 0, 1, \dots, N/2.$$

If the time series is defined in  $[0, T = N\Delta t]$ , the limit of the summation has to be changed to [from  $m = 0$  to  $m = N$ ].

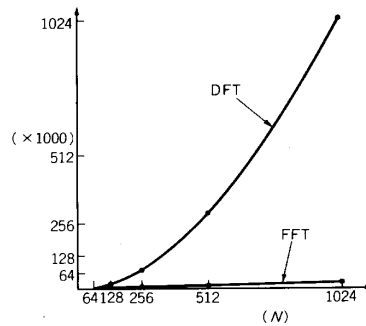


Fig. 13 Comparison of the calculation time of DFT and FFT.

## 2.4. Algorithm for FFT

The calculation of discrete Fourier transform by using Eq. (7) (denoted by DFT) uses considerable time. The time that is necessary for the calculation increases proportionally with  $N^2$ , where  $N$  is the number of samples. For example, the time is 40 s for  $N = 1024$  on a 486 DX2 50 MHz PC.

Fast Fourier transform (FFT) is a technique to compute the Fourier transform of a time series efficiently; this was invented by J. W. Cooley and J. W. Tukey. The calculation time increases proportionally with  $(N/2)\log_2 N$  (Fig. 13).

### 2.4.1. Removal of DC Component and Linear Trend

When the time series has a DC component or a linear trend, the Fourier spectrum cannot be estimated correctly because of the assumption of periodicity. Therefore, it is necessary to remove them before the application of the FFT. It is usually sufficient to remove the straight line connecting the first and last data; however, least square fitting is a recommended method.

Let the discrete time variable  $t_n = n\Delta t$ , and the objective time series  $x_n = x(t_n)$ . The misfit function  $S$  is defined as

$$S = \sum_{n=1}^N (r_n)^2 = \sum_{n=1}^N \{x_n - (at_n + b)\}^2,$$

where  $r_n$  denotes the residual of the fitting;  $a$ , the linear trend; and  $b$ , the DC component. The minimum value of this misfit function is given at

$$\frac{\partial S}{\partial a} = 0, \quad \frac{\partial S}{\partial b} = 0.$$

Thus,

$$\begin{cases} \left( \sum_{n=1}^N n^2 \right) \Delta t \cdot a + \left( \sum_{n=1}^N n \right) \cdot b = \sum_{n=1}^N (nx_n), \\ \left( \sum_{n=1}^N n \right) \Delta t \cdot a + N \cdot b = \sum_{n=1}^N x_n. \end{cases} \quad (8)$$

where the formulas

$$\sum_{n=1}^N n = \frac{N(N+1)}{2}, \quad \text{and} \quad \sum_{n=1}^N n^2 = \frac{N(N+1)(2N+1)}{6},$$

can make the calculation easier.

The coefficients of these linear simultaneous equations  $a$  and  $b$  can be easily obtained.

### 2.4.2. Tapering or Windowing

After removing the DC offset and linear trend, the processed time series begins with one value and ends with another. This causes an unexpected jump or step due to the implicit assumption of periodicity by FFT and results in a bad influence on the estimation of the Fourier transform. A method of preventing such an artificial effect is tapering or windowing, which causes the time series to start with zero and end with zero.

The processing comprises a multiplication with the window function  $w(t)$  in the time domain. The box car window is the simplest solution, but it also has the abovementioned problem.

The tapered window is given as follows:

$$w(t) = \begin{cases} t/t_{taper} & \text{for } 0 \leq t < t_{taper}, \\ 1 & \text{for } t_{taper} \leq t \leq t_{win} - t_{taper}, \\ (t_{win} - t)/t_{taper} & \text{for } t_{win} - t_{taper} < t \leq t_{win}, \\ 0 & \text{for } t_{win} < t. \end{cases} \quad (9.1)$$

The sine tapered window is given as follows:

$$w(t) = \begin{cases} \sin(\pi t/2t_{taper}) & \text{for } 0 \leq t < t_{taper}, \\ 1 & \text{for } t_{taper} \leq t \leq t_{win} - t_{taper}, \\ \sin\{\pi(t_{win} - t)/2t_{taper}\} & \text{for } t_{win} - t_{taper} < t \leq t_{win}, \\ 0 & \text{for } t_{win} < t. \end{cases} \quad (9.2)$$

The tapering time length  $t_{taper}$  is usually approximately one tenth of the time window length  $t_{win}$ .

The following two windowing functions are also used popularly.

Hanning Window:

$$w(t) = \begin{cases} 0.5 - 0.5 \cos(\pi t/t_{taper}) & \text{for } t < t_{taper}, \\ 1 & \text{for } t_{taper} \leq t \leq t_{win} - t_{taper}, \\ 0.5 - 0.5 \cos\{\pi(t_{win} - t)/t_{taper}\} & \text{for } t_{win} - t_{taper} < t \leq t_{win}, \\ 0 & \text{for } t_{win} < t. \end{cases} \quad (9.3)$$

Hamming Window:

$$w(t) = \begin{cases} 0.54 - 0.46 \cos(\pi t/t_{taper}) & \text{for } t < t_{taper}, \\ 1 & \text{for } t_{taper} \leq t \leq t_{win} - t_{taper}, \\ 0.54 - 0.46 \cos\{\pi(t_{win} - t)/t_{taper}\} & \text{for } t_{win} - t_{taper} < t \leq t_{win}, \\ 0 & \text{for } t_{win} < t. \end{cases} \quad (9.4)$$

### 2.4.3. Zero padding

The FFT can be performed efficiently when the number of data  $N$  is  $2^n$  where  $n$  is an integer. Otherwise, zeros must be padded up to the nearest  $2^n$ . Usually, the zeros are padded at the end of the time series. They are not padded at the beginning of the time series, even though this is acceptable theoretically. This is because padding them at the beginning apparently changes the arrival time and causes confusion.

#### 2.4.4. Algorithm

Ohsaki(1976) explained the Algorithm for performing the FFT as a disassembling process.

First, the coefficients of the Fourier expansion of the original time series are given by

$$C_k = \frac{1}{N} \sum_{m=0}^N x_m e^{-i(2\pi km/N)},$$

Disassemble the time series  $x_m$  into two time series in the following manner:

$$\begin{cases} y_m = x_{2m}, & m = 0, 1, 2, \dots, \frac{N}{2} - 1. \\ z_m = x_{2m+1}, & \end{cases}$$

The coefficients of the Fourier expansion of the disassembled time series are

$$\begin{cases} Y_k^{(N/2)} = \frac{2}{N} \sum_{m=0}^{N/2-1} y_m e^{-i[2\pi km/(N/2)]}, \\ Z_k^{(N/2)} = \frac{2}{N} \sum_{m=0}^{N/2-1} z_m e^{-i[2\pi km/(N/2)]}, \end{cases} \quad k = 0, 1, 2, \dots, \frac{N}{2} - 1. \quad (10.1)$$

The original definition becomes

$$\begin{aligned} C_k^{(N)} &= \frac{1}{N} \sum_{m=0}^N x_m e^{-i(2\pi km/N)} \\ &= \frac{1}{N} \left\{ \sum_{m=0}^{N/2-1} y_m e^{-i(2\pi k(2m)/N)} + \sum_{m=0}^{N/2-1} z_m e^{-i(2\pi k(2m+1)/N)} \right\} \\ &= \frac{1}{N} \sum_{m=0}^{N/2-1} y_m e^{-i(2\pi k(2m)/N)} + e^{-i[\pi k/(N/2)]} \frac{1}{N} \sum_{m=0}^{N/2-1} z_m e^{-i(2\pi km/(N/2))} \\ &= \frac{1}{2} Y_k^{(N/2)} + \frac{1}{2} e^{-i[\pi k/(N/2)]} Z_k^{(N/2)}, \end{aligned} \quad (10.2)$$

for  $k = 0, 1, 2, \dots, N/2 - 1$ .

Replace  $k$  in Eq. (10.1) with  $k + N/2$ ; then

$$\begin{cases} Y_{k+N/2}^{(N/2)} = \frac{2}{N} \sum_{m=0}^{N/2-1} y_m e^{-i[2\pi(k+N/2)m/(N/2)]} = \frac{2}{N} \sum_{m=0}^{N/2-1} y_m e^{-i[2\pi km/(N/2)+2\pi m]} \\ \quad = \frac{2}{N} \sum_{m=0}^{N/2-1} y_m e^{-i[2\pi km/(N/2)]} = Y_k^{(N/2)}, \\ Z_{k+N/2}^{(N/2)} = \frac{2}{N} \sum_{m=0}^{N/2-1} z_m e^{-i[2\pi(k+N/2)m/(N/2)]} = \frac{2}{N} \sum_{m=0}^{N/2-1} z_m e^{-i[2\pi km/(N/2)+2\pi m]} \\ \quad = \frac{2}{N} \sum_{m=0}^{N/2-1} z_m e^{-i[2\pi km/(N/2)]} = Z_k^{(N/2)}, \end{cases}$$

for  $k = 0, 1, 2, \dots, \frac{N}{2} - 1$ .

Eq. (10.2) gives



$$\begin{aligned}
C_{k+N/2}^{(N)} &= \frac{1}{2} Y_{k+N/2}^{(N/2)} + \frac{1}{2} e^{-i[\pi(k+N/2)/(N/2)]} Z_{k+N/2}^{(N/2)} \\
&= \frac{1}{2} Y_{k+N/2}^{(N/2)} + \frac{1}{2} e^{-i[\pi k/(N/2)]} e^{-i\pi} Z_{k+N/2}^{(N/2)} \\
&= \frac{1}{2} Y_k^{(N/2)} - \frac{1}{2} e^{-i[\pi k/(N/2)]} Z_k^{(N/2)}.
\end{aligned}$$

This last change is due to Euler's formula. Therefore,

$$\begin{cases} 2C_k^{(N)} = Y_k^{(N/2)} + e^{-i[\pi k/(N/2)]} Z_k^{(N/2)}, \\ 2C_{k+N/2}^{(N)} = Y_k^{(N/2)} - e^{-i[\pi k/(N/2)]} Z_k^{(N/2)}. \end{cases} \quad k = 0, 1, 2, \dots, \frac{N}{2} - 1 \quad (10.3)$$

This process clearly shows that the coefficients of the Fourier expansion of the original time series  $x_m$  can be easily given by the coefficients of the Fourier expansion of the two time series  $y_m$  and  $z_m$  obtained by disassembling.

By applying the same process repetitively,  $N$  time series, each with only one sample, are obtained. The coefficient of the Fourier expansion of the time series having only one sample is the sample itself, as shown by

$$C_0^{(1)} = \frac{1}{1} \sum_{m=0}^0 x_m e^{-i(2\pi km/1)} = x_0. \quad (10.4)$$

Thus, the coefficients of the Fourier expansion of the original time series  $x_m$  can be obtained by using Eq. (10.3) repetitively. Let us check the process by using an example of a time series of 8 samples.

### Example

Consider a time series of 8 samples, as shown in Table 1.

Table 1 (After Ohsaki(1976))

$m$	0	1	2	3	4	5	6	7
$x_m$	5	32	38	-33	-19	-10	1	-8

The first disassembling gives two series of 4 samples, as shown in Table 2. The second disassembling gives four time series of 2 samples, as shown in Table 3. The third disassembling gives eight series that has only one sample, as shown in Table 4.

Table 2 (After Ohsaki(1976))

$m$	0	1	2	3
$y_m$	5	38	-19	1
$z_m$	32	-33	-10	-8

Table 3 (After Ohsaki(1976))

$m$	0	1
$y_m'$	5	-19
$Z_m'$	38	1
$y_m''$	32	-10
$z_m''$	-33	-8

Table 4 (After Ohsaki(1976))

$m$	0
$y_m'''$	5
$Z_m'''$	-19
$y_m''''$	38
$z_m''''$	1
$y_m'''''$	32
$Z_m'''''$	-10
$y_m''''''$	-33
$Z_m''''''$	-8

Table 5. (After Ohsaki(1976))

$k$	0
$y_k'''$	5
$Z_k'''$	-19
$y_k''''$	38
$z_k''''$	1
$y_k'''''$	32
$Z_k'''''$	-10
$y_k''''''$	-33
$Z_k''''''$	-8

Table 6 (After Ohsaki(1976))

$k$	0	1
$Y_k'$	-7.0	12.0
$Z_k'$	19.5	18.5
$Y_k''$	11.0	21.0
$Z_k''$	-20.0	-12.5

The coefficients of the Fourier expansion of these 8 series of only one sample are given in Table 5. These are the same series as shown in Table 4.

By using relations such that

$$e^{-i0} = 1.0,$$

$$e^{-i[\pi/4]} = 0.7071 - 0.7071i,$$

$$e^{-i[\pi/2]} = -1.0i,$$

$$e^{-i[3\pi/4]} = -0.7071 - 0.7071i,$$

the coefficients of the Fourier expansion of the four time series of two samples in Table 3 are given in Table 6.

The coefficients of the Fourier expansion of the two time series of four samples in Table 2 are given in Table 7.

Table 7 (After Ohsaki(1976))

$k$	0	1	2	3
$y_k$	6.25	6.00 - 9.25i	-13.25	6.00 + 9.25i
$Z_k$	-4.75	10.50 + 6.25i	15.75	10.50 - 6.25i

Finally, the coefficients of the Fourier expansion of the original time series in Table 1 are given in Table 8.

Table 8 (After Ohsaki(1976))

$k$	0	1	2	3	4	5	6	7
$C_k$	0.75	8.922 - 6.128i	6.625 - 7.875i	-2.922 + 3.122i	5.5	-2.922 - 3.122i	-6.625 + 7.875i	8.922 + 6.128i

For the time series having a large number of samples, disassembling process will take time. The special feature, however, can shorten the disassembling process drastically. A comparison of Table 1 and Table 4 is shown in Table 9. The order numbers  $m'$  (binary) in Table 4 are completely *bit-reversed* ones of those in Table 1. This provides an efficient strategy to obtain pivoted time series like those in Table 4.

Table 9 (After Ohsaki(1976))

Table 1			Table 4			
$x_m$	$m$	$m$ (binary)	$m'$ (binary)	$m'$	$(X_m)$	$m$
5	0	000	000	0	5	0
32	1	001	100	4	-19	1
38	2	010	010	2	38	2
-33	3	011	110	6	1	3
-19	4	100	001	1	32	4
-10	5	101	101	5	-10	5
1	6	110	011	3	-33	6
-8	7	111	111	7	-8	7

Express the index  $m$  of  $x_m$  in binary, and then reverse their bit order to obtain the pivoted index  $m'$ . The pivoted time series  $X_m$  is given by  $X_m = x_{m'}$ . Begin the process to obtain the coefficients of a Fourier expansion by using Eq. (10.3).

### 2.4.5. Time Window and Periodicity

As mentioned above, the FFT can be applied only to a time series within a limited range of time. This time window of finite length can affect the estimation of the Fourier spectra. Consider the following time window defined in  $[-T/2, T/2]$ :

$$B(t) = \begin{cases} 1, & |t| \leq T/2, \\ 0, & |t| > T/2, \end{cases}$$

the Fourier transform of which is given by

$$B(\omega) = \int_{-\infty}^{\infty} B(t) e^{-i\omega t} dt = \int_{-T/2}^{T/2} B(t) e^{-i\omega t} dt = \left[ \frac{1}{-i\omega} e^{-i\omega t} \right]_{-T/2}^{T/2} = \frac{2}{\omega} \sin\left(\frac{\omega T}{2}\right),$$

as shown in Fig. 14. Note that this depends on the window length  $T$ .

Applying the time window  $B(t)$  to the original function  $x_0(t)$  implies the product of these two time functions:

$$x(t) = B(t) \cdot x_0(t).$$

The mathematically defined Fourier transform of the time windowed function  $x(t)$  is given by

$$x(\omega) = B(\omega) * x_0(\omega) = \frac{2}{\omega} \sin\left(\frac{\omega T}{2}\right) * x_0(\omega),$$

$$x(\omega) = B(\omega) * x_0(\omega) = \frac{2}{\omega} \sin\left(\frac{\omega T}{2}\right) * x_0(\omega),$$

where  $\omega$  is the angular frequency and  $*$  denotes convolution. Since  $B(\omega)$  depends on  $T$ ,  $x(\omega)$  also depends on  $T$ . The length of the time window can affect the result of the estimation of the Fourier transform for a time windowed function.

As mentioned previously, the coefficients of the Fourier expansion for a time series of a finite length  $C_k$  implicitly satisfy the assumption of periodicity outside the time window, whereas the discrete Fourier transform  $TC_k$  assumes zeros outside of the time window.

For a sinusoid that is time windowed by the same time length as its period multiplied by an integer, the coefficients of the Fourier expansion  $C_k$  is not affected by the length of the time window, whereas the discrete Fourier transform  $TC_k$  changes its value depending on  $T$ .

In contrast, for an impulse function, the coefficients of the Fourier expansion

$$C_k = I/N\Delta t = I/T,$$

changes, whereas

$$TC_k = 1,$$

does not depend on  $T$ .

Since the effect of the time window length is an artifact, it is better to select a measure that is not influenced by the window length in order to estimate the Fourier spectra of a given time series. The examples explained above suggest that the coefficients of the Fourier expansion  $C_k$  is a good measure for time series assumed to be a digitized part of a periodic function, because the feature of the original time

function coincides with the assumption accompanying  $C_k$ . Further, it is suggested that the discrete Fourier transform  $TC_k$  is a better measure of the estimation of the Fourier spectra of an impulse function.

However, actual seismic signals, are transient and neither periodic nor impulsive. Thus, there is an ambiguity with respect to the selection of a measure for estimating the frequency components of time series, that is, a time windowed and discretized function. It is important to recognize these characteristics of a discrete Fourier transform and the coefficients of Fourier expansion and to select an appropriate one for each problem.

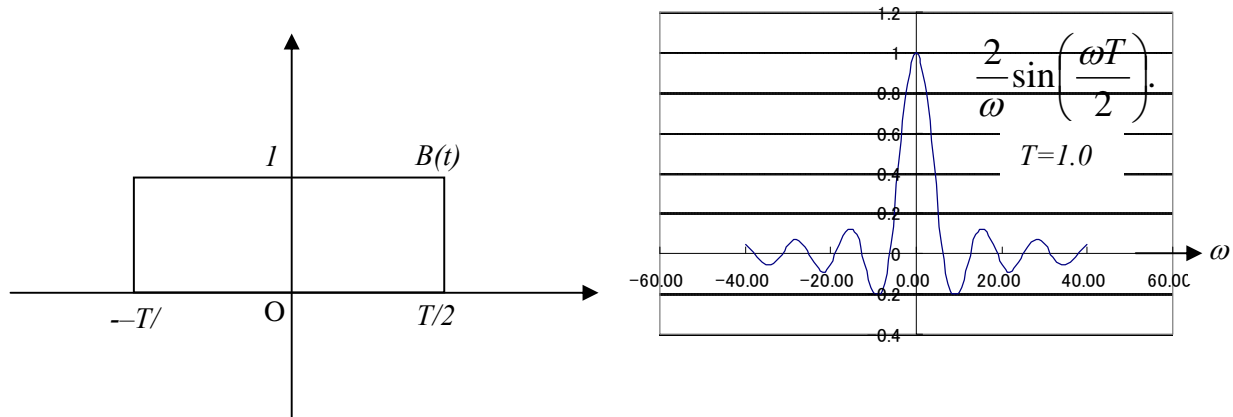


Fig. 14 Boxcar function and its amplitude spectra.

#### 2.4.6 Spectral Smoothing

The Fourier spectra of real seismograms deviate considerably. Since the result of the FFT analysis are obtained for constantly sampled frequencies, the deviation is emphasized at higher frequency ranges. If plotted on a full logarithmic chart, the high frequency portion is almost completely painted black. This makes it difficult to observe a general tendency. Moreover, they occasionally take very small values. This causes instability of the spectral ratio, simply when required. Therefore, the spectral smoothing techniques are applied widely.

In order to plot them on a linear logarithmic chart, a simple moving average over the frequency works well typically.

$$Y(f_i) = \sum_{\substack{j < i+u \\ j > i-u}} y(f_j),$$

where  $u$  denotes half bandwidth.

The following are examples for the weighted moving average that are applied repeatedly until the processed spectra become sufficiently smooth.

$$Y(f_i) = 0.25y(f_{i-1}) + 0.5y(f_i) + 0.25y(f_{i+1}),$$

$$Y(f_i) = 0.23y(f_{i-1}) + 0.54y(f_i) + 0.23y(f_{i+1}).$$

The weight coefficients can be given in the form of specially selected functions  $w(f)$ . For example,

$$w(f) = u \cdot \left( \sin \frac{\pi f}{2} / \frac{\pi f}{2} \right)^2 \quad : \text{Bartlett window}$$

$$w(f) = 0.75u \cdot \left( \sin \frac{\pi f}{2} / \frac{\pi f}{2} \right)^4 \quad : \text{Parzen window}$$

$$w(f) = a \left[ \sin \left\{ b \log_{10} (f / f_c) \right\} / b \log_{10} (f / f_c) \right]^4 \quad : \text{Logarithmic window}$$

There is a trade-off relation between the capacity of smoothing techniques to a stabilizing spectra and the resolution of the processed spectra. Thin peaks may be smoothed out and diminished by efficient smoothing. Insufficient smoothing cannot be used to show the general features of spectra. The objectives of smoothing are not achieved in the both extreme cases. The only way to find an appropriate smoothing technique is the trial-and-error approach.

## 2.5. Practice for FFT

The topics in this chapter can be understood more easily if the distributed programs are used for practice. In the following pages, the coefficients of the Fourier expansion obtained by FFT, amplitude spectra, and phase spectra are calculated for various test signals. *Ghost View* and *Ghost Script* can draw G.PS on the computer.

The following six programs have been prepared for practice:

<i>TESTSIG.EXE</i>	prepares test signals such as cosine or sine function, impulse, etc.
<i>PTIME.EXE</i>	plots the signal.
<i>FFT.EXE</i>	calculates Fourier coefficients by using FFT.
<i>PCFFT.EXE</i>	plots raw coefficients of FFT.
<i>PSPEC.EXE</i>	plots discrete Fourier transforms (Fourier spectra),
<i>IFFT.EXE</i>	calculate the inverse Fourier transform from data given in the frequency domain.

- (1) Assume that *UT* is a filename for the time series  $u(t)$  and *UF* its FFT coefficients  $U(f)$ . First make the file *UT* by using *TESTSIG*. The output from *TESTSIG* is *UT*.
- (2) Draw *UT* in the file *G.PS* by using *PTIME*. The input file name for *PTIME* is *UT* and the output file name is *G.PS*.
- (3) Calculate the coefficients of the Fourier expansion for the time series stored in the file *UT* by using *FFT*. The input file is *UT* for *FFT* and the output file is *UF*.
- (4) Draw the coefficients of the Fourier expansion stored in the file *UF* by using *PCFFT*. The input file is *UF* for *PCFFT* and the output is *G.PS*.
- (5) Draw the Fourier spectra for the data stored in the file *UF* by using *PSPEC*. The input file is *UF* for *PSPEC* and the output is *G.PS*.

The data in the file *UT* consist of number of data  $N$  and the sampling interval  $\Delta t$  or the frequency interval  $\Delta f$  in the header, followed by data in “one-data-a-line” format. The programs are prepared separately so that you could use each of them as a basic tool of data processing.

### 2.5.1. Cosine and Sine Wave

#### Exercise: Cosine wave

Compute the FFT of a cosine wave with  $\Delta t = 1.0$  s,  $N = 32$ , period = 16.0 s, amplitude = 10.0, phase = 0.0, and damping = 0.0.

Draw the time series, the FFT coefficients, and Fourier spectra.

Note that the original cosine wave is decomposed into two cosine waves of a half amplitude having positive and negative frequencies as

$$u(t) = 5.0e^{i2\pi ft} + 5.0e^{-i2\pi ft} = 10.0\cos 2\pi ft. \quad (11.1)$$

Note that the calculated FFT coefficients have these values, and the Fourier spectra has the value of  $5.0 \times 32.0 = 160.0$

Repeat the same procedure but with  $N = 64$  and check the amplitude of the Fourier spectra and FFT coefficients.

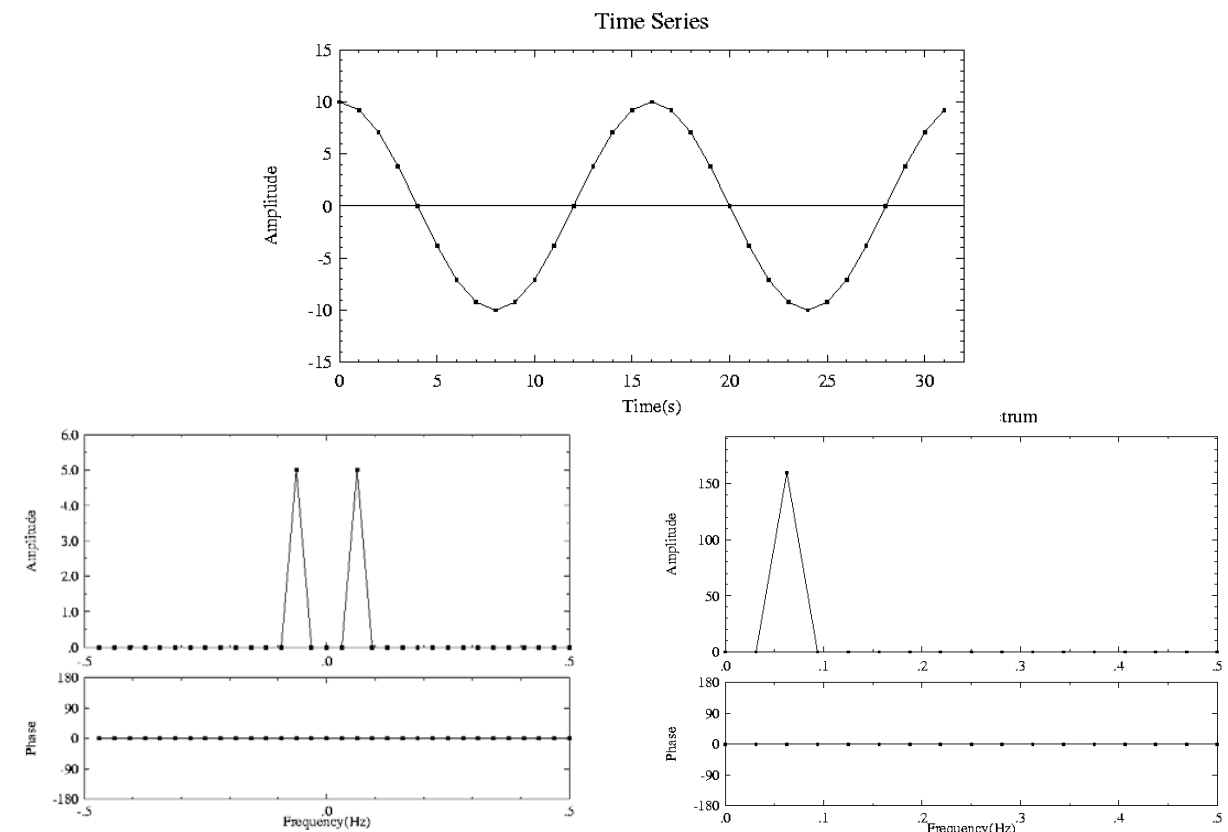


Fig. 15 Time series (*top*), coefficients of FFT (*left bottom*), and discrete Fourier spectra (*right bottom*) of the given time series, i.e., a cosine function.



**Exercise: Sine wave**

Calculate the FFT of a sine wave by making phase = 90.0 (that is, the progress of a phase in deg.) with  $\Delta t = 1.0$  s,  $N = 32$ , period = 16.0 s, amplitude = 10.0, and damping = 0.0.

Draw the time series, FFT coefficients, and Fourier spectra.

Note that the phase  $\phi$  for the positive and negative frequencies have opposite signs, i. e.,

$$\begin{aligned}
 u(t) &= 5.0e^{i\phi}e^{i2\pi ft} + 5.0e^{-i\phi}e^{-i2\pi ft} = 5.0e^{i(2\pi ft + \phi)} + 5.0e^{-i(2\pi ft + \phi)} \\
 &= 10.0\cos(2\pi ft + \phi) = -10.0\sin 2\pi ft.
 \end{aligned}
 \tag{11.2}$$

The last change is due to  $\phi = 90.0$  degrees. A complex conjugate relation of the FFT coefficients for the negative frequencies with those for the corresponding positive ones are required to ensure that the original time series is real.

Note again that the calculated FFT coefficients have such values, the Fourier spectra has the value of  $5.0 \times 32.0 = 160.0$ , and the phase is 90.0 degrees.

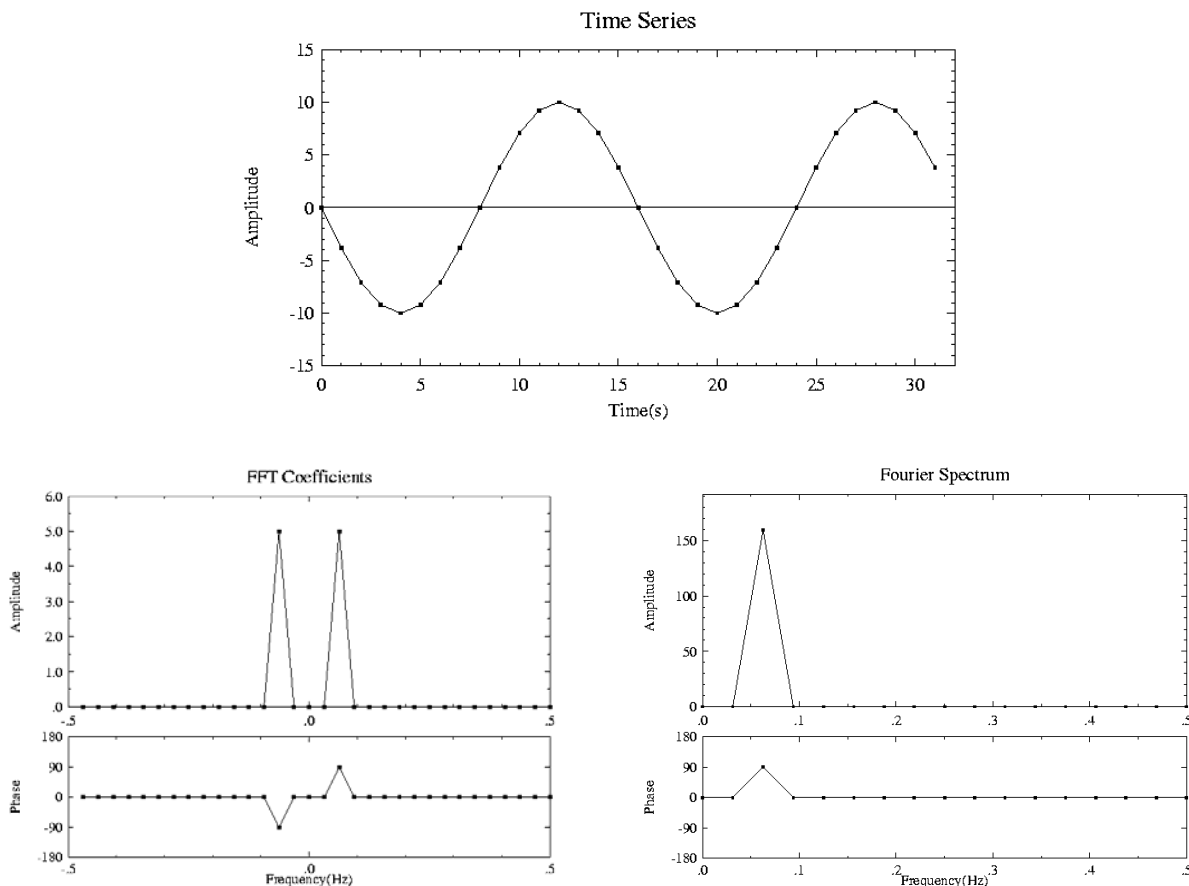


Fig. 16 Time series (*top*), coefficients of FFT (*left bottom*), discrete Fourier spectra (*right bottom*) of the given time series, i. e., a sine function.

### Exercise: Cosine wave at the Nyquist frequency

Calculate the FFT of a cosine wave by using phase = 0.0 with  $\Delta t = 1.0$  s,  $N = 32$ , amplitude = 10.0, and damping = 0.0, and with the period corresponding to the Nyquist frequency  $f_{Nyquist} = 1/2\Delta t = 0.5$  Hz.

Draw the time series, FFT coefficients, and Fourier spectra.

Note that the coefficient for  $f_{Nyquist} = 10.0$ , i. e., it does not share the amplitude with the coefficient for  $-f_{Nyquist}$ .

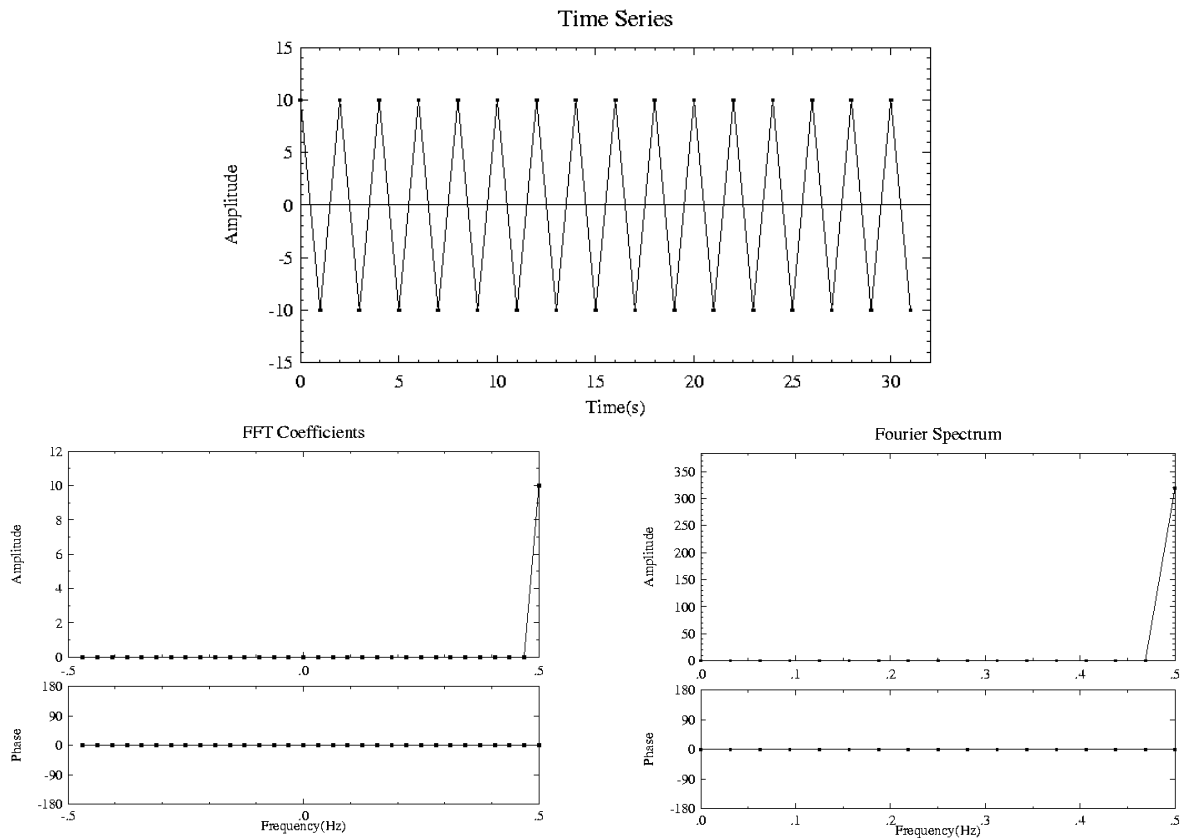


Fig. 17 Time series (*top*), coefficients of FFT (*left bottom*), discrete Fourier spectra (*right bottom*) of the given time series, i. e., a cosine function at the Nyquist frequency.

### Exercise: Summation

Calculate the FFT of two cosine waves with ( $T = 16.0$  s,  $A = 10.0$ ) and ( $T = 2.0$  s,  $A = 10$ ). Other parameters are common, i. e.,  $\Delta t = 1.0$  s,  $N = 32$ , amplitude = 10, phase = 0, and damping = 0.0.

Draw the time series, FFT coefficients, and Fourier spectra for each cosine wave and for the summed one. Additivity is one of the basic characteristics of Fourier transforms.

$$u(t) + v(t) = \int_{-\infty}^{\infty} [U(f) + V(f)] e^{-2\pi i f t} df. \quad (11.3)$$

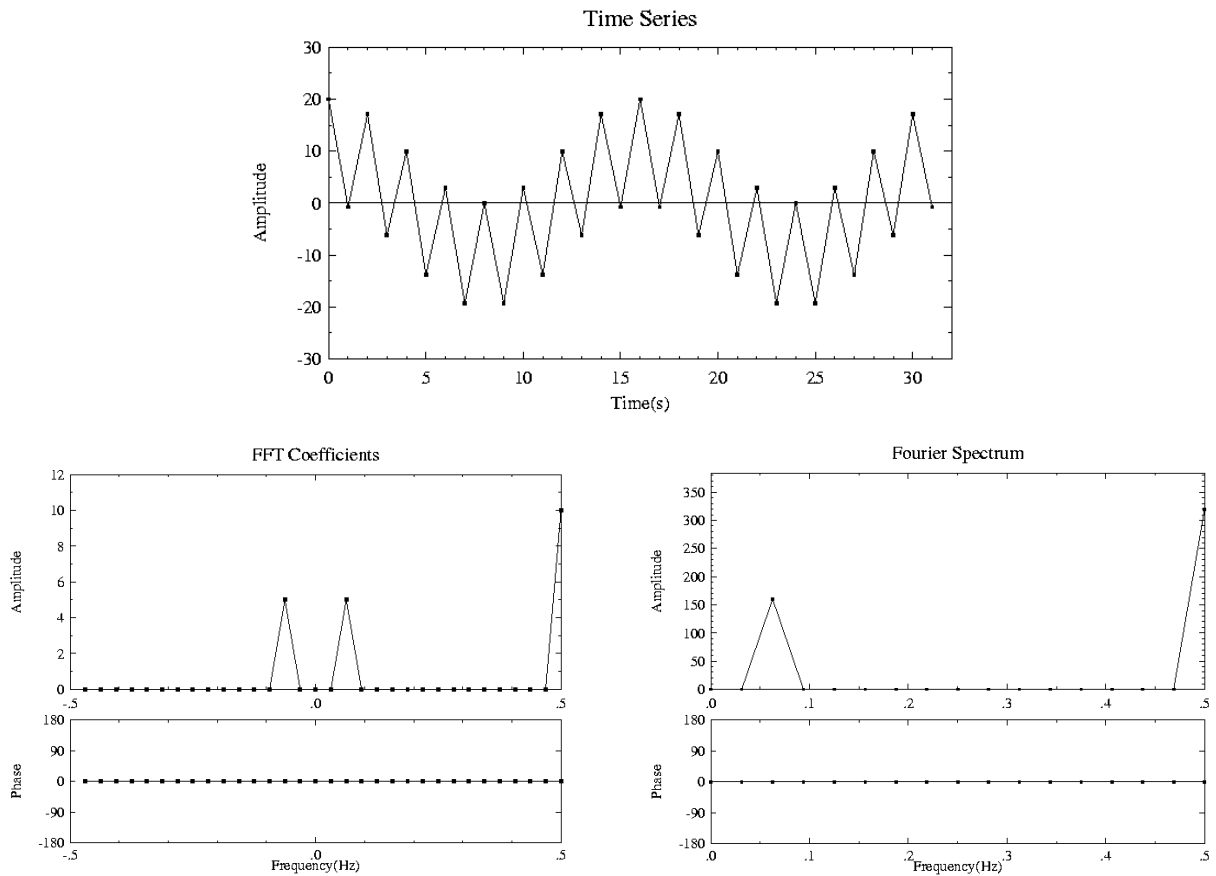


Fig. 18 Time series (*top*), coefficients of FFT (*left bottom*), discrete Fourier spectra (*right bottom*) of the given time series, i. e., two superposed cosine functions with different frequencies.

### 2.5.2. Constant

#### Exercise: Constant

Calculate the FFT for a constant with the constant = 10.0 and  $\Delta t = 1.0$  s,  $N = 32$ .

Draw the time series, FFT coefficients, and Fourier spectra.

Note that the amplitude of the Fourier spectra at zero frequency is

$$TC_k = N\Delta t C_k = 32 \times 1.0 \times 10.0 = 320.0,$$

whereas the FFT coefficient at zero frequency is  $C_k = 10.0$ .

The FFT coefficient and Fourier spectra at the zero frequency correspond to the DC component of the time series.

Repeat the same procedure but with  $N = 64$  and check the amplitude of the Fourier spectra and FFT coefficients.

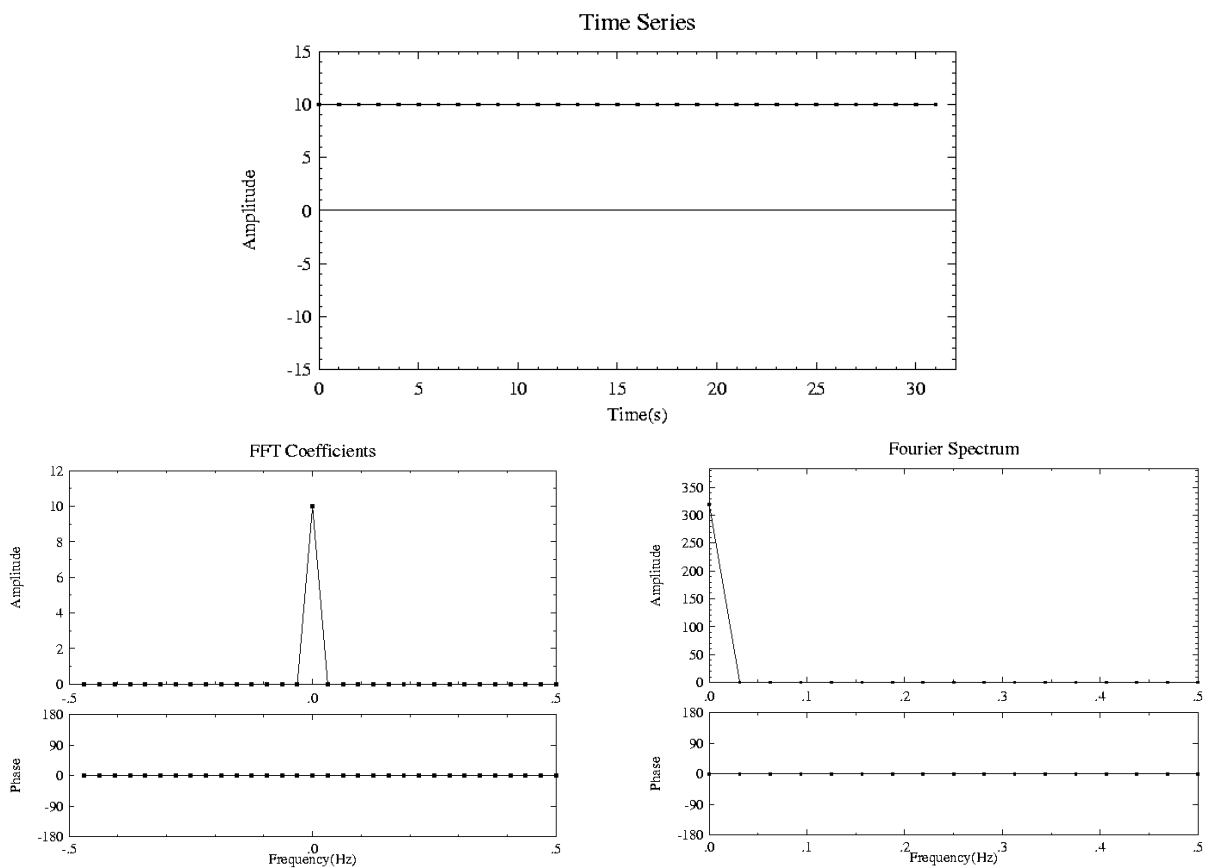


Fig. 19 Time series (*top*), the coefficients of FFT (*left bottom*), the discrete Fourier spectra (*right bottom*) of the given time series, i. e., a constant function.

### 2.5.3. Impulse

#### Exercise: Calculate the FFT for an impulse at $t = 0$ .

Calculate the FFT for an impulse of amplitude = 10.0 with  $\Delta t = 1.0$  s,  $N = 32$ , amplitude = 10.0, and location = 0.0.

Draw the time series, FFT coefficients, and Fourier spectra.

For a “continuous-infinite” case, the Fourier transform of a Dirac delta function

$$\delta(t) = \begin{cases} +\infty & t = 0 \\ 0 & t \neq 0 \end{cases}, \quad \int_{-\infty}^{\infty} \delta(t) dt = 1 \quad (11.4)$$

is known to be unity.

$$\int_{-\infty}^{\infty} \delta(t) e^{-i2\pi ft} dt = 1.$$

Therefore, the mathematical Fourier transform of an impulse, i. e., the delta function at  $t = 0$  is a constant and has zero phase.

Note that the amplitude of the Fourier spectra, which is constant for all frequencies, coincides with that of the impulse = 10.0, whereas the FFT coefficients have the value of  $A_{impulse}/N = 10.0/32 = 0.3125$ .

Repeat the same procedure but with  $N = 64$  and check the amplitude of the Fourier spectra and FFT coefficients.

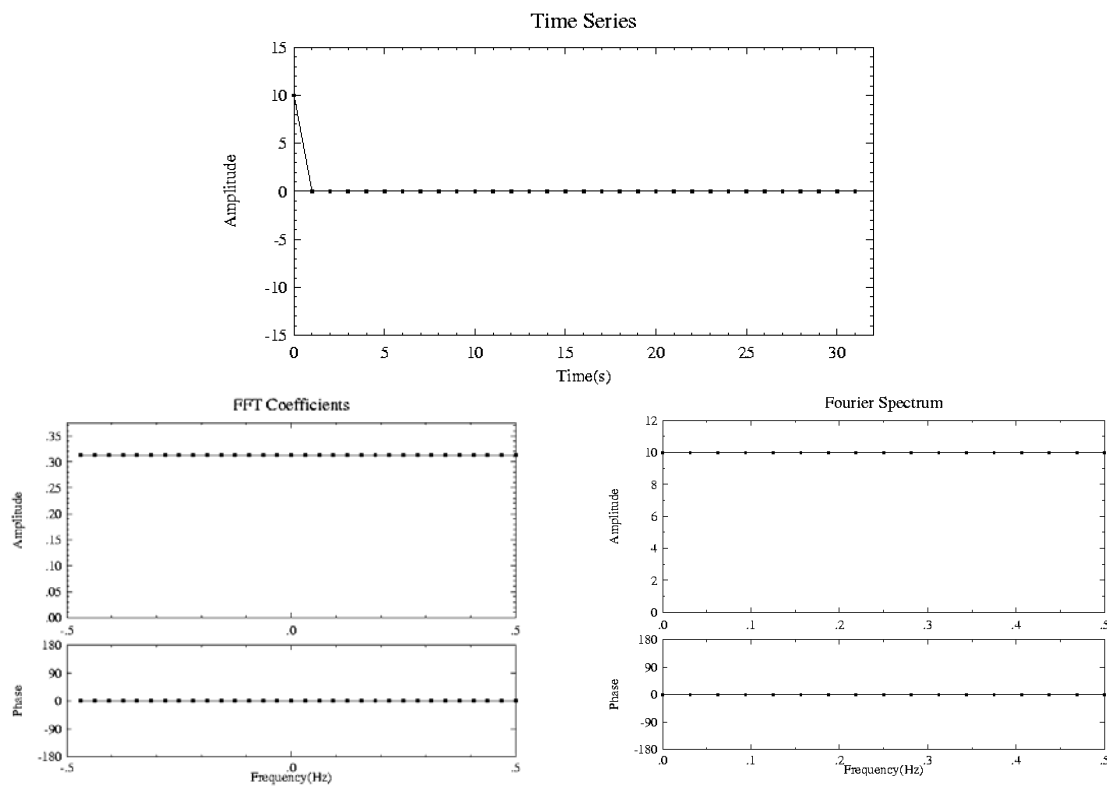


Fig. 20 Time series (*top*), coefficients of FFT (*left bottom*), discrete Fourier spectra (*right bottom*) of the given time series, i. e., an impulse at  $t = 0.0$ .

### 2.5.4. Time Shift

#### Exercise: Compute FFT of an impulse at $t = 2.0, 4.0, 6.0, \dots$

Calculate the FFT of an impulse of amplitude = 10.0 with  $\Delta t = 1.0$  s,  $N = 32$ , amplitude = 10.0, and location = 2.0, 4.0, 6.0...

Draw the time series, FFT coefficients, and Fourier spectra.

The Fourier transform of a time-shifted signal  $u(t - \tau)$  is known to be as follows:

$$\int_{-\infty}^{\infty} u(t - \tau) e^{-i2\pi ft} dt = e^{-i2\pi f\tau} U(f), \quad (11.5)$$

which introduces  $2\pi f\tau$  phase shifts to  $U(f)$ .

The Fourier transform of a time-shifted signal is  $u(t - \tau)$  is known to be as follows:

$$\int_{-\infty}^{\infty} u(t - \tau) e^{-i2\pi ft} dt = e^{-i2\pi f\tau} U(f), \quad (11.6)$$

which introduces  $2\pi f\tau$  phase shifts to  $U(f)$ .

Note that there is no difference with respect to the amplitude Fourier spectra and FFT coefficients among the shifted impulses, and that the slope of the phase spectra increases with the time shift given to the impulse in the time domain.

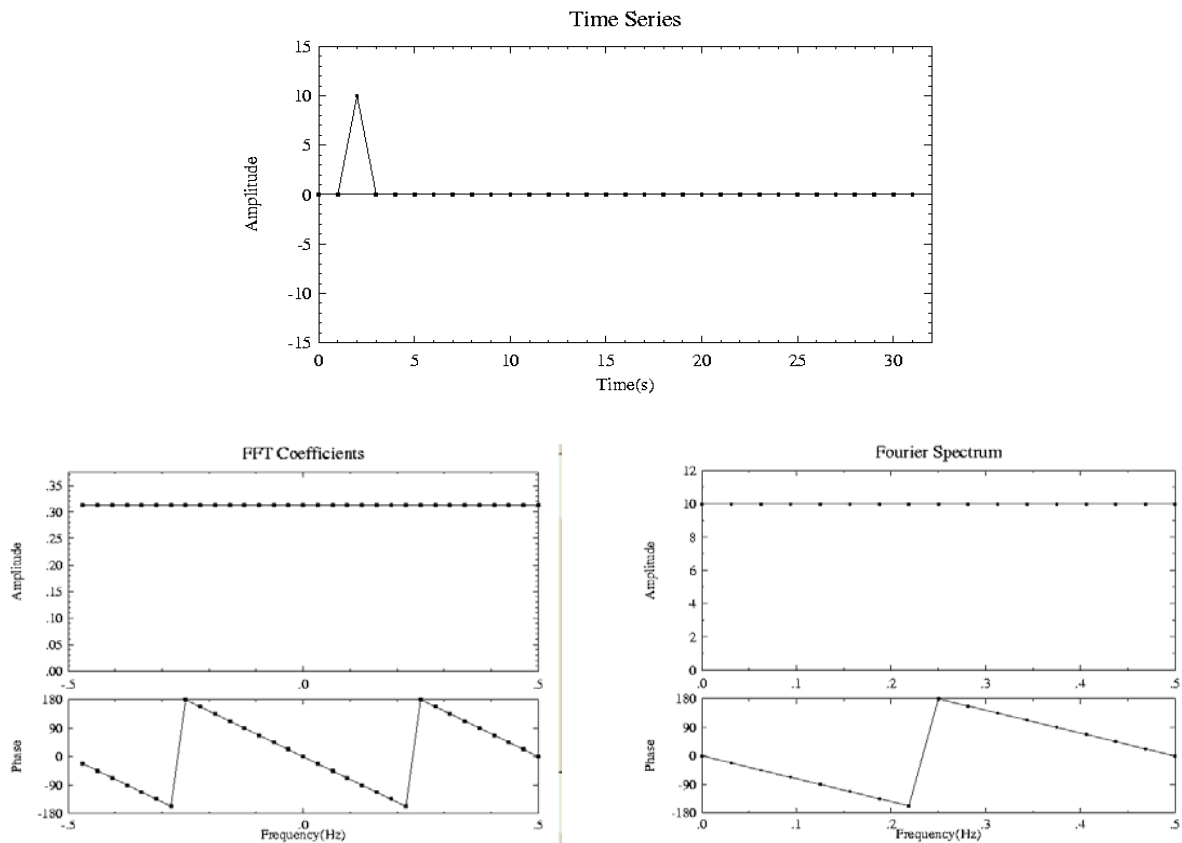


Fig. 21.1 Time series (*top*), coefficients of FFT (*left bottom*), discrete Fourier spectra (*right bottom*) of the given time series, i. e., an impulse at  $t = 2.0$ .

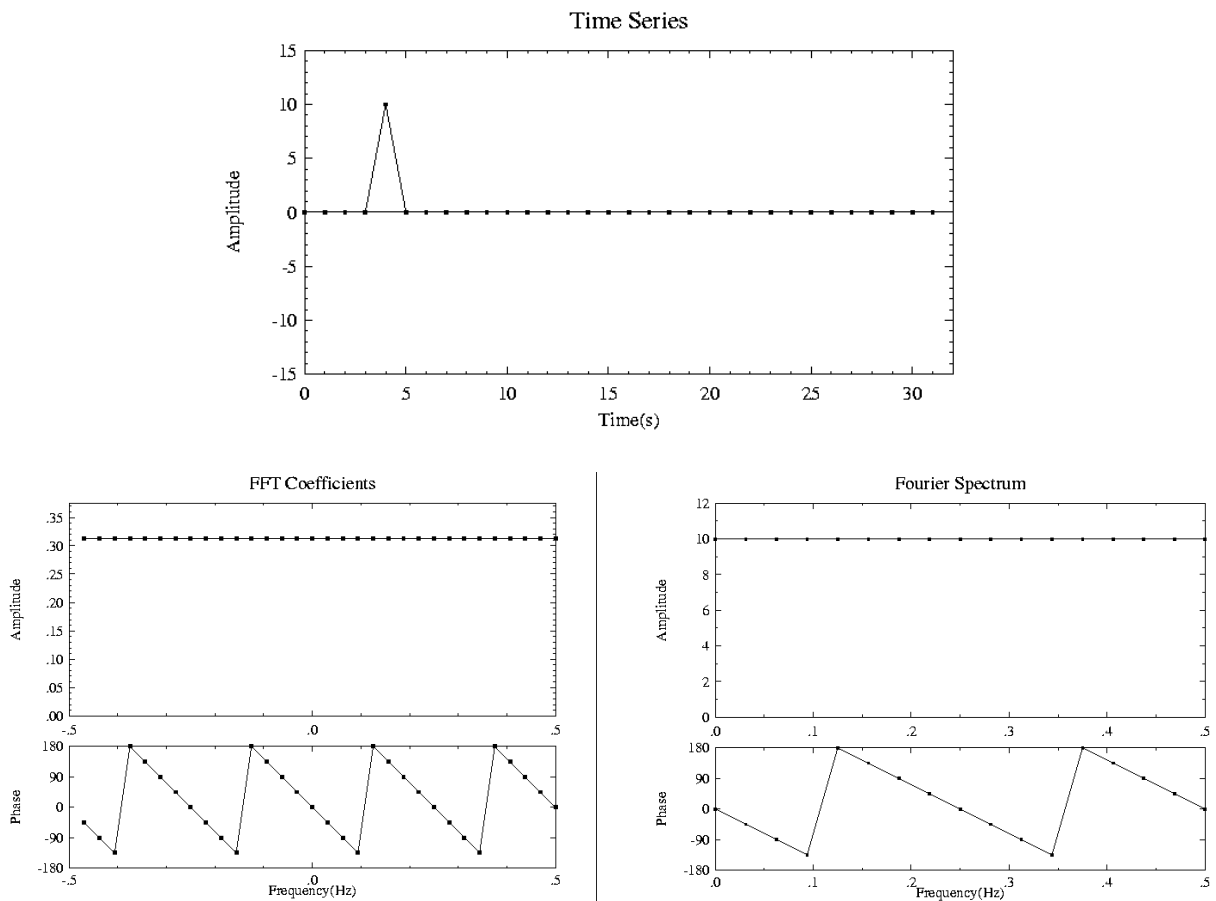


Fig. 21.2 Time series (*top*), Coefficients of FFT (*left bottom*), discrete Fourier spectra (*right bottom*) of the given time series, i. e., an impulse at  $t = 4.0$ .

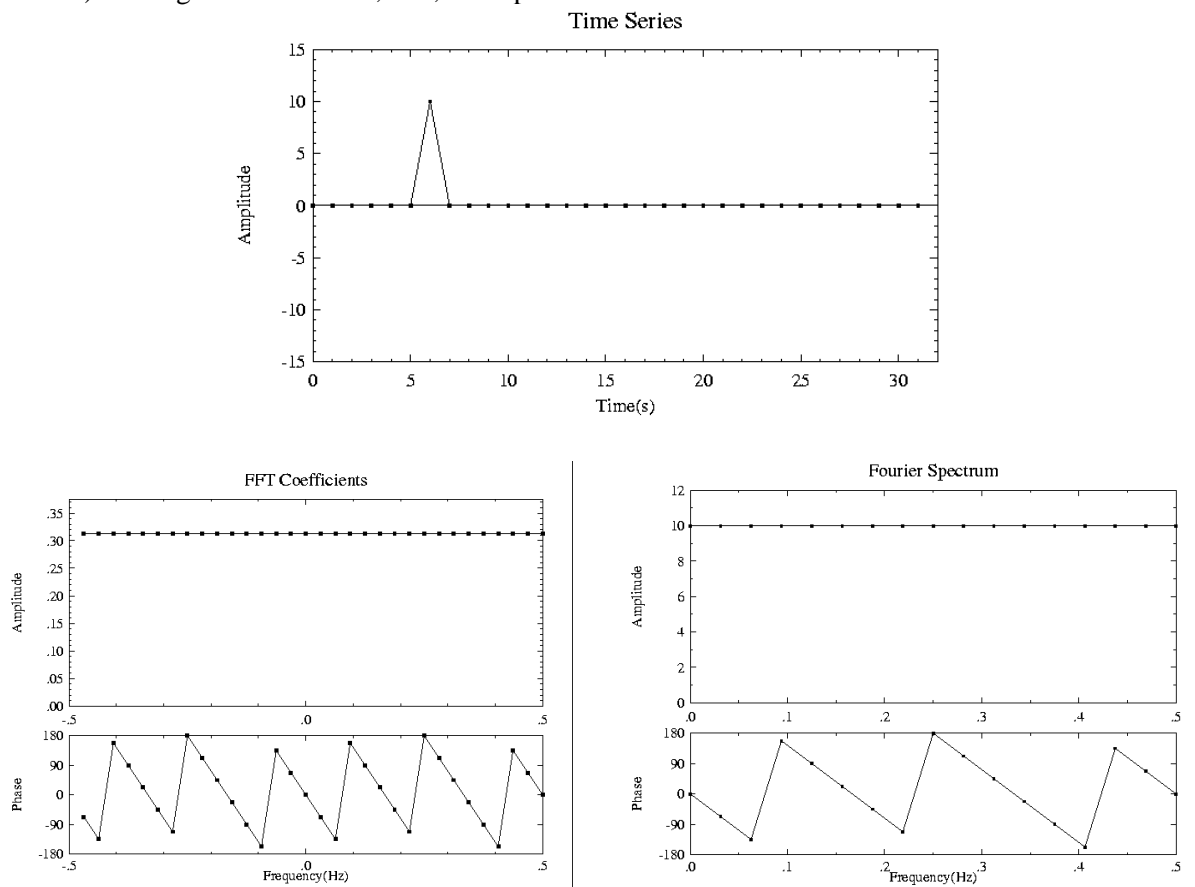


Fig. 21.3 Time series (*top*), Coefficients of FFT (*left bottom*), discrete Fourier spectra (*right bottom*) of the given time series, i. e., an impulse at  $t = 6.0$ .

### 2.5.5. Aliasing

#### Exercise: Simulate aliasing effect.

Calculate the cosine wave having a period of 0.8 s (1.25 Hz) with  $\Delta t = 1.0$  s,  $N = 32$ , amplitude = 10.0, damping = 0.0, and phase = 0.0.

The frequency of 1.25 Hz is higher than the Nyquist frequency 0.5 Hz. The time series panel below does not appear to be a 1.25-Hz wave. A false peak due to aliasing is observed in the spectrum panels.

Aliasing occurs due to the  $f = \pm(f_0 \pm 2f_{Nyquist})$  ambiguity of the frequency. In this example, the peak at 1.25 Hz is folded into  $1.25 - 2*0.5$  Hz = 0.25 Hz.

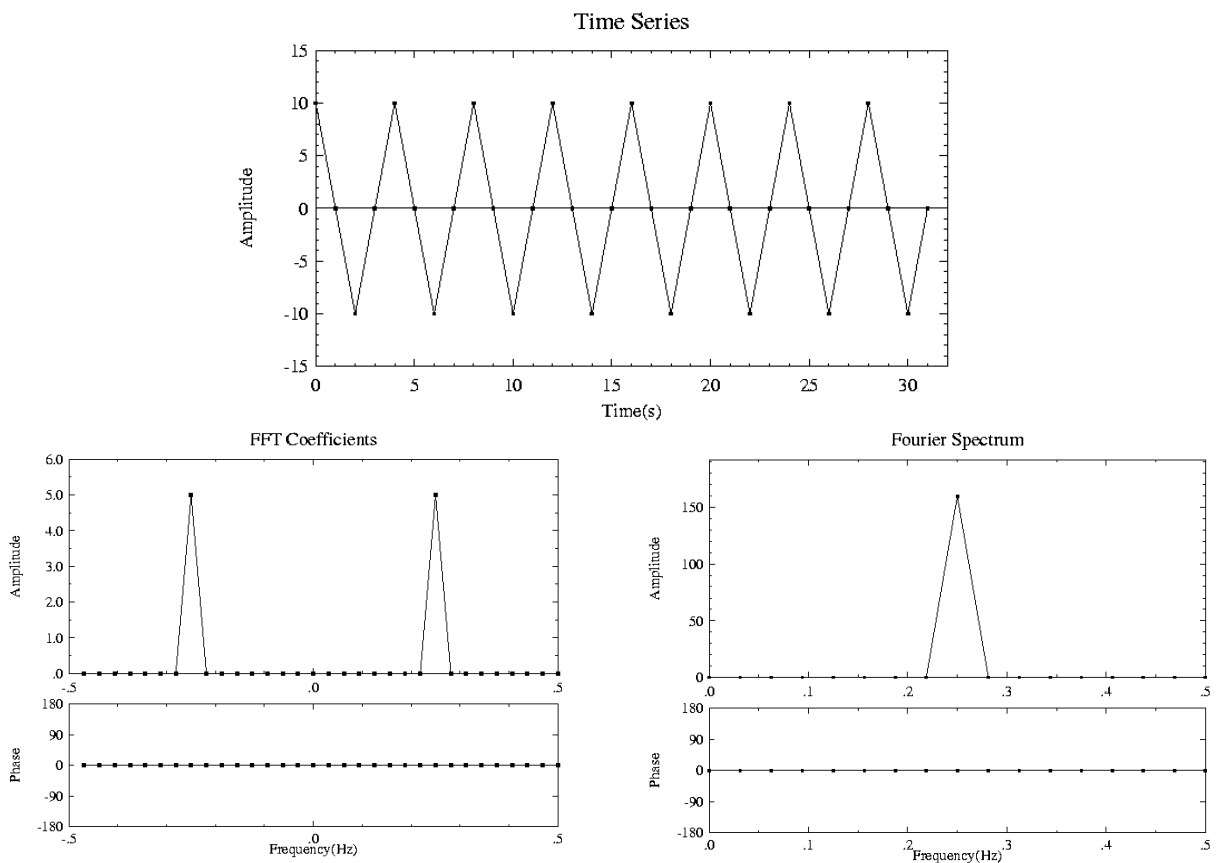


Fig. 22 Time series (*top*), coefficients of FFT (*left bottom*), discrete Fourier spectra (*right bottom*) of the given time series, i. e., a cosine function at the frequency higher than the Nyquist frequency.



### 2.5.6. Assumption of Periodicity

#### Exercise:

Calculate a cosine wave period = 7.0 s (0.14 Hz)  $\Delta t = 1.0$  s,  $N = 32$ , amplitude = 10.0, damping = 0.0, phase = 0.0, and consider why a line spectrum at  $f = 0.14$  Hz could not be obtained.

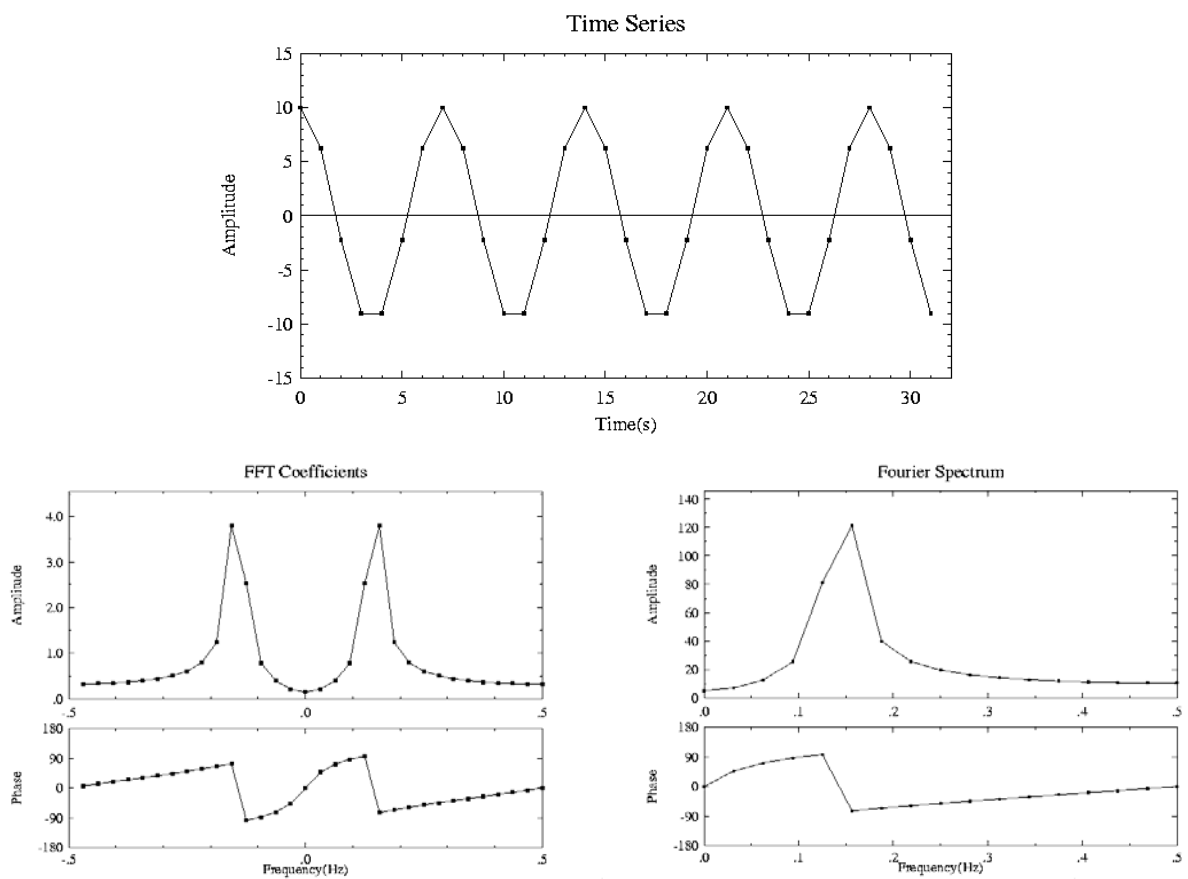


Fig. 23 Time series (*top*), coefficients of FFT (*left bottom*), discrete Fourier spectra (*right bottom*) of the given time series, i. e., a cosine function at the frequency mentioned above.

### 3. Filtering Techniques

Recorded signals are often contaminated by AC noise or high-frequency ground noise from nearby stations. Therefore, various filtering techniques are essential for digital data processing.

#### 3.1. Weighted Moving Average

The simplest method to suppress high-frequency components included in a given time series  $x_m = x(t_m)$  may be moving averages such as

$$y_m = \frac{x_{m-1} + x_m + x_{m+1}}{3}.$$

This equation shows a three-point moving average. The output  $y_m$  is defined by the one-step previous term of the input  $x_{m-1}$ , the present term  $x_m$ , and the future term  $x_{m+1}$ .

The five-point moving average is given by

$$y_m = \frac{x_{m-2} + x_{m-1} + x_m + x_{m+1} + x_{m+2}}{5}.$$

The performance of the moving average can be estimated by applying it to an impulse of the unit amplitude located at  $t = 0.0$ , because the moving average belongs to linear systems. Fig. 24 shows the Fourier spectra of the output corresponding to the impulse input for the two moving averages mentioned above. As shown here, the moving average can certainly be used to eliminate high-frequency components. However, it is difficult to control the performance and value of parameters such as the cut-off frequency, the slope of the cut-off, etc.

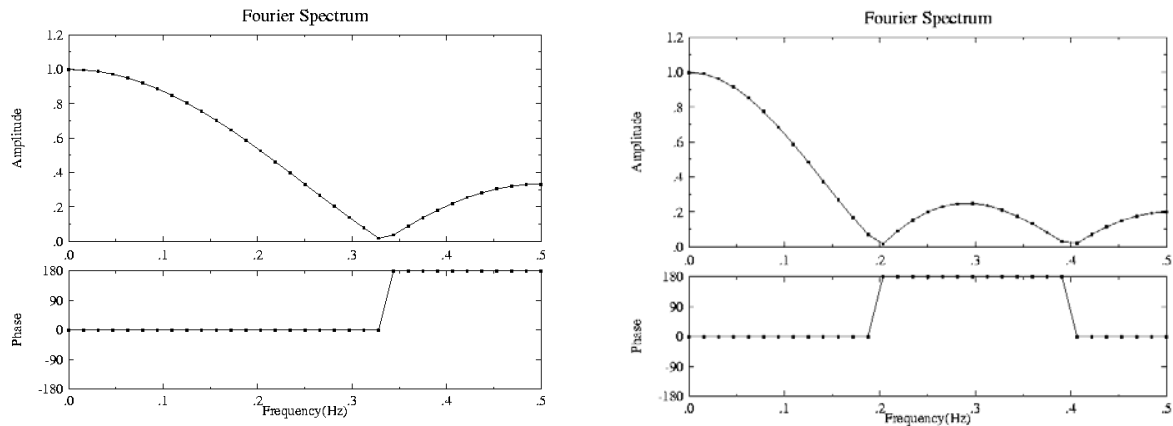


Fig. 24 Fourier spectra of the output from the three-point moving average corresponding to the impulse input (*left*) and that from the five-point moving average (*right*). These examples are calculated with  $\Delta t = 1.0$  s and  $N = 64$ .

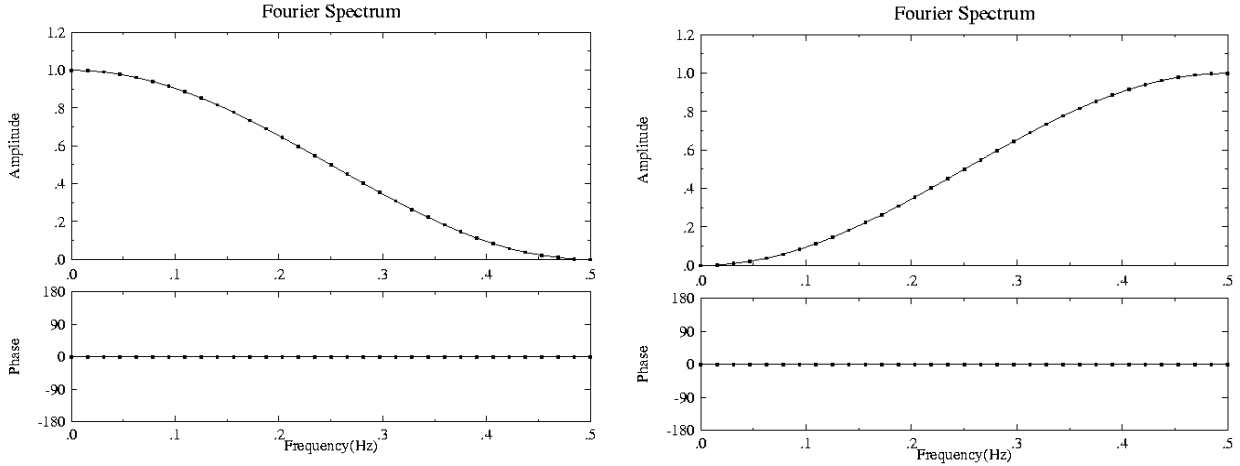


Fig. 25 Fourier spectra of the output from the weighted moving average corresponding to the impulse input. *Left*: for the weight coefficients (0.25, 0.5, and 0.25). *Right*: for the weight coefficients (-0.25, 0.5, -0.25). These examples are calculated with  $\Delta t = 1.0$  s and  $N = 64$

The weighted moving average is a similar procedure but with different weight coefficients. This gives a better performance than that given by the simple moving average.

For example, the three-point average

$$y_m = 0.25x_{m-1} + 0.5x_m + 0.25x_{m+1},$$

can eliminate high-frequency components, as shown in Fig. 25 (*left*). Note that this maintains the phase lag at zero for all frequencies.

The other example,

$$y_m = -0.25x_{m-1} + 0.5x_m - 0.25x_{m+1},$$

can eliminate low-frequency components as shown in Fig. 25 (*right*). Note that the phase lag is maintained at zero for all frequencies.

The differentiation of a continuous function  $x(t)$ ,

$$y(t) = \frac{d}{dt} x(t),$$

can be approximated by the finite difference,

$$y_m = \frac{x_m - x_{m-1}}{\Delta t} = \left(-\frac{1}{\Delta t}\right)x_{m-1} + \left(\frac{1}{\Delta t}\right)x_m.$$

This also belongs to the weighted moving average.

These examples show that the weighted moving average can have a good performance. In other words, we can arrange its characteristics by selecting the weight coefficients. By using the idea of the impulse response, we can check the characteristics of its performance. However, we must design the weighted moving average by selecting the weight coefficients in such a way that the characteristics of the performance are obtained as desired.

## 3.2. Convolution—Filtering in the Time Domain

In order to understand the methods that are used to design weight coefficients, in this chapter, the procedure for “convolution” in the time domain is examined.

### 3.2.1. Convolution

Suppose that the Fourier transform of the time dependent functions  $f(t)$  and  $g(t)$  are  $F(\omega)$  and  $G(\omega)$ , respectively. The inverse Fourier transform of the product of  $F(\omega)$  with  $G(\omega)$  is given by

$$\begin{aligned} \frac{1}{2\pi} \int_{-\infty}^{\infty} F(\omega)G(\omega)e^{i\omega t} d\omega &= \frac{1}{2\pi} \int_{-\infty}^{\infty} F(\omega)e^{i\omega t} d\omega \left( \int_{-\infty}^{\infty} g(\tau)e^{-i\omega\tau} d\tau \right) \\ &= \int_{-\infty}^{\infty} g(\tau) d\tau \frac{1}{2\pi} \int_{-\infty}^{\infty} F(\omega)e^{i\omega(t-\tau)} d\omega \\ &= \int_{-\infty}^{\infty} g(\tau)f(t-\tau) d\tau. \end{aligned}$$

This integration is referred to as the **convolution** of two functions  $f(t)$  and  $g(t)$  in the range  $(-\infty, \infty)$ . Convolution is usually expressed by an asterisk between two functions.

$$f(t) * g(t) = \int_{-\infty}^{\infty} f(\tau)g(t-\tau) d\tau = \int_{-\infty}^{\infty} f(t-\tau)g(\tau) d\tau$$

If  $f(t)$  has a non-zero value only in the range  $t_1 < t < t_2$ ,

$$f(t) * g(t) = \int_{t_1}^{t_2} f(\tau)g(t-\tau) d\tau = \int_{t-t_2}^{t-t_1} f(t-\tau)g(\tau) d\tau$$

Mathematically, convolution in the time domain corresponds to the product of the Fourier transforms in the frequency domain. The time domain operation may have advantages when one of the two time series has a short duration. In such a case, the short time series  $f(t)$  is considered as a filter for modifying the input signal  $g(t)$ . The effect of filtering must be controlled by the spectrum of the filter  $f(t)$ . Note that the amplitude spectrum of the output is the product of the amplitude spectra of the two original signals and the phase spectrum of the output is the sum of their phase spectra.

### 3.2.2. Filtering in the Time Domain by Convolution

Let us examine the calculation procedures in a computer for the convolution of two time series. Assume that  $f(t)$  has a short duration with non-zero values only within  $[0, t_M]$ .

$$h(t) = f(t) * g(t) = \int_{t-t_M}^t f(t-\tau)g(\tau) d\tau.$$

In a discretized form,

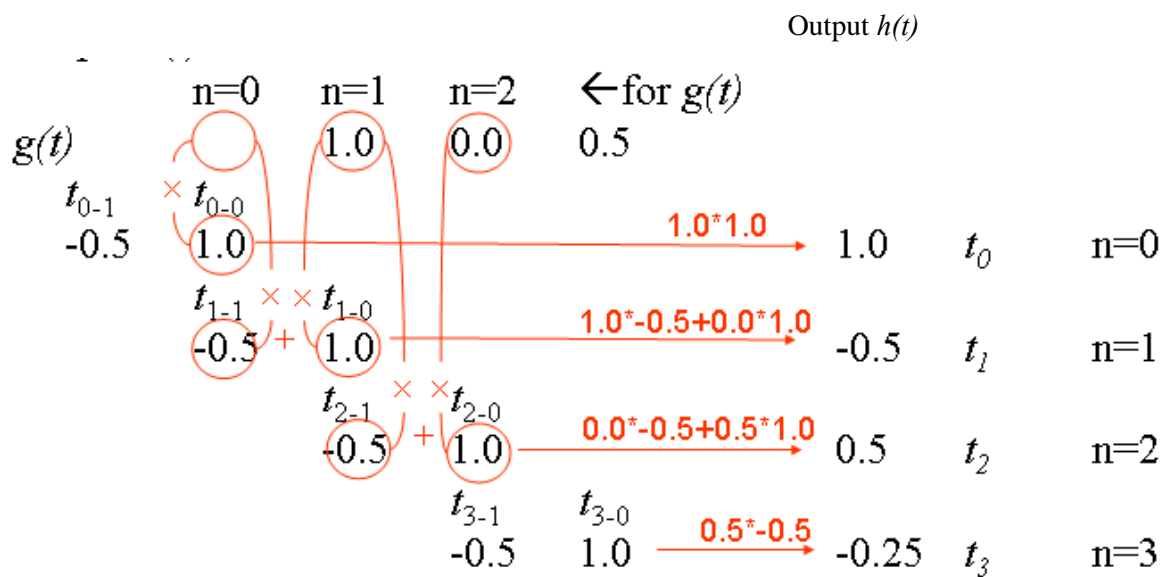
$$\begin{aligned}
 h(t_n) &= \Delta \tau \sum_{m=n-M}^n f(t_{n-m})g(t_m) \\
 &= \Delta \tau f(t_M)g(t_{n-M}) + \Delta \tau f(t_{M-1})g(t_{n-M+1}) + \dots + \Delta \tau f(t_0)g(t_n).
 \end{aligned}$$

This means that the filter time series is reversed and used as the weight coefficients for the weighted moving average, as explained in the previous chapter.

### Example: Time domain operation

Yilmaz(1994) show a graphical explanation of the convolution as follows. The convolution of a filter  $f(t) = (1.0, -0.5)$  with a signal  $g(t) = (1.0, 0.0, 0.5)$ . Assume that the sampling interval is equal to 1.0 s. Further, note that  $M = 1$ .

Reversing the filter  $f(t)$ :  $(1, -0.5)$  changes it into  $(-0.5, 1)$ .



Add the product  $f(t_{n-m})g(t_m)$  for  $m = 0$  to  $M$ . The sum gives the value  $h(t_n)$ .

Shift the moving array one sample to the right and repeat the procedure for adding products. Try to examine whether the same result is given if  $f(t)$  and  $g(t)$  exchange their roles. (After Yilmaz (1994)).

### 3.3. Feature of Filter Wavelets

When a signal is composed of only a few cycles in the time domain, it is called a “wavelet.” A wavelet is usually considered a transient signal. Undoubtedly, convolution with a wavelet can be considered equivalent to filtering. The characteristics of this filter are directly defined by the frequency spectrum of the filter wavelet. Every weighted moving average belongs to this category. The discussion in the previous chapter gave some typical examples.

Again, it must be noted that the amplitude spectrum of the output is the product of the amplitude spectra of the filter and the input signal, while the phase spectrum of the output is the sum of their phase spectra.

#### 3.3.1. Phase

The time series are composed of limited and discrete numbers of sinusoidal functions with a constant interval of frequency  $\Delta f$ , that is, the reciprocal of the duration  $T$ . The wavelet that shows a symmetry around  $t = 0$  and has a positive peak amplitude is the *zero phase* wavelet. Fig. 26 shows the decomposition of a wavelet into various sinusoids. Note that all the component sinusoids have *zero* time shifts. The phase lag is defined by  $2\pi f t_{shift}$ , where  $t_{shift}$  is the time shift. If the time shift is zero for all the frequencies, the wavelet is called “zero phase wavelet.” If the time shift is a constant for all frequency components, it is equivalent to a *linear phase shift*, an example of which is given in Fig. 27. The tangent of

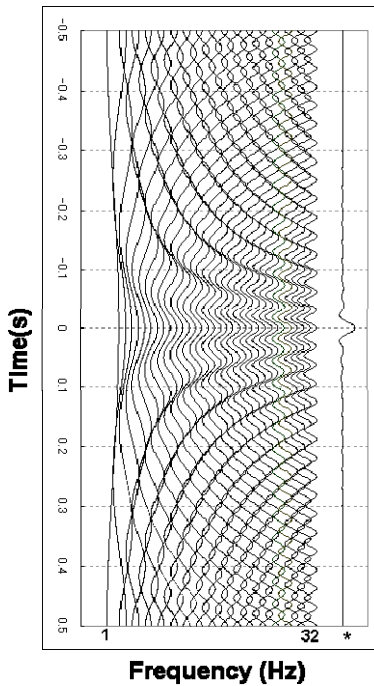


Fig. 26 Decomposition of a band limited symmetric (zero-phase) wavelet (denoted by an asterisk) into a discrete number of sinusoids with no phase lag, but with the same peak amplitude. Produced newly based on the concept of Yilmaz(1994).

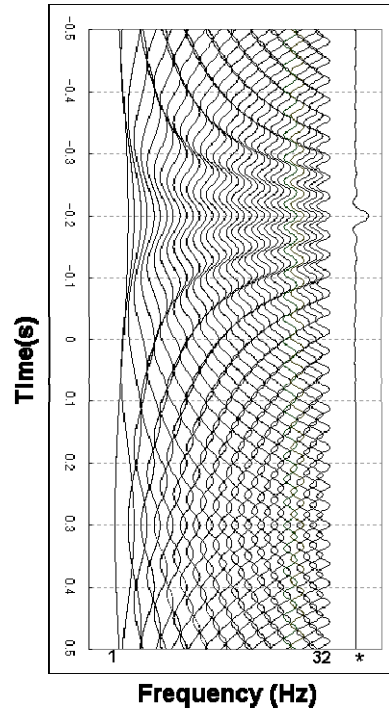


Fig.27 Constant time delay -0.2 sec given to the sinusoidal components same as those in Fig. 26 results in a wavelet of the same shape as that in Fig.26 (denoted by an asterisk), except that it is shifted in time by -0.2 sec. Produced newly based on the concept of Yilmaz(1994).

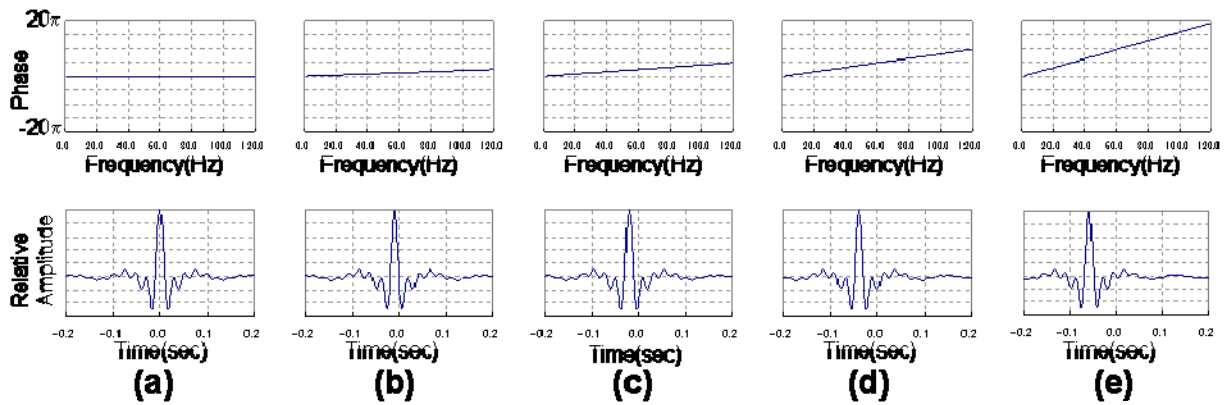


Fig.28 Wavelets in the time domain are shifted by the linear phase shift starting with a zero-phase wavelet (a). The slope of the linear phase function is related to the time shift. Produced newly based on the concept of Yilmaz(1994).

the line for the phase spectra is proportional to the time shift for the linear phase shift, as shown in Fig. 28. Note that the linear phase shift keeps the waveform constant.

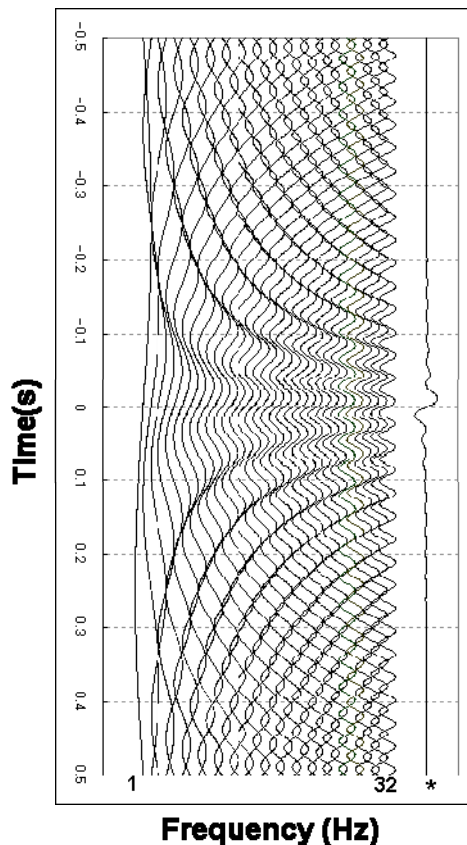


Fig. 29 Constant 90-degree phase shift given to the sinusoidal components same as those in Fig. 26 results in asymmetric wavelet but the zero crossing at  $t = 0$  (indicated by an asterisk). Produced newly based on the concept of Yilmaz(1994).

In contrast to the linear phase shift, a constant phase shift changes the waveform. The wavelet shown in Fig. 29 has the same amplitude spectrum as that in Fig. 26. The difference is their phase spectrum. Note that zero crosses are aligned in Fig. 29, whereas the peaks are aligned in Fig. 26. Fig. 30 shows the way in which the waveform is changed by a constant phase shift. Note that a constant 180-degree phase shift changes the sign of the wavelet. Note the relation of the panels (a) and (c) and panels (b) and (d). The constant phase shift of 180 degree implies a reversal of sign.

The linear and constant phase shifts are two basic examples of phase change. The combined operation is defined as  $a + b \cdot \text{frequency}$ , where  $a$  is the constant phase shift and  $b$  is the tangent of the linear phase shift, gives a time shift with a waveform change. The result is a combination of both effects, as shown in Fig. 31.

Note that the shape of the wavelet can be changed by modifying the phase spectrum even while keeping the amplitude spectrum constant. Several examples for this combination are shown

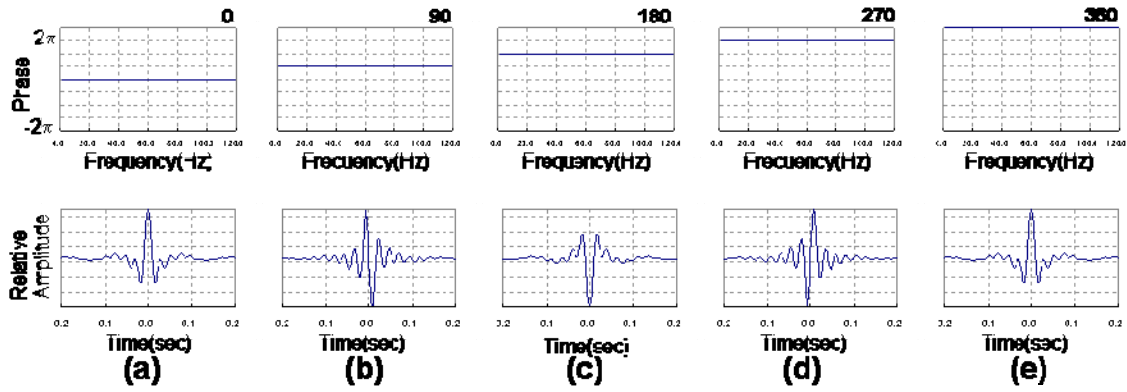


Fig. 30 Series of waveform change caused by a constant phase shift starting with the zero-phase wavelet (a). A 90-degree phase shift converts the zero-phase wavelet to an antisymmetric wavelet (b), while a 180-degree phase shift reverses its polarity (c). A 270-degree phase shift reverses the polarity, while making the wavelet antisymmetric (d). Finally, a 360-degree phase shift does not influence the wavelet (e). Produced newly based on the concept of Yilmaz(1994).

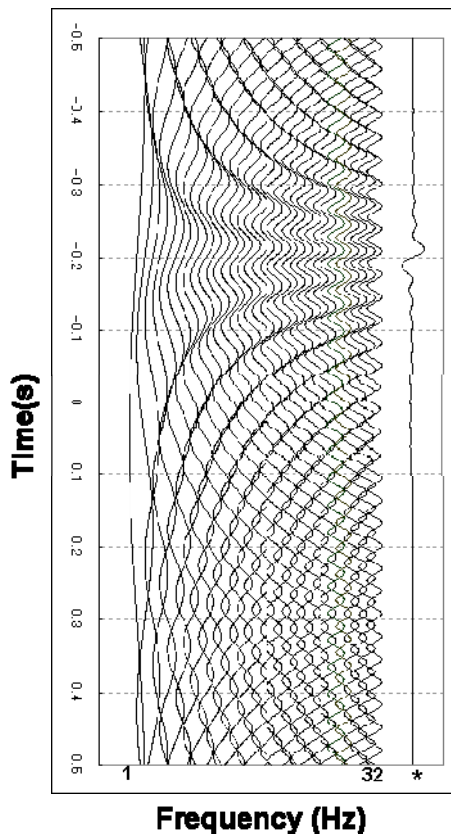


Fig. 31 Time-shifted antisymmetric wavelet (denoted by an asterisk) caused by a linear phase shift combined with a constant phase shift for the component sinusoids same as those in Fig.26. Produced newly based on the concept of Yilmaz(1994).

in Figs. 32 (a), (b), and (c). An arbitrary change in the phase spectra, however, can break the wavelet.

The tangent of the phase shift is called “delay.”

$$delay = -\frac{d\phi}{d\omega}$$

The linear phase shift is an example of constant delay for all frequencies. In general, “delay” can be dependent on the frequency.



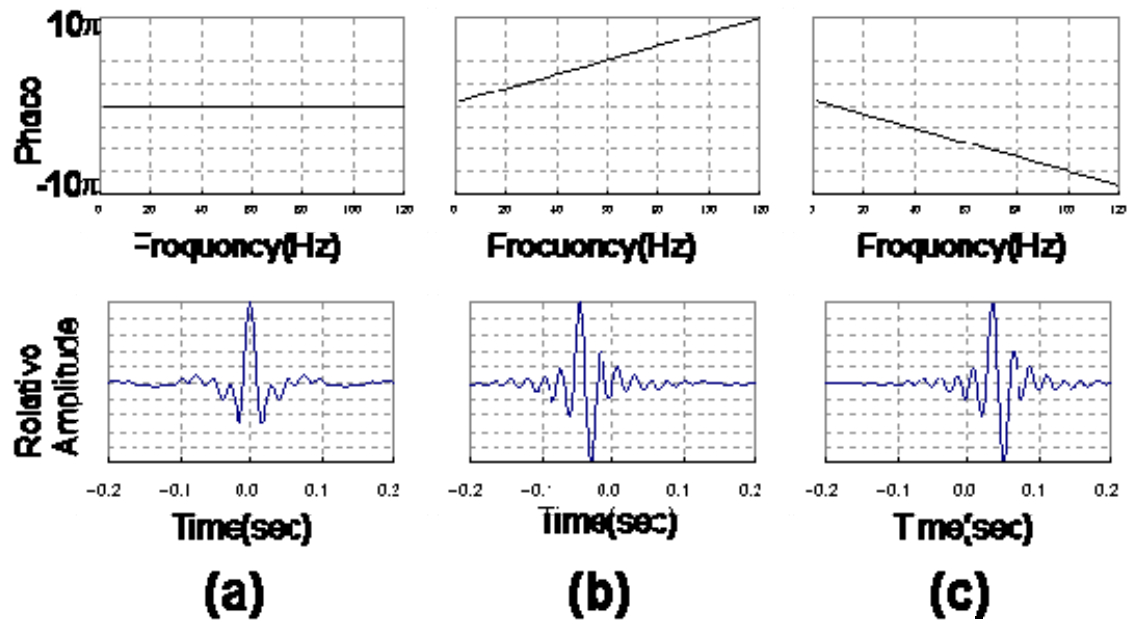


Fig. 32 a non-zero-phase spectrum of any form in (b) and (c) modifies the shape of a zero-phase wavelet (a). Produced newly based on the concept of Yilmaz(1994).

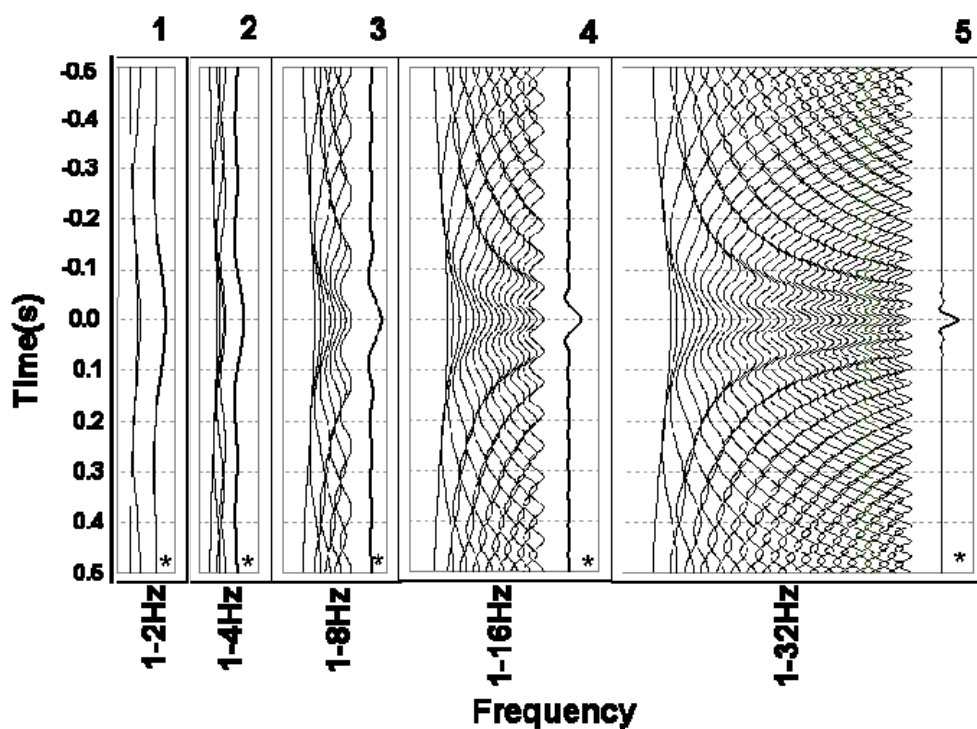


Fig. 33 The summation of zero-phase sinusoids with an identical peak amplitude shows that the increasing frequency bandwidth results in the synthesized zero-phase wavelet increasingly compressed. Produced newly based on the concept of Yilmaz(1994).

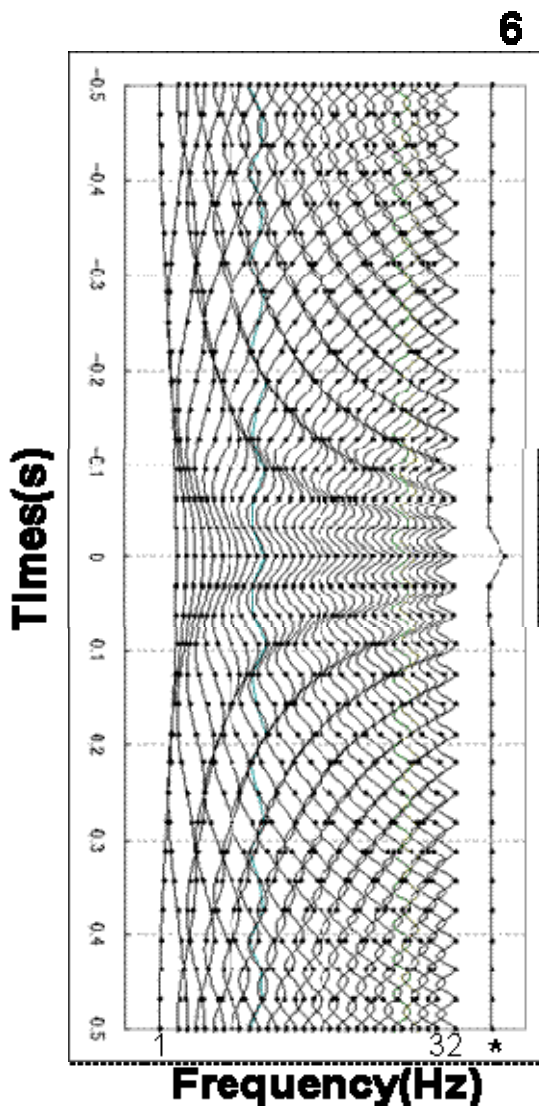
### 3.3.2. Frequency Components

In the previous chapter, wavelets with varying phase spectra and fixed amplitude spectra are observed. By changing the amplitude spectrum or selecting the frequency contents, the wavelet changes its shape even when its phase spectrum is maintained constant. Here, zero-phase wavelets are used for simplicity.

Fig. 33 shows a clear example of the changes in zero-phase wavelets by the selection of frequency contents. As more frequency components are summed, the synthesized zero-phase wavelet is increasingly compressed. If they are summed till the Nyquist frequency, a spike is formed (Fig. 34).

The broader the bandwidth, the more compressed the wavelet; in other words, a shorter wavelet is obtained. This property also follows from the fundamental concept that the effective time span of a time series is inversely proportional to its effective spectral bandwidth (Fig. 35).

The shape of the frequency spectrum also influences the wavelet shape. Fig. 36 shows a typical case. A short wavelet requires a tapered amplitude spectrum, although the width of the passband for all cases is identical.



Filtering in the frequency domain can be performed by the inverse Fourier transform of the product of the Fourier transforms of the filter and input time series. This is equivalent to the filtering in the time domain that is performed by the weighted moving average of the input time series, the weight coefficients of which are the reversed filter time series. This, in general, can be written as

$$y_i = \dots + a_2 x_{i-2} + a_1 x_{i-1} + a_0 x_i + a_{-1} x_{i+1} + a_{-2} x_{i+2} + \dots,$$

where  $(\dots, a_{-2}, a_{-1}, a_0, a_1, a_2, \dots)$  is the filter time series.

If we consider causality, i. e., the idea that the results cannot proceed to the cause, we cannot use the terms of the future  $x_{i+1}, x_{i+2}, \dots$  to obtain the present output  $y_i$ . Thus,

$$y_i = \dots + a_2 x_{i-2} + a_1 x_{i-1} + a_0 x_i.$$

Fig. 34 The output wavelet becomes a spike when the summation includes sinusoids at all frequencies up to the Nyquist frequency. Small dots denote the sampling points at 64Hz. Produced newly based on the concept of Yilmaz(1994).

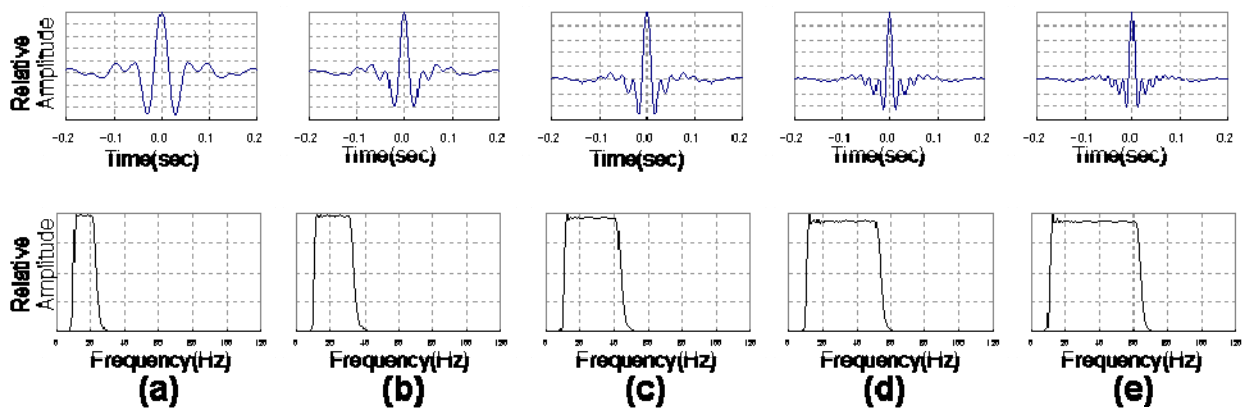


Fig. 35 Increasing bandwidth in the frequency domain (*bottom panels*) corresponds to more compressed wavelet in the time domain(*top panels*). Produced newly based on the concept of Yilmaz(1994).

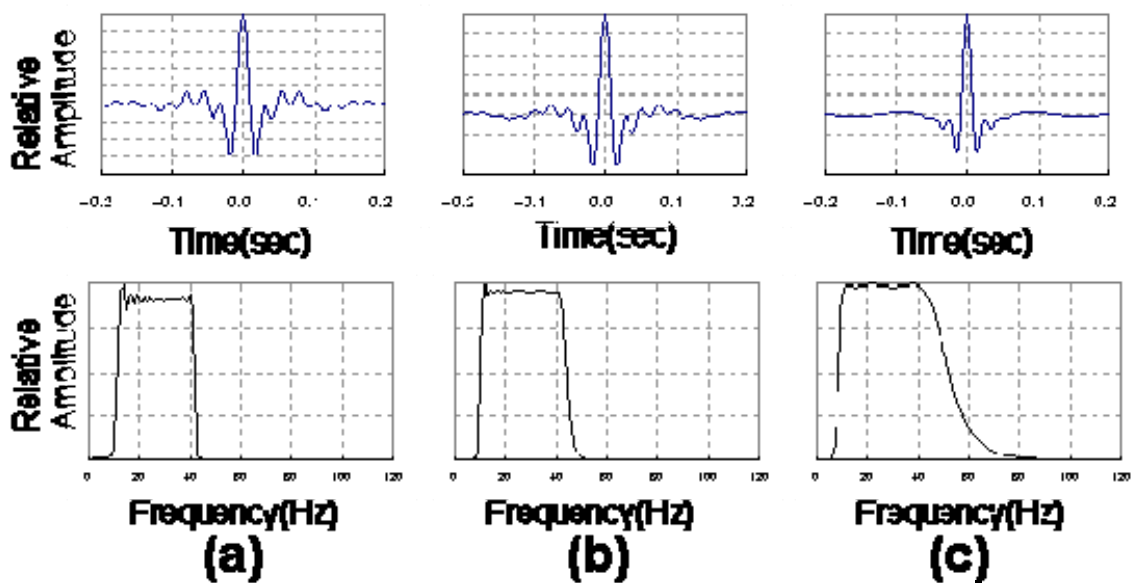


Fig. 36 More gentle slope in the frequency domain(*bottom panels*) corresponds to smoother wavelet in the time domain (*top panels*). (a) The steep slopes of the passband cause ripples in the wavelet and the actual amplitude spectrum. (b) A moderate and (c) gentle slope help eliminate the ripples. Produced newly based on the concept of Yilmaz(1994).

### 3.3.3. Causality or Non-Causal Filtering

A phenomenon that is the result of another phenomenon (the cause) never occurs before the cause itself does. This is called a “causal relation” or “causality” and it is strictly maintained in the real world. However, in a computer, this relation can be broken. Such a breakage often affects the seismological analyses. The following shows us examples.

#### Exercise: Filtering in the time domain and in the frequency domain: an example

The topics in this chapter can be learned much better by practicing with the distributed software. Here, the following programs are prepared for practice.

*FFILT.EXE* creates a set of coefficients of Fourier expansion from given bandpass characteristics.

*FPRDCT.EXE* calculates the product of two given sets of the coefficients of Fourier expansion, i. e., filtering in the frequency domain.

*FWVLET.EXE* creates a filter wavelet from the given time series that may be obtained by the inverse Fourier transform of a given set of the coefficients of Fourier expansion.

*FCONV.EXE* calculates convolution, i. e., filtering in the time domain for a given filter wavelet and input signal.

The programs *TESTSIG.EXE*, *PTIME.EXE*, *FFT.EXE*, *PCFFT.EXE*, *PSPEC.EXE*, and *IFFT.EXE* are also used.

- (0) Prepare the test input signal *UTI*, that is a unit impulse located at  $t = 8.0$  s of the time series with  $N = 64$ ,  $\Delta t = 1.0$ , and its Fourier transform *UF1* by using *TESTSIG.EXE* and *FFT.EXE* (Fig. 36a).

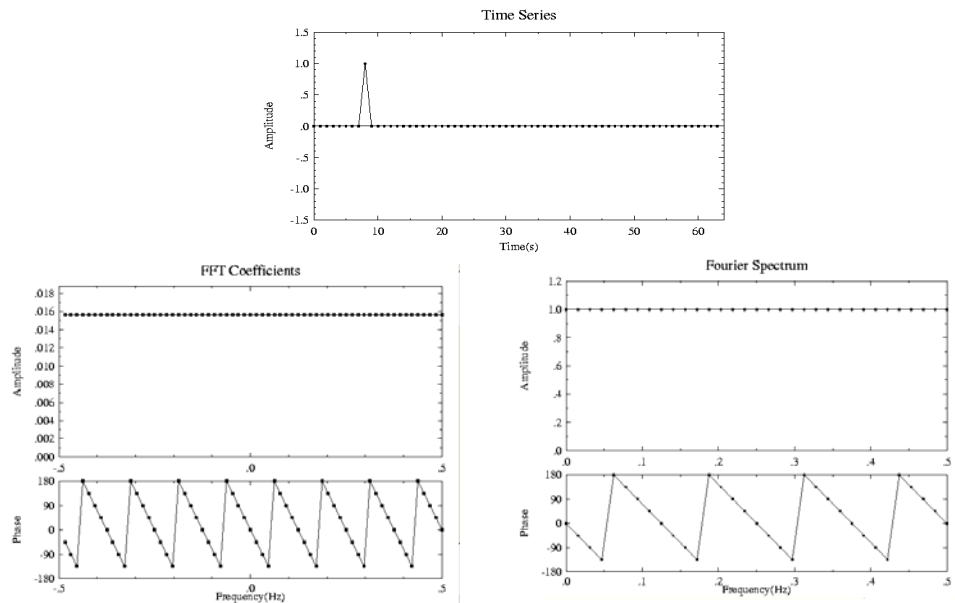


Fig. 36a Test input signal for the exercise *UTI* and its spectra. Time dependence (*upper panel*), the coefficients of Fourier expansion (*lower left panel*), and Fourier spectra (*lower right panel*). The impulse is located at  $t = 8.0$  s of the time series with  $N = 64$ ,  $\Delta t = 1.0$ . The linear phase shift due to the shifted location of the impulse is shown.

**Filtering in the frequency domain:**

(1) Design a band pass filter in the frequency domain by using *FFILT*. The number of data for the corresponding time series  $N$ , its sampling interval  $\Delta t$ , and four frequencies  $f_1, f_2, f_3$ , and  $f_4$  must be given when we run *FFILT*. Let  $N = 64$ ,  $\Delta t = 1.0$ ,  $f_1 = 0.1$  Hz,  $f_2 = 0.2$  Hz,  $f_3 = 0.3$  Hz, and  $f_4 = 0.4$  Hz and the output file name be *UF2*. Draw the coefficients of the Fourier expansion by using *PCFFT* and Fourier spectra by using *PSPEC* (Fig. 37.1). The figure is stored in the PostScript file *G.PS*. The following steps are schematically shown in Fig. 37.2.

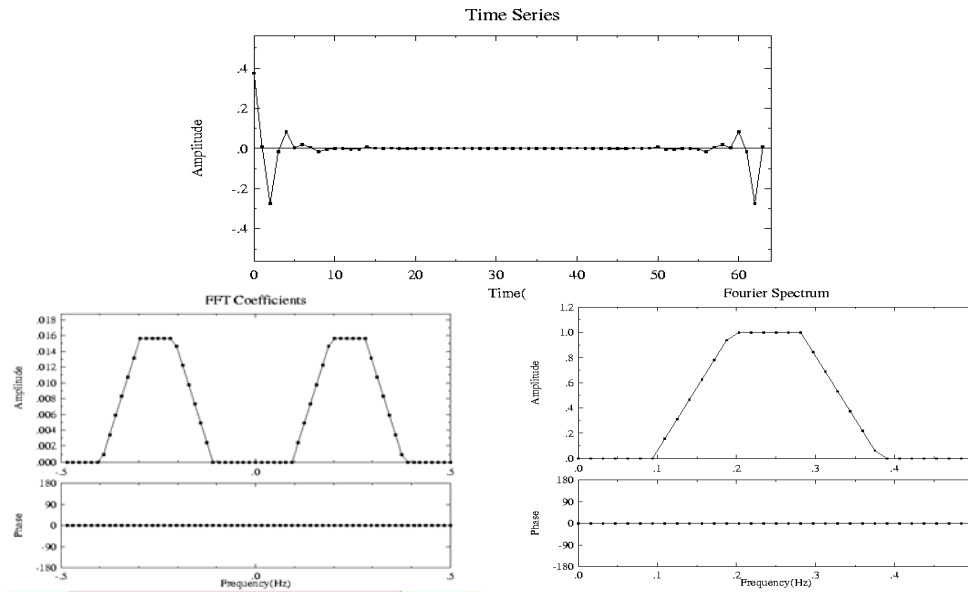


Fig. 37.1 Filter given in the frequency domain *UF2* (lower panels) and the corresponding time series *UT1* (upper panel). Since it is assumed that phase lag is zero at all frequencies, the time series is symmetrical till the point  $t = 0.0$  s. Remember the implicitly assumed periodicity. The next step of  $t = 63.0$  s is  $t = 0.0$  s.

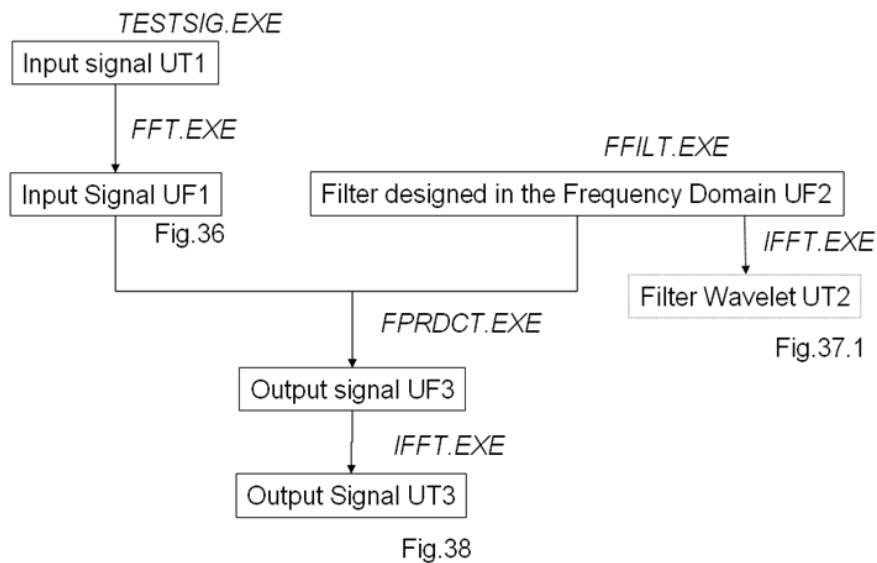


Fig. 37.2 Procedure for designing a filter in the frequency domain.

- (2) Apply the filter in the frequency domain by using *FPRDCT*. Use the stored data in *UF2* for the filter in the frequency domain. Give the output the file name *UF3*.
- (3) Apply an inverse Fourier transform by using *IFFT* and store the results in the output file *UT3*. Draw this time series and its Fourier components by using *PTIME*, *PCFFT*, and *PSPEC* (Fig. 38).
- (4) Obtain the time series corresponding to the band pass filter in the frequency domain designed by the procedure in (1) by using *IFFT*. Give the input file the name *UF2* and the output file the name *UT2*. Draw this time series by using *PTIME* (Fig. 37.1). The figure is stored in the PostScript file *G.PS*.

Remember that the input impulse is located at  $t = 8.0$  and notice that there are signals in Fig. 38 before  $t = 8.0$ . This means that the causal relation is broken because the result (output) is occurring before the cause (input) does. As shown in Fig. 37, the phase of the filter wavelet is assumed to be zero for all frequencies. Since there has been breakage of causal relation, this assumption may not be a valid one. It may be necessary to arrange the phase of the filter wavelet in order to maintain the causal relation. However, this is difficult to achieve this in the design of the filter wavelet in the frequency domain.

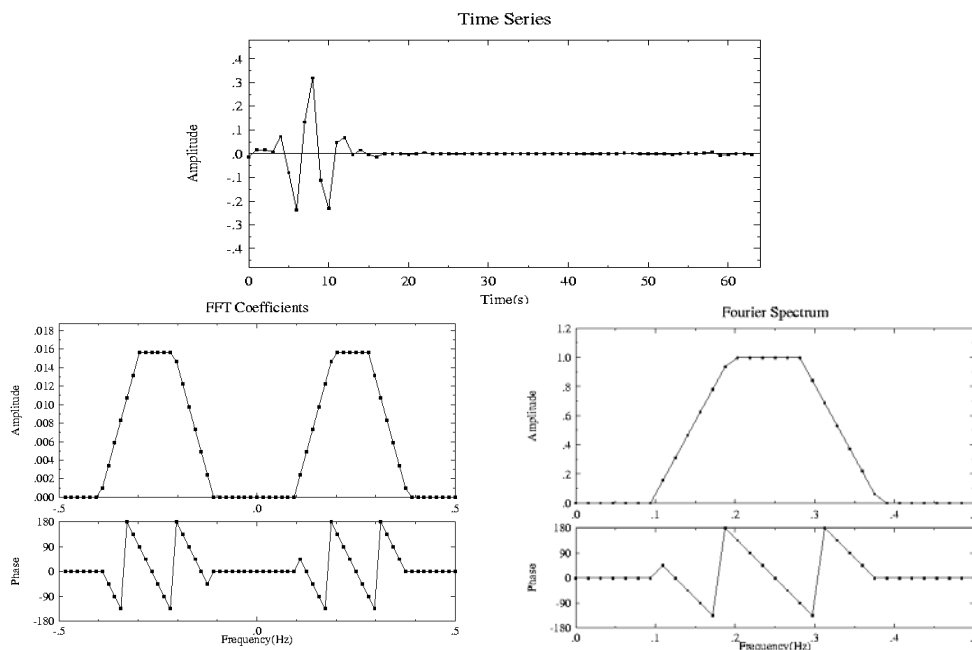


Fig. 38 The spectra of filtered signal *UF3*, which is obtained by the product of the Fourier transform of the filter shown in Fig. 37 and that of the input time series shown in Fig. 36 (*lower panels*). The upper panel shows the time series obtained by their inverse Fourier transform, *UT3*.

**Filtering in the time domain:**

(5) Extract the filter wavelet from the time series stored in the file *UT2* by using *FWVLET*. Give the output the file name *FWV1*. We have to select either a causal filter or a zero-phase filter. Here, we select a zero-phase filter with 13 coefficients.

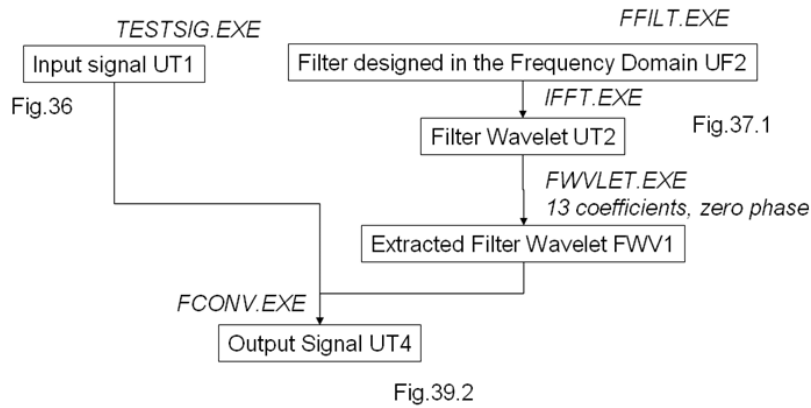


Fig. 39.1 Procedure for designing a filter in the time domain

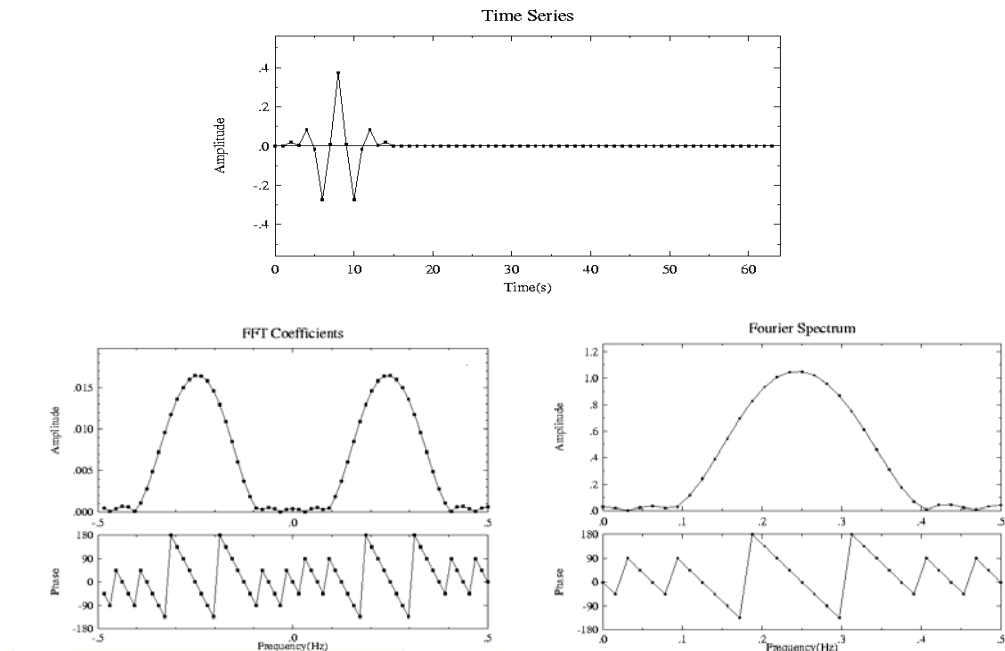


Fig. 39.2 Output time series from the filtering in the time domain with the truncated filter wavelet designed for zero phase filtering, *UT4* (upper panel), and its spectra, *UF4* (lower panels). Note the stability of the output time series and the negligible phase change in the pass band in comparison with Fig. 38. The change in the amplitude spectral shape is the effect of the truncation of the wavelet.

(6) Apply the filter in the time domain obtained in (5) by using *FCONV*. The input filter file name is *FWV*, the file name of the input time series is *UT1*, and the output file name is *UT4*. Draw the time series stored in *UT4* by using *PTIME* and compare it with the figures for *UT3* obtained in (3). Due to the truncation of the filter time series, the time series stored in *UT4* is slightly different from that in *UT3*. Check the performance of the filtering by using *FFT*, *PCFFT*, and *PSPEC* with *UT4* (Fig. 39.2).

(7) Extract the filter wavelet from the time series stored in the file *UT2* by using *FWVLET*. Give the output file the name *FWV2*. We must select either a causal filter or a zero-phase filter. In this case, we select a causal filter with 6 coefficients.

(8) Apply the filter in the time domain obtained in (5) by using *FCONV*. The input filter file name is *FWV*, the file name of the input time series is *UT1*, and the output file name is *UT5*. Draw the time series stored in *UT4* by using *PTIME* and compare it with the figure for *UT3* obtained in (3). The output in this case is clearly different from the results obtained in the time domain in (6) and from those in (3). Note that the causality is satisfied in the time domain. Check the performance of the filtering by using *FFT*, *PCFFT*, and *PSPEC* with *UT5* (Fig.40).

(9) Repeat the procedures explained above after changing the number of weight coefficients for *FWV* and compare them.

This example shows the problems with designing a filter wavelet with desirable characteristics both in the time frequency domains. In general, however, the filtering in the frequency domain works well for zero-phase filtering.

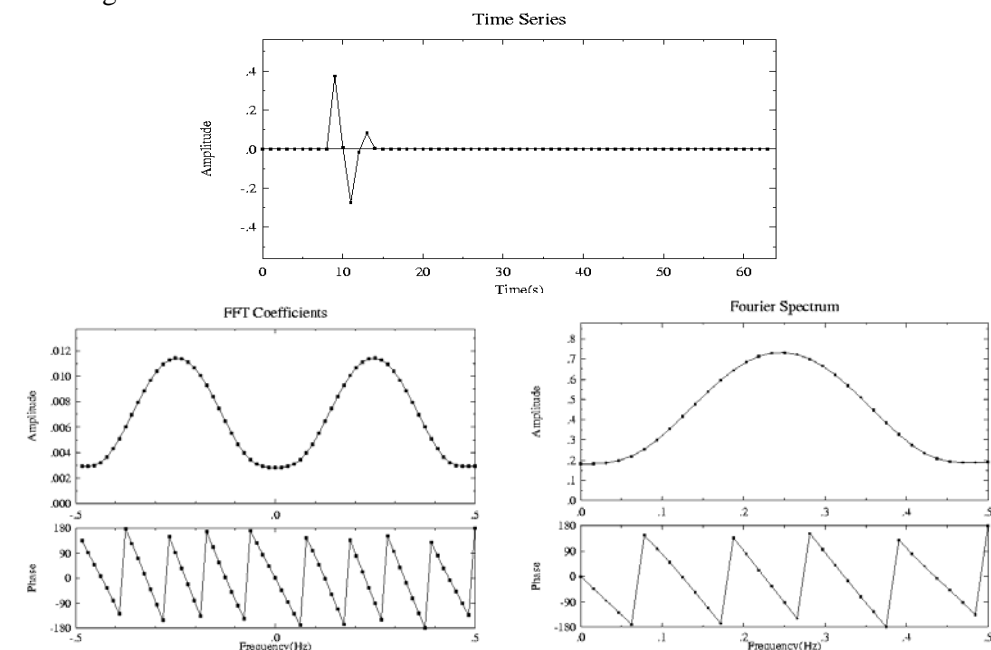


Fig. 40 Output time series from the filtering in the time domain with the truncated filter wavelet designed for zero phase filtering *UT5* (upper panel) and its spectra *UF5* (lower panels). Note that the causality with Fig.36 is maintained. The change in the amplitude and phase spectral shape is the effect of the truncation of wavelet. The truncation of a former half of the filtering wavelet results in the phase shift by filtering and reduction in the amplitude spectra even in the pass band.



### 3.4. Recursive Filter

In the previous chapter, we have checked the features of the filter wavelet that can be replaced by using the weighted moving average. The general formula of this filter is given by the following equation after taking causality into account.

$$y_i = \dots + a_2 x_{i-2} + a_1 x_{i-1} + a_0 x_i = a_0 x_i + a_1 x_{i-1} + a_2 x_{i-2} \dots$$

We have checked that the differentiation can be expressed by the weighted moving average that belongs to this category, i. e.,

$$y_m = \frac{x_m - x_{m-1}}{\Delta t} = \left(-\frac{1}{\Delta t}\right)x_{m-1} + \left(\frac{1}{\Delta t}\right)x_m.$$

Let us consider the integration given by

$$y(t) = \int_0^t x(\tau) d\tau.$$

The discretization gives

$$y_m = \Delta t \cdot \sum_{n=0}^m x_n.$$

This implies the relation,

$$y_m = y_{m-1} + \Delta t \cdot x_m.$$

Note that the term in the output that corresponds to the single step after  $y_{m-1}$  is used to construct the output for the current  $y_m$ . The filter that has had such a recursive usage of the output in the past is called a “recursive filter” (Fig. 40). The general formula for the recursive filter after taking the causality into account is given by

$$b_0 y_i = a_0 x_i + a_1 x_{i-1} + a_2 x_{i-2} \dots - (b_1 y_{i-1} + b_2 y_{i-2} + \dots).$$

#### Filter

$$y_i = a_0 x_i + a_1 x_{i-1} + a_2 x_{i-2} + \dots$$



#### Recursive Filter

$$y_i = a_0 x_i + a_1 x_{i-1} + a_2 x_{i-2} + \dots - (b_1 y_{i-1} + b_2 y_{i-2} + \dots).$$



Fig. 41 Block diagram for the filter that can be expressed by the weighted moving average (upper) and the recursive filter (lower).

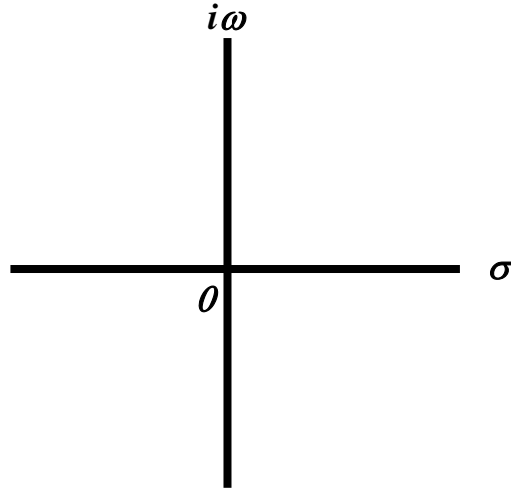


Fig.42 Complex s-plane

### 3.4.1. Laplace Transform

The Fourier transform of a continuous function has been defined previously. The meaning of Fourier transform is basically an expansion of the function on the basis of the sinusoidal function  $\exp(i\omega t)$ . The sinusoidal function with an exponential decay or amplitude  $\exp((\sigma + i\omega)t)$  can also be used as the basis for expansion. Such an integral transform is called *Laplace Transform*.

The Laplace transform and its inverse transform are defined by the following (here  $s = \sigma + i\omega$ ).

$$F(s) = \int_0^{\infty} f(t) e^{-st} dt ,$$

$$f(t) = \frac{1}{2\pi i} \int_{\gamma - i\omega}^{\gamma + i\omega} F(s) e^{st} ds .$$

The Laplace transform  $F(s)$  is defined in the complex  $s$ -domain (Fig. 42).

### 3.4.2. Filter Operation in the $s$ -domain

The instrument characteristics of seismometers, seismographs, and every electronic circuit can be described by an appropriate transfer function. The analog transfer function may be given by using the variable for the Laplace transform as follows:

$$T(s) = \frac{A_L s^L + A_{L-1} s^{L-1} + \dots + A_2 s^2 + A_1 s + A_0}{B_M s^M + B_{M-1} s^{M-1} + \dots + B_2 s^2 + B_1 s + B_0} , \quad s = \sigma + i\omega . \quad (12)$$

The stability of this analog filter is obtained simply when all the solutions of the equation,

$$B_M s^M + B_{M-1} s^{M-1} + \dots + B_2 s^2 + B_1 s + B_0 = 0 ,$$

$s_n$  has to satisfy the following condition for the stability of the system.

$$\text{Re}(s_n) = \sigma < 0 .$$

Otherwise, the circuit becomes a noise generator.

**Example: Simple Moving Coil Type Seismometer (Transfer function in the s-domain)**

The equation of motion for a pendulum's displacement in a seismometer relative to the ground  $x(t)$  induced by the ground motion  $y(t)$  is given by

$$\frac{d^2x}{dt^2} + 2h\omega_0 \frac{dx}{dt} + \omega_0^2 x = -\frac{d^2y}{dt^2}, \quad (13)$$

where  $\omega_0$  denotes the natural frequency of the pendulum and  $h$  denotes the damping factor. Applying the Fourier transform to both sides yields

$$-\omega^2 x_m + 2ih\omega_0 \omega x_m + \omega_0^2 x_m = \omega^2 y_m \quad .$$

Thus, the response in the frequency domain is given by

$$-x_m / y_m = \frac{1}{1 - 2ih(\omega_0/\omega) - (\omega_0/\omega)^2}.$$

This response belongs to a high pass filter, and therefore, the seismometer has an equivalent digital filter.

Define the transfer function in the frequency domain,

$$T(i\omega) = -x_m / y_m = \frac{1}{1 - 2ih(\omega_0/\omega) - (\omega_0/\omega)^2} = \frac{(i\omega)^2}{(i\omega)^2 + 2h\omega_0(i\omega) + \omega_0^2}. \quad (14)$$

The Laplace transform of Eq. (14) gives the transfer function in the  $s$ -domain. The substitution of  $i\omega$  with  $s$  in Eq. (15) gives the result:

$$T(s) = -\frac{X(s)}{Y(s)} = \frac{s^2}{s^2 + 2h\omega_0 s + \omega_0^2}. \quad (15)$$

The solutions obtained when the denominator = 0 are called “poles.” In contrast, the solutions of the numerator = 0 are called “zeros,” because these cause the transfer function to be equal to zero. Eq. (12) can be factorized by using these poles and zeros as follows.

$$T(s) = \frac{A_L}{B_M} \cdot \frac{(s - s_L^0) \cdots (s - s_2^0)(s - s_1^0)}{(s - s_M) \cdots (s - s_2)(s - s_1)} = G_0 \cdot \frac{(s - s_L^0) \cdots (s - s_2^0)(s - s_1^0)}{(s - s_M) \cdots (s - s_2)(s - s_1)}. \quad (16)$$

The suffix  $0$  denotes the “zero” point. As shown, “zeros” and “poles” determine the transfer function with a constant  $G_0$ . If  $s$  coincides with one of the poles,  $T(s)$  becomes infinite. If  $s$  coincides with one of the zeros,  $T(s)$  becomes zero. Actually,  $s = \sigma + i\omega$  moves only along the imaginary axis in the complex  $s$ -plane, and poles must locate at  $\sigma < 0$  in stable systems. Then,  $s$  cannot coincide exactly with any of the poles. Poles located near the imaginary axis can induce resonance. If one of the zeros lies on the imaginary axis,  $T(s)$  becomes zero sharply at the corresponding frequency. This feature is important for the design of a notch filter. The pole-zero representation of an analog transfer function provides a method for the design of circuits. Readers are recommended to study books on electronics, especially on *active filters* for more information. Several examples for simple transfer functions will be given in the following description.

Today, many seismic observation organizations release their data to the public via the Internet so that any researcher can use them. Some of these organizations provide information on instrumental characteristics using the pole-zero representation. Hence, it may be useful to show the method of reconstructing the transfer function in the frequency domain from a given value of poles and zeros.

At an angular frequency  $\omega$ , the variable of the Laplace transform  $s$  is located at  $(0, i\omega)$ .  $(s - s_m)$  in the denominator of Eq. (13) implies the distance between  $s = (0, i\omega)$  and the pole  $s_m$  taking phase into consideration as well. Namely,

$$(s - s_m) = |(s - s_m)| \exp(\arg(s - s_m)).$$

When all poles and zeros are similarly considered, the following relations are obtained:

$$\begin{aligned} T(s) &= G_0 \cdot \frac{(s - s_L^0) \cdots (s - s_2^0)(s - s_1^0)}{(s - s_M) \cdots (s - s_2)(s - s_1)} \\ &= G_0 \frac{|s - s_L^0| \cdots |s - s_2^0| |s - s_1^0|}{|s - s_M| \cdots |s - s_2| |s - s_1|} \exp\{i \arg(s - s_L^0) + \cdots + i \arg(s - s_1^0) - i \arg(s - s_M) - \cdots - i \arg(s - s_1)\} \end{aligned}$$

Suppose that  $X_L \dots X_2, X_1$  denote the absolute values of  $(s - s_m)$  and  $\Theta_L \dots \Theta_2, \Theta_1$  denote the absolute values of their phases. Similarly,  $x_M \dots x_2, x_1$  and  $\theta_M \dots \theta_2, \theta_1$  for poles. Then,

$$T(s) = G_0 \frac{X_L \cdots X_2 X_1}{x_M \cdots x_2 x_1} \exp\{i\Theta_L + \cdots + i\Theta_1 - i\theta_M - \cdots - i\theta_1\}.$$

This shows that the transfer function can be reconstructed from the given values of poles and zeros by a direct graphical measurement on the complex  $s$ -plane without any special software. This simple feature is one of the advantages of introducing Laplace transforms in the analysis of transfer functions.

### Example: Simple Moving Coil Type Seismometer (Poles and Zeros)

Eq. (15) is factorized in the following manner.

$$T(s) = -\frac{X(s)}{Y(s)} = \frac{s^2}{s^2 + 2h\omega_0 s + \omega_0^2} = \frac{(s - s_2^0)(s - s_1^0)}{(s - s_2)(s - s_1)}.$$

The solutions of the equation, achieved by equating the denominator to zero are  $s = \omega_0(-h \pm \sqrt{h^2 - 1})$ .

This gives the pole position at

$$\begin{aligned} & \left(-\omega_0 h, \omega_0 \sqrt{1 - h^2}\right), \left(-\omega_0 h, -\omega_0 \sqrt{1 - h^2}\right) && \text{for } h < 1.0, \text{ under-damped case,} \\ & (-\omega_0, 0.) && \text{doubled for } h = 1.0, \text{ critically damped case,} \\ & \left(-\omega_0(h - \sqrt{h^2 - 1}), 0.\right), \left(-\omega_0(h + \sqrt{h^2 - 1}), 0.\right) && \text{for } h > 1.0, \text{ over-damped case.} \end{aligned}$$

For all these three cases, the poles are located in the left half of the  $s$ -plane. This guarantees the stability of the system. The doubled zeros are located at  $(0,0)$ .

### 3.4.3. Z-transform

Remember the discrete Fourier transform:

$$\begin{aligned}
 X_k &= TC_k = \Delta t \sum_{m=1}^N x_m e^{-im\omega_k \Delta t}, \\
 x_m &= \frac{1}{N\Delta t} \sum_{k=1}^N X_k e^{im\omega_k \Delta t}, \\
 \text{where } \omega_k &= \frac{2\pi k}{N\Delta t}.
 \end{aligned} \tag{17}$$

$X_k$  has a certain physical meaning. Let us change Eq. (17) slightly in the following way.

$$\begin{aligned}
 \tilde{X}_k &= \sum_{m=1}^N x_m e^{-im\omega_k \Delta t}, \\
 x_m &= \frac{1}{N} \sum_{k=1}^N \tilde{X}_k e^{im\omega_k \Delta t}.
 \end{aligned}$$

This gives an abstract quantity in the transformed domain. A new variable is introduced as

$$z = e^{i\omega_k \Delta t}. \tag{18}$$

Thus,

$$\begin{aligned}
 Z(x_m) &= \tilde{X}_k = \sum_{m=1}^N x_m z^{-m}, \\
 x_m &= \frac{1}{N} \sum_{k=1}^N \tilde{X}_k z^k.
 \end{aligned} \tag{19}$$

This new integral transform for discrete systems is called *z-transform*. Eq. (18) can be extended to relate the discrete z-transform with a continuous Laplace transform with  $s = \sigma + i\omega$ .

$$z = e^{s\Delta t}. \tag{20}$$

The product with  $z$  implies a time shift of  $\Delta t$  toward the future, whereas that with  $z^{-1}$  implies one toward the

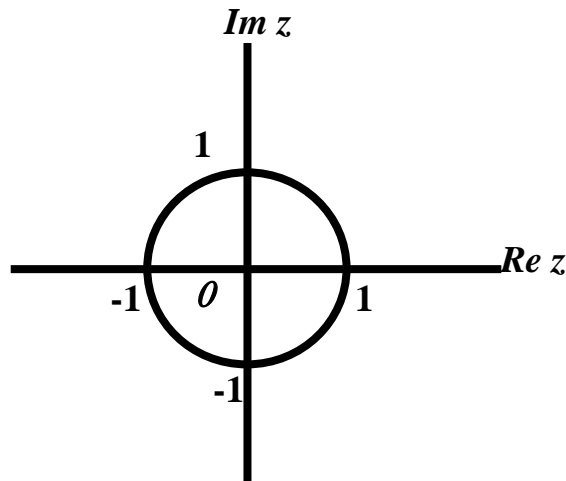


Fig. 43 Complex z-plane. The left half of the complex s-plane is mapped into a unit circle centered at the origin by  $z = \exp(st)$ . The points  $(0, i\omega/2)$ ,  $(0, -i\omega/2)$  on the s-plane are mapped to  $(-1, 0)$ . In other words, the positive and negative parts of the imaginary axis on the s-plane are mapped to the upper and lower halves of the unit circle on the z-plane, respectively. The origin of the s-plane is mapped to  $(1, 0)$  on the z-plane.

past.

### 3.4.4. Filter Operator in the Z-domain

Suppose  $x(t)$  denotes the input time series;  $X(\omega)$ , its Fourier spectrum;  $y(t)$ , the filtered output;  $Y(\omega)$ , its spectrum; and  $F(\omega)$ , the spectrum of the applied filter. Then,

$$Y(\omega) = F(\omega)X(\omega). \quad (21)$$

Suppose the filter spectra can be written, *e. g.*, in the following form in order to facilitate ease of discussion.

$$F(z(\omega)) = \frac{a_0 + a_1z^{-1} + a_2z^{-2}}{b_0 + b_1z^{-1} + b_2z^{-2}}, \quad (22)$$

Eq. (21) gives the relation

$$[b_0 + b_1z^{-1} + b_2z^{-2}]Y(\omega) = [a_0 + a_1z^{-1} + a_2z^{-2}]X(\omega).$$

The inverse Fourier transform of both sides gives

$$b_0y(t) + b_1y(t - \Delta t) + b_2y(t - 2\Delta t) = a_0x(t) + a_1x(t - \Delta t) + a_2x(t - 2\Delta t),$$

because of the relation

$$\frac{1}{2\pi} \int_{-\infty}^{\infty} z^{-n} Y(\omega) e^{i\omega t} d\omega = \frac{1}{2\pi} \int_{-\infty}^{\infty} Y(\omega) e^{i\omega(t-n\Delta t)} d\omega = y(t - n\Delta t) \quad .$$

Then, the filtered output can be calculated rapidly with a defined value of coefficients, a few preceding data of the input time series, and a few preceding data of the output. For Eq. (22),

$$\left. \begin{aligned} y_0 &= \frac{a_0}{b_0} x_0 \\ y_1 &= \frac{1}{b_0} (a_0 x_1 + a_1 x_0 - b_1 y_0) \\ y_j &= \frac{1}{b_0} (a_0 x_j + a_1 x_{j-1} + a_2 x_{j-2} - b_1 y_{j-1} - b_2 y_{j-2}), \quad j \geq 2 \end{aligned} \right\} . \quad (23)$$

This shows an example for a recursive filter operating in the time domain. The filter  $F(\omega)$  might give a phase lag. In order to compensate for the phase lag, namely, in order to apply a zero-phase filter, inverse the time axis and apply the same filter in such way that

$$\left. \begin{aligned} y_N &= \frac{a_0}{b_0} x_{N-1} \\ y_{N-2} &= \frac{1}{b_0} (a_0 x_{N-2} + a_1 x_{N-1} - b_1 y_{N-1}) \\ y_{N-j} &= \frac{1}{b_0} (a_0 x_{N-j} + a_1 x_{N-j+1} + a_2 x_{N-j+2} - b_1 y_{N-j+1} - b_2 y_{N-j+2}), \quad j \geq 3 \end{aligned} \right\} . \quad (24)$$

Of course, we can employ more coefficients  $a_i$  and  $b_m$  if necessary. The general form of Eq. (22) may be

$$F(z) = \frac{a_L z^{-L} + a_{L-1} z^{-(L-1)} + \cdots + a_2 z^{-2} + a_1 z^{-1} + a_0}{b_M z^{-M} + b_{M-1} z^{-(M-1)} + \cdots + b_2 z^{-2} + b_1 z^{-1} + b_0}, z = e^{i\omega\Delta t}. \quad (25)$$

The corresponding recursive filter may be

$$y_j = \frac{1}{b_0} \left\{ a_0 x_j + a_1 x_{j-1} + a_2 x_{j-2} + \cdots + a_{L-1} x_{j-L+1} + a_L x_{j-L} \right. \\ \left. - (b_1 y_{j-1} + b_2 y_{j-2} + \cdots + b_{M-1} y_{j-M+1} + b_M y_{j-M}) \right\}. \quad (26)$$

This filter given by Eq. (25) is sometimes called a *transfer function* by analogy with the filters of electronic circuits, which are analog filters. Since the denominator of Eq. (25) controls the feedback part of Eqs. (23) and (24),  $b_m$  must be selected carefully in order to avoid any instability in the filtering. The transform  $z = \exp(s\Delta t)$  maps the left half of the complex  $s$ -plane to the unit circle centered at the origin on the complex  $z$ -plane. Therefore, a stable and causal filtering requires that all the solutions of the equation

$$b_M z^{-M} + b_{M-1} z^{-(M-1)} + \cdots + b_2 z^{-2} + b_1 z^{-1} + b_0 = 0$$

$z_n$  must satisfy the condition

$$|z_n| < 1. \quad (27)$$

Additionally, if there are no zeros outside the unit circle on the complex  $z$ -plane, it is called minimum phase condition.

Focus on the direct coincidence of the coefficients of a filter wavelet in the time domains shown in Eq. (23) and Eq. (24) with the coefficients used in the transfer function in the  $Z$ -domain that is shown in Eq. (22). This shows that the analysis of the transfer function in the  $Z$ -domain gives the value of the coefficients for recursive filtering in the time domain.

The transfer function given in the frequency domain and that given in the  $s$ -domain are analogous functions, whereas that represented by a recursive filter is applied to a discrete time series in the computer.  $Z$ -transform behaves like an interpreter at the border between two worlds that are different each other—one a continuous world and the other a digital one.

The relation between the transfer function in the  $Z$ -domain and that in the  $s$ -domain is given approximately by the so called ***bilinear transform***:

$$s = \frac{2}{\Delta t} \cdot \frac{1 - z^{-1}}{1 + z^{-1}}, \quad z = \frac{1 + \frac{\Delta t \cdot s}{2}}{1 - \frac{\Delta t \cdot s}{2}}. \quad (28)$$

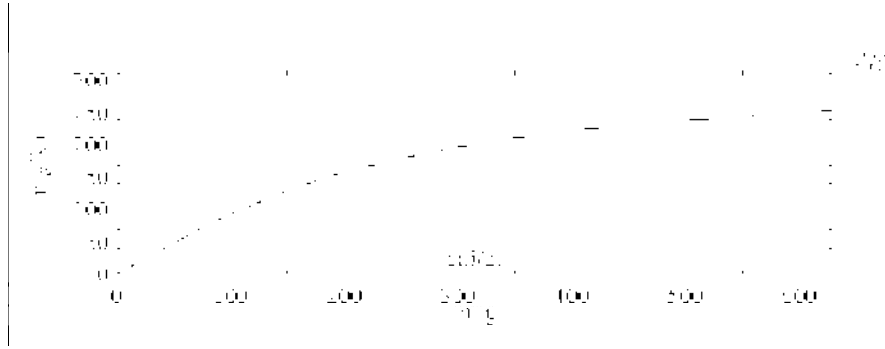


Fig. 44.1 Mapping of the continuous angular frequency  $\omega$  on to the discrete frequency  $\omega_k$ . This is an example of  $\Delta t=0.01$ ,  $N=1024$ . The Nyquist angular frequency is  $\pi/\Delta t$ .

This is an approximation of Eq. (18) that is equivalent to the following:

$$s = \frac{1}{\Delta t} \ln(z).$$

The discrete angular frequency  $\omega_k$  is also transformed. For  $z = \exp(i\omega_k \Delta t)$ , the bilinear transform gives

$$s = \frac{2}{\Delta t} \cdot \frac{1 - e^{-i\omega_k \Delta t}}{1 + e^{-i\omega_k \Delta t}} = \frac{2i}{\Delta t} \cdot \tan(\omega_k \Delta t / 2). \text{ Since } s = i\omega \text{ for } \sigma = 0,$$

$$\omega = \frac{2}{\Delta t} \cdot \tan(\omega_k \Delta t / 2) \text{ or } \omega_k = \frac{2}{\Delta t} \cdot \tan^{-1}(\omega \Delta t / 2) \quad (29)$$

This shows the distortion of the continuous angular frequency  $\omega$  by the bilinear transform to the discrete one  $\omega_k$  (Fig. 44.1). To compensate this distortion, a warped angular frequency

$$\omega'_c = \frac{2}{\Delta t} \cdot \tan(\omega_c \Delta t / 2) \quad (30)$$

is introduced (Sherbaum(1996)). The critical angular frequencies  $\omega_c$  of the continuous transfer function are first converted to the corresponding warped angular frequencies  $\omega'_c$ , and then the bilinear transform using the warped ones is applied to obtain the equivalent discrete transfer function (Fig. 44.2). Eq. (30) shows that  $\omega'_c$  tends to  $\omega_c$  for a small value of

$$\frac{\omega_c \Delta t}{2} = \frac{\pi f_c}{f_{\text{Sampling}}} = 2\pi \cdot \frac{f_c}{f_{\text{Nyquist}}}$$

This means that considering the warped frequency is not necessary for frequencies that are considerably smaller than the Nyquist frequency. The natural frequency of a seismometer is usually much smaller than the Nyquist frequency, whereas the anti-aliasing filter has a cut-off frequency that is comparable with the Nyquist one.



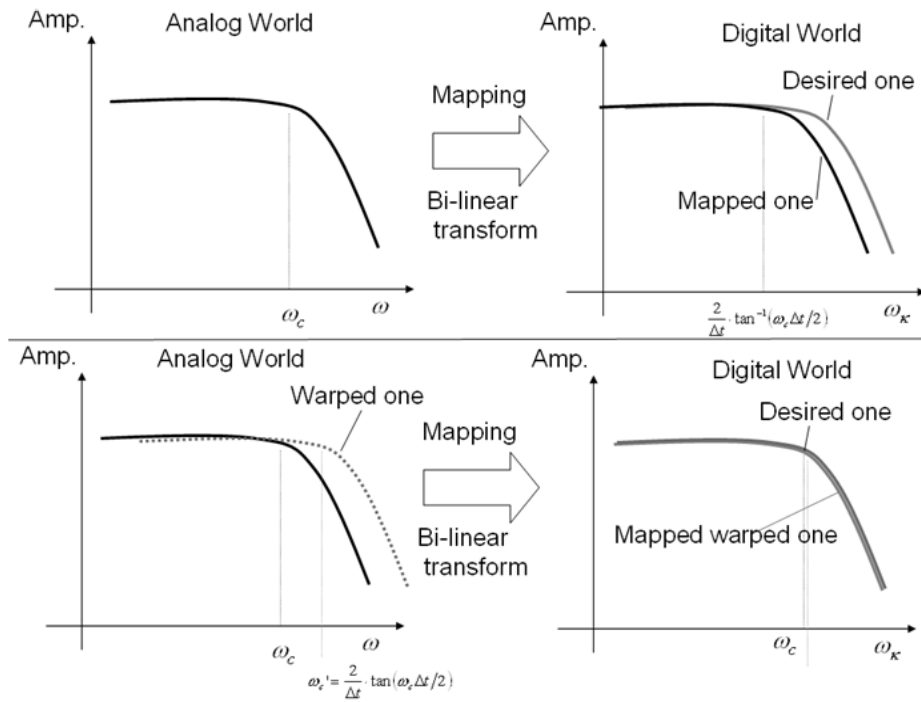


Fig. 44.2 Schematic drawing that shows how warped frequency works.

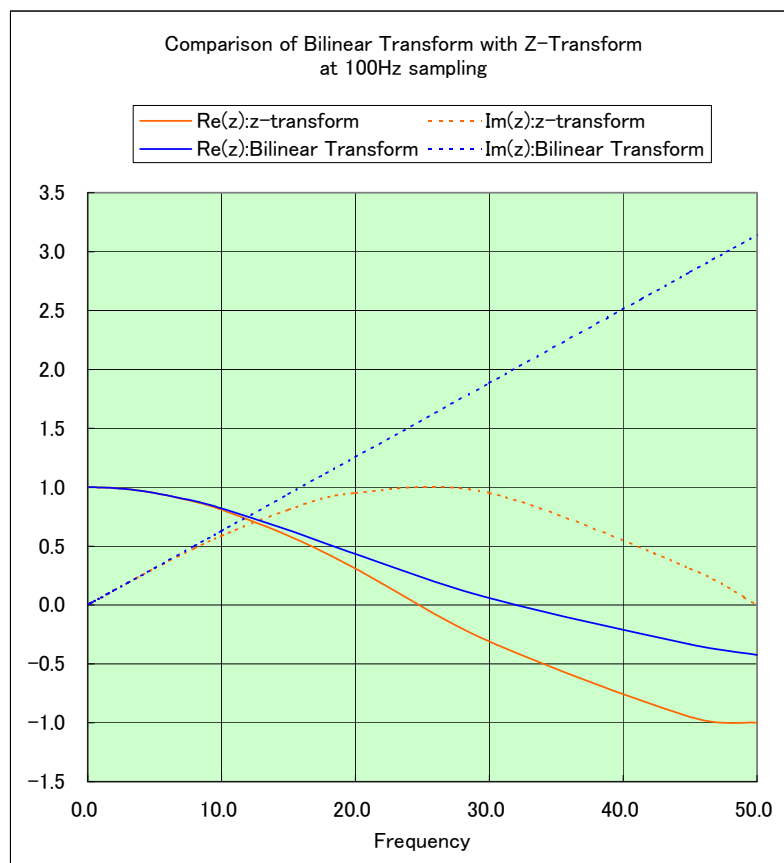


Fig. 44a Example of a comparison of bi-linear transform with z-transform for  $\Delta t = 0.01$  s. The difference is negligible at frequencies less than 10 Hz.

**Example: Simple Moving Coil Type Seismometer (z-transform and equivalent recursive filter)**

Eq. (15) is factorized in the following manner:

$$T(s) = -\frac{X(s)}{Y(s)} = \frac{s^2}{s^2 + 2h\omega_0 s + \omega_0^2} = \frac{(s - s_2)(s - s_1)}{(s - s_2)(s - s_1)}.$$

This reduces to the following with frequency warping as given by Eq. (30).

$$T(s) = \frac{s^2}{s^2 + 2h\omega'_0 s + \omega_0'^2}. \quad (31)$$

Applying the bilinear transformation to Eq. (31) gives a simulated transfer function in the z-domain,

$$\begin{aligned} T(z) &= \frac{\left(\frac{1-z^{-1}}{1+z^{-1}}\right)^2}{\left(\frac{1-z^{-1}}{1+z^{-1}}\right)^2 + 2h \tan\left(\frac{\omega_0 \Delta t}{2}\right) \left(\frac{1-z^{-1}}{1+z^{-1}}\right) + \tan^2\left(\frac{\omega_0 \Delta t}{2}\right)} \\ &= \frac{(1-z^{-1})^2}{(1-z^{-1})^2 + 2h \tan\left(\frac{\omega_0 \Delta t}{2}\right) (1-z^{-1})(1+z^{-1}) + \tan^2\left(\frac{\omega_0 \Delta t}{2}\right) (1+z^{-1})^2} \\ &= \frac{a_0 + a_1 z^{-1} + a_2 z^{-2}}{b_0 + b_1 z^{-1} + b_2 z^{-2}}, \end{aligned} \quad (32)$$

where

$$\begin{aligned} a_0 &= 1.0, a_1 = -2.0, a_2 = 1.0, \\ b_0 &= 1 + 2h \tan(\omega_0 \Delta t / 2) + \tan^2(\omega_0 \Delta t / 2), \\ b_1 &= -2 + 2 \tan^2(\omega_0 \Delta t / 2), \\ b_2 &= 1 - 2h \tan(\omega_0 \Delta t / 2) + \tan^2(\omega_0 \Delta t / 2). \end{aligned}$$

These coefficients of the recursive filter give an approximately equivalent discrete transfer function.

### Exercise: Filter Equivalent to a Simple Moving Coil Type Seismometer—1.

The program *DSEISM.EXE* calculates the coefficients of the recursive filter equivalent to the relative motion of the pendulum mass of a seismometer and applies the recursive filter to an input time series.

- (1) Prepare an input time series by using *TESTSIG.EXE*. An impulse of a unit amplitude at  $t = 0.0$  s will give you the response characteristics of the filter. For example, use  $\Delta t = 0.05$  s and  $N = 128$ .
- (2) Run *DSEISM.EXE* with a natural period  $T_0 = 0.5$  s, damping factor  $h = 0.71$ ,  $\Delta t = 0.05$ , and gain  $G_0 = 1.0$ .
- (3) Draw the filtered time series by using *PTIME.EXE* and its Fourier spectra by using *FFT.EXE* and *PSPEC.EXE*. An example is shown in Fig. 45.1.
- (4) Repeat the above procedure with different values of the natural period and damping factor.

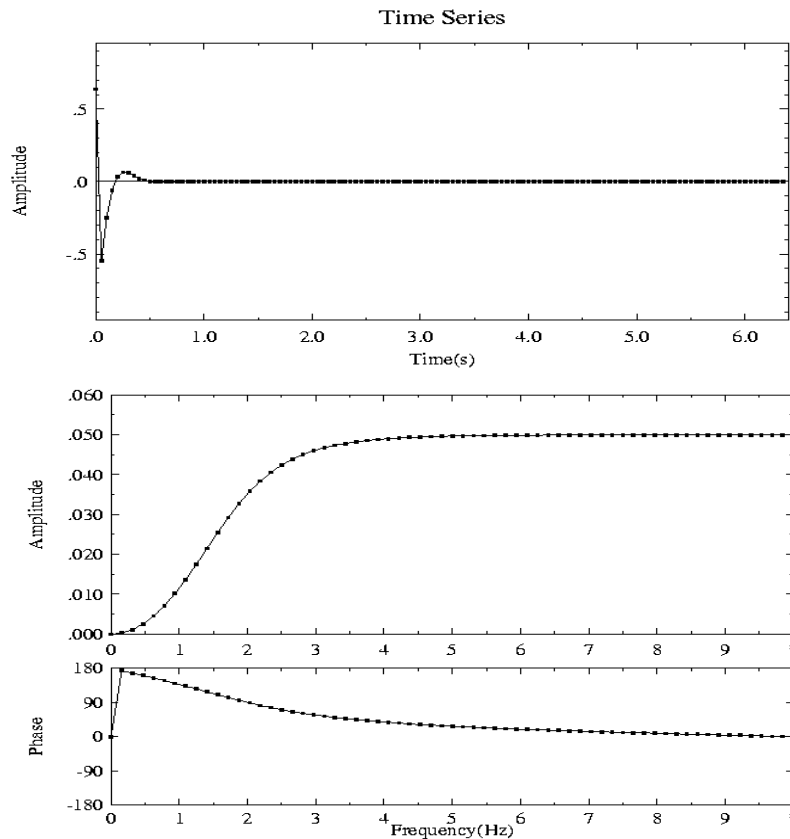


Fig. 45.1 Filtered time series and its Fourier spectra obtained by *DSEISM.EXE*. These are equivalent to the impulse response of the relative motion of a pendulum mass of a seismometer with the natural period  $T_0 = 0.5$  s, the damping factor  $h = 0.71$ ,  $\Delta t = 0.05$ , and gain  $G_0 = 1.0$ . The phase at zero frequency should converge to 180 degrees. However, to avoid division by zero, it is forced to be zero.

**Example: Filter equivalent to a Simple Moving Coil Type Seismometer— 2**

In the previous example, the recursive filter that gives the relative displacement of a pendulum  $-x_m$  for ground displacement  $y_m$  is given. Usually, the data obtained by a digital recorder are given numerically and a constant is given for conversion into volts. The potential difference, which is the output from seismometer, is given as follows:

$$e_m = \frac{G_0 R_s}{R_0 + R_s} (i\omega)x_m,$$

where  $(i\omega)$  shows the effect of differentiation due to a moving coil type transducer;  $R_0$ , the coil resistance;  $R_s$ , the shunt resistance; and  $G_0$ , the product of the sensitivity of the seismometer with the conversion constant of a digital recorder. Therefore, the system response is

$$\frac{e_m}{y_m} = \frac{G_0 R_s}{R_0 + R_s} \cdot \frac{(i\omega)}{1 - 2ih(\omega_0/\omega) - (\omega_0/\omega)^2}.$$

Fig. 45.2 shows its frequency dependency.

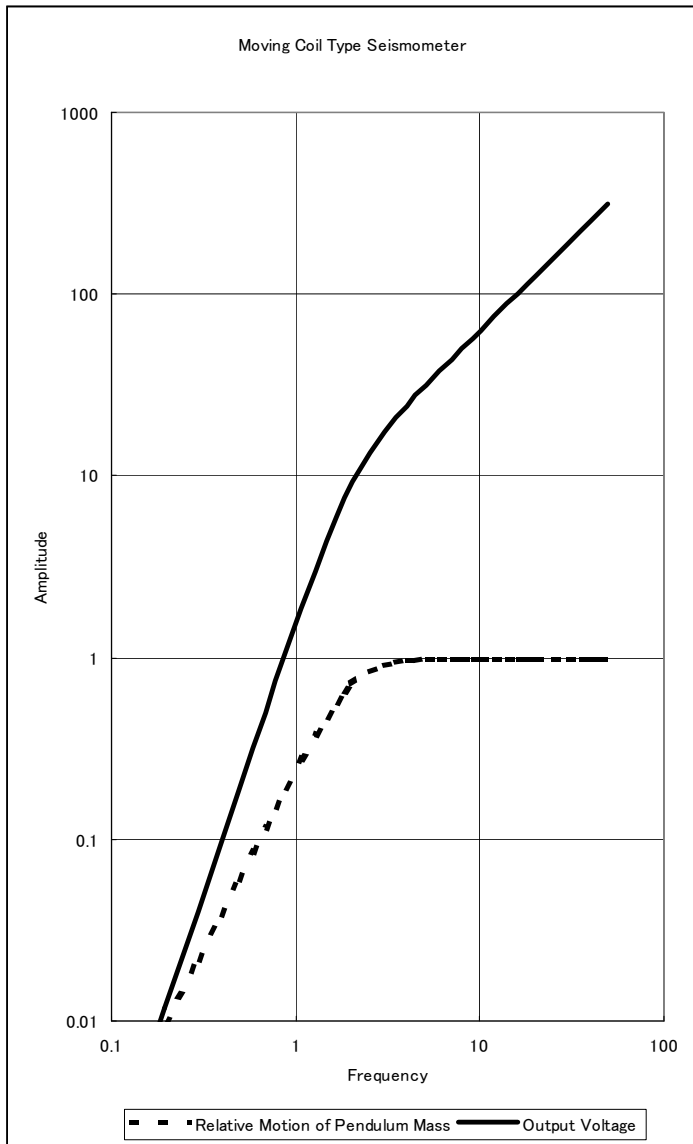


Fig. 45.2 Response of the recursive filters calculated by the formula written above in this page. *Dashed line*: Filter equivalent to the relative motion of pendulum mass of a simple moving coil type seismometer, against ground displacement. *Solid line*: Filter equivalent to the voltage change between two output terminals of a simple moving coil type seismometer, against ground displacement. The parameters used are  $T_0=0.5, h=0.71$ .

This has an equivalent digital filter. The corresponding transfer function in the  $s$ -domain is

$$T(s) = G \cdot \frac{s^3}{s^2 + 2h\omega_0 s + \omega_0^2}, \quad G = \frac{G_0 R_s}{R_0 + R_s}. \quad (33.1)$$

The solutions of the equation, i. e., the denominator is equal to zero, are

$$s = \omega_0 \left( -h \pm \sqrt{h^2 - 1} \right).$$

This gives the pole position at

$$\left( -\omega_0 h, \omega_0 \sqrt{1 - h^2} \right), \left( -\omega_0 h, -\omega_0 \sqrt{1 - h^2} \right) \quad \text{for } h < 1.0, \text{ under-damped case,}$$

$$\left( -\omega_0, 0 \right) \quad \text{doubled for } h = 1.0, \text{ critically damped case,}$$

$$\left( -\omega_0 \left( h - \sqrt{h^2 - 1} \right), 0 \right), \left( -\omega_0 \left( h + \sqrt{h^2 - 1} \right), 0 \right) \quad \text{for } h > 1.0, \text{ over-damped case.}$$

However, the numerator gives a tripled zero at  $(0, 0)$ .

By using the warping frequency, the transfer function is given approximately as follows:

$$T(z) = \frac{a_0 + a_1 z^{-1} + a_2 z^{-2} + a_3 z^{-3}}{b_0 + b_1 z^{-1} + b_2 z^{-2} + b_3 z^{-3}}, \quad (33.2)$$

where

$$a_0 = G, a_1 = -3G, a_2 = 3G, a_3 = -G, .$$

$$b_0 = (\Delta t/2) \cdot \left\{ 1 + 2h \tan(\omega_0 \Delta t/2) + \tan^2(\omega_0 \Delta t/2) \right\},$$

$$b_1 = (\Delta t/2) \cdot \left\{ -1 + 2h \tan(\omega_0 \Delta t/2) + 3 \tan^2(\omega_0 \Delta t/2) \right\},$$

$$b_2 = (\Delta t/2) \cdot \left\{ -1 - 2h \tan(\omega_0 \Delta t/2) + 3 \tan^2(\omega_0 \Delta t/2) \right\},$$

$$b_3 = (\Delta t/2) \cdot \left\{ 1 - 2h \tan(\omega_0 \Delta t/2) + \tan^2(\omega_0 \Delta t/2) \right\}.$$

**Example: Reconstruction of Transfer Function from Given Poles and Zeros**

Suppose that the following data are given for an observation system. In fact, these data are for a STS—2 feedback type seismometer.

Normalization factor:  $A_0 = 5.42787E+07$   
 Normalization frequency: 0.02 (Hz)

**Complex zeros:**

i	real part	imaginary part	Index
0	0.000000E+00	0.000000E+00	$X_0$
1	0.000000E+00	0.000000E+00	$X_1$

**Complex poles:**

i	real part	imaginary part	Index
0	-1.247510E+02	-4.171480E+02	$x_0$
1	-1.247510E+02	4.171480E+02	$x_1$
2	-4.873870E-02	-1.552120E-02	$x_2$
3	-4.873870E-02	1.552120E-02	$x_3$
4	-2.513300E+02	0.000000E+00	$x_4$

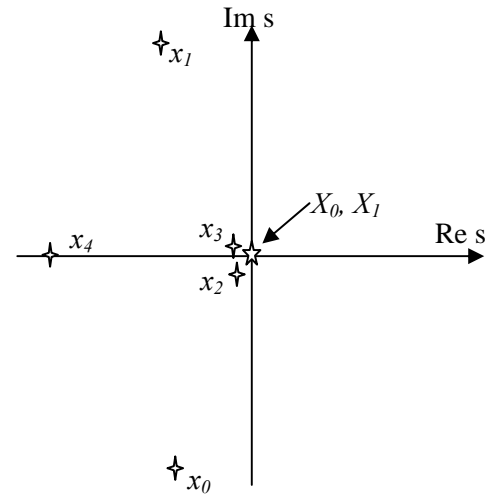


Fig. 45.3 Configuration of poles and zeros, for the example.

Sensitivity:  $G_0 = 6.291456E+08$  (digit/(M/s))

Frequency of sensitivity: 0.02 (Hz)

The normalized amplitude spectra of the transfer function are given as follows:

$$|\tilde{T}(\omega)| = \frac{|s - X_0| \cdot |s - X_1|}{\prod_{m=0}^4 |s - x_m|} = \frac{\omega^2}{\prod_{m=0}^4 \sqrt{(0 - \text{Re}(x_m))^2 + (\omega - \text{Im}(x_m))^2}}$$

The normalizing factor  $A_0$  is the reciprocal of this value at  $f = 0.02$  (Hz).  $A_0 |\tilde{T}(\omega)| = 1.0$  at  $f = 0.02$  Hz.

The reconstructed transfer function is given as follows:

$$|T(\omega)| = G_0 A_0 |\tilde{T}(s)| = \frac{G_0 A_0 \omega^2}{\prod_{m=0}^4 \sqrt{(\text{Re}(x_m))^2 + (\omega - \text{Im}(x_m))^2}}$$

For phase spectra,

$$\text{Arg}(T(\omega)) = \sum_{l=0}^1 \frac{\pi}{2} - \sum_{m=0}^4 \left\{ \tan^{-1} \left( \frac{\omega - \text{Im}(x_m)}{0 - \text{Re}(x_m)} \right) \right\}$$

The first term corresponds to two zeros at the origin. These spectra can be calculated even with a handy calculator.

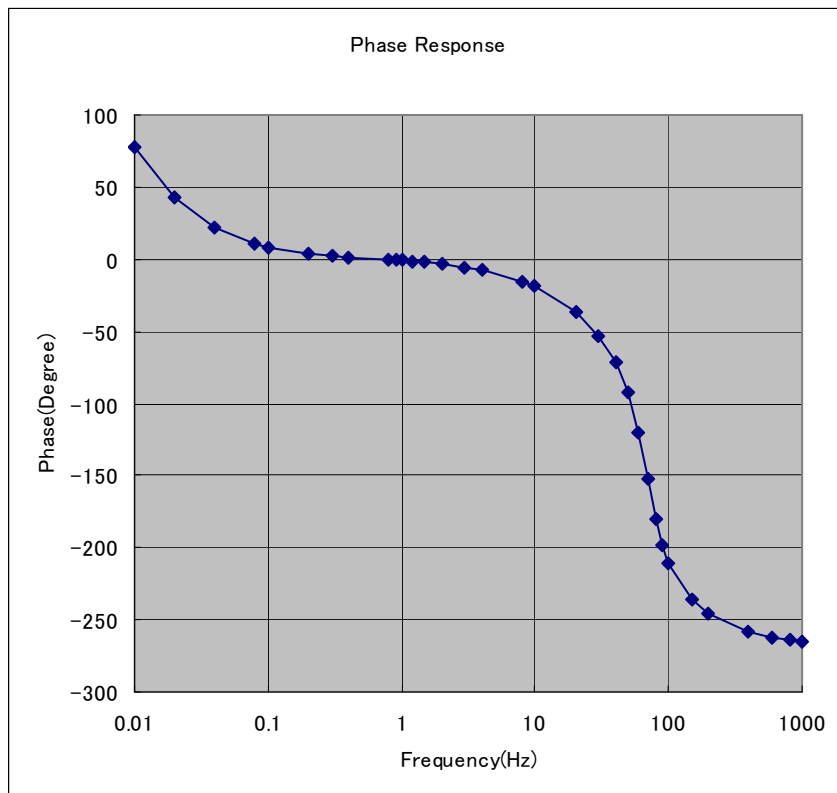
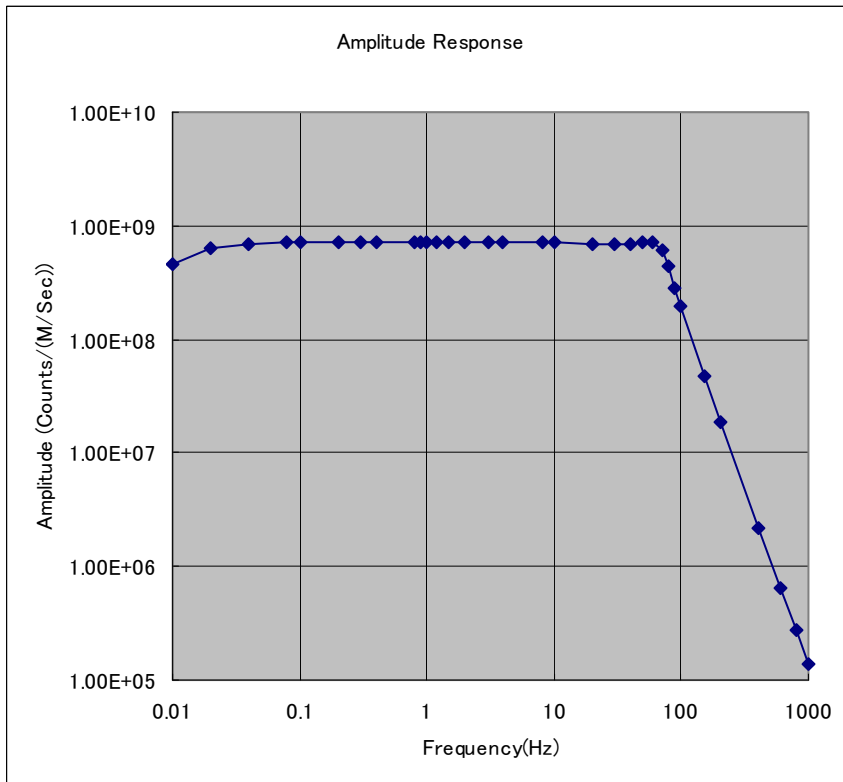


Fig. 45.4 Amplitude (top) and phase (bottom) response of an STS-2 feed back type seismometer. These are calculated by the formula using poles and zeros data given in the previous page. The calculations were performed easily in Microsoft Excel.

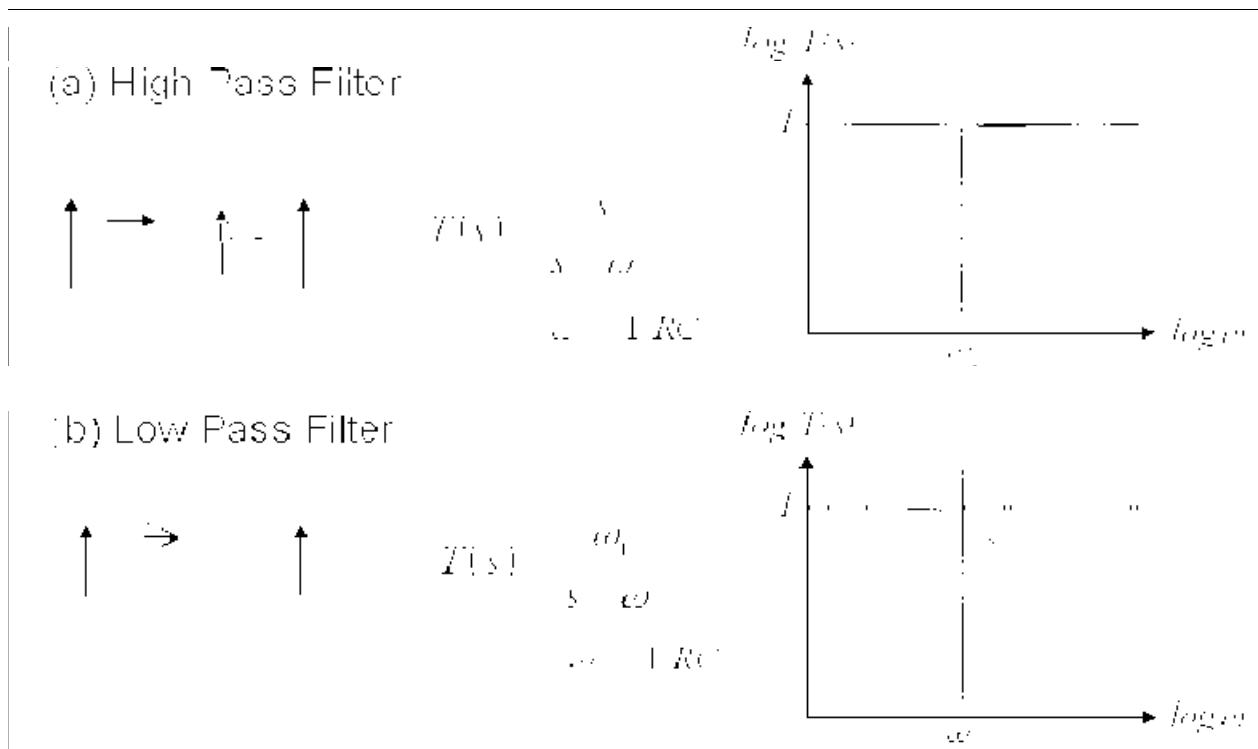


Fig. 46 (Top) RC high pass filter. Its circuit, transfer function, and response characteristics. (Bottom) RC low pass filter.

### 3.4.5. Analog Filters and their Transfer Functions

Several electronic analog filters are commonly used in seismometry. Their transfer functions are introduced here. Naturally, each circuit has its equivalent digital filter.

(1) **RC filter:** One of the simplest electronic circuits is an RC filter that is composed of a resistor and a capacitor, as shown in Fig. 46.

#### High Pass RC Filter

If the resistor is connected parallel to the output (Fig. 46 top), it is a high pass filter. Suppose  $x(t)$  and  $y(t)$  denotes the input and output voltage imbalance, respectively. The equation for a high pass filter is as follows:

$$\begin{aligned}
 y(t) + V_c(t) &= x(t), \\
 I(t) &= C \frac{d}{dt} V_c(t), \\
 RI(t) &= y(t),
 \end{aligned}
 \tag{34.1}$$

where  $R$  and  $C$  denotes the resistance and the capacitance integrated in the circuit (Fig. 46 top left);  $I(t)$ , the current that passes through  $R$  and  $C$ ; and  $V_c$ , the voltage across the capacitor. This formula gives

$$\frac{d}{dt} y(t) + \frac{1}{CR} y(t) = \frac{d}{dt} x(t).
 \tag{34.2}$$



The transfer function is given by the Laplace transform of Eq. (34.2).

$$T(s) = \frac{s}{s + \omega_0}, \text{ where } \omega_0 = 1/RC. \quad (34.3)$$

The amplitude spectra are drawn schematically in Fig. 46 top right.

**Low Pass RC Filter**

If the capacitor is connected parallel to the output, it works as a low pass filter (Fig. 46 bottom).

The low pass filter shown in Fig. 46 bottom left corresponds to the following equation.

$$RI(t) + y(t) = x(t),$$

$$I(t) = C \frac{d}{dt} y(t), \quad (35.1)$$

This can be combined into

$$RC \frac{d}{dt} y(t) + y(t) = x(t). \quad (35.2)$$

The transfer function is given by the Laplace transform of (35.2).

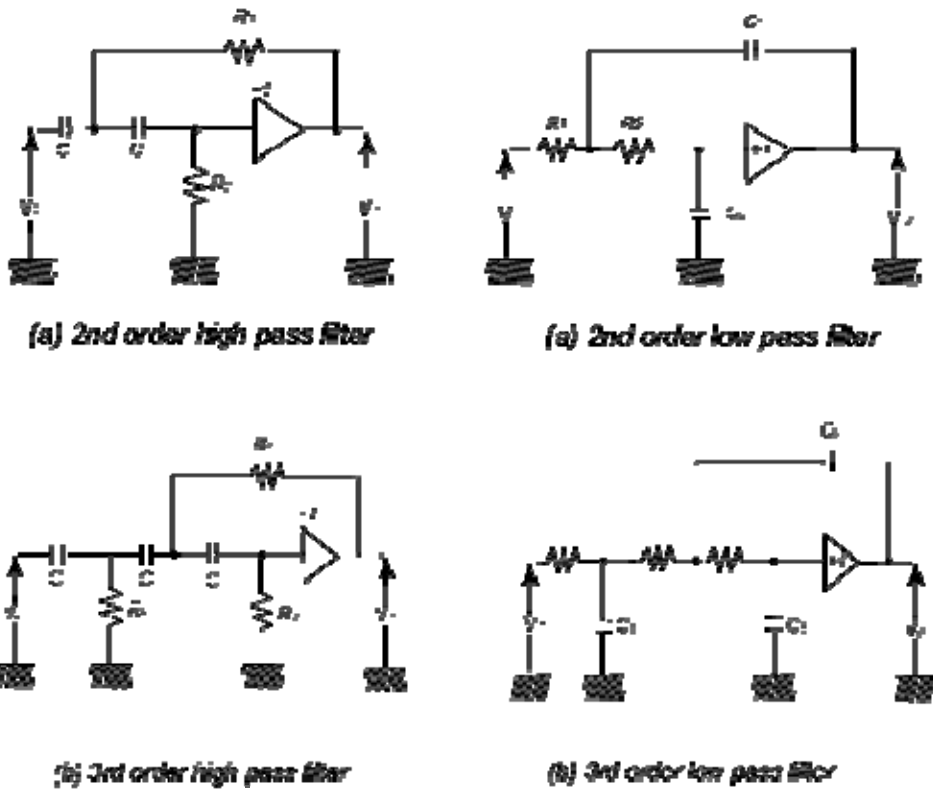


Fig. 47 Schematic circuit of an active filter composed of OpAmps. Reproduced based on Yanagisawa and Kanematsu (1980)

$$T(s) = \frac{\omega_c}{s + \omega_c}, \text{ where } \omega_0 = 1/RC. \quad (35.3)$$

The amplitude spectra are drawn schematically in Fig. 46 *bottom right*.

**(2) Active filter.**

Today, many electronic circuits with known characteristics are widely used for filtering. Some of those characteristics have their proper names. The requirements for the filters may be flat amplitude characteristics in the pass band, sharp cutoff, and flat delay characteristics in the pass band. The last implies linear phase characteristics because delay is defined as phase differentiated by the angular frequency.

A **Butterworth filter** has its amplitude characteristics as

$$|T(s)| = \frac{1}{\sqrt{1 + (\omega/\omega_c)^{2n}}},$$

where  $\omega_c$  is the cutoff angular frequency.  $n$  denotes the order of the filter and the slope of amplitude characteristics in the stop frequency band. In other words, an  $n$ -order filter has a decay of  $20 n$  db/oct. The transfer function itself is given for a low pass filter

$$T(s) = \Pi \frac{\Omega_0^2}{(s/\omega_c)^2 + (\Omega_0/Q)(s/\omega_c) + \Omega_0^2}, \text{ for even } n,$$

$$T(s) = \frac{\Omega_a}{(s/\omega_c) + \Omega_a} \Pi \frac{\Omega_0^2}{(s/\omega_c)^2 + (\Omega_0/Q)(s/\omega_c) + \Omega_0^2}, \text{ for odd } n,$$

and for a high pass filter

$$T(s) = \Pi \frac{\Omega_0^2}{(\omega_c/s)^2 + (\Omega_0/Q)(\omega_c/s) + \Omega_0^2}, \text{ for even } n,$$

$$T(s) = \frac{\Omega_a}{(\omega_c/s) + \Omega_a} \Pi \frac{\Omega_0^2}{(\omega_c/s)^2 + (\Omega_0/Q)(\omega_c/s) + \Omega_0^2}, \text{ for odd } n,$$

where  $\omega_c$  is the cut-off angular frequency; other coefficients are given in the following table.

Table 10 Coefficients for *Butterworth* filter

$n$	$Q$	$\Omega_0$	$\Omega_a$	$C_1$	$C_2$	$C_3$
2	0.707107	1.000000	---	1.4142	0.7071	---
3	1.000000	1.000000	1.000000	1.3926	3.5468	0.2025
4	0.541196	1.000000	---	1.0824	0.9239	---
	1.306536	1.000000	---	2.6131	0.3827	---
5	0.618034	1.000000	1.000000	1.3541	1.7529	0.4213
	1.618034	1.000000	---	3.2361	0.3090	---

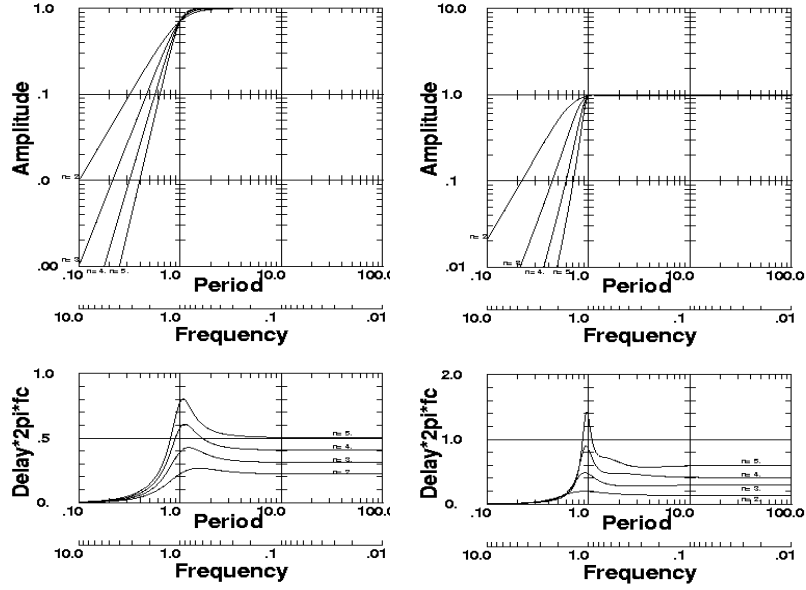


Fig. 48 Amplitude and delay of (a) Butterworth filter and (b) Chebyshev filter in the frequency domain. Both are plotted for the case of a low pass filter

The value of  $C_n$  ( $n = 1, 2, 3$ ) in Table 10 corresponds to the actual value of capacitor  $C_a$  ( $a = 1, 2, 3$ ) in Fig. 47 (right) but normalized one, i.e.,  $C_a = C_n/\omega_c R$ , where  $R$  is the resistance selected in advance. For high pass filter (Fig. 47 left),  $R_n$  ( $n = 1, 2, 3$ ) are calculated by  $1/C_n$  ( $n = 1, 2, 3$ ) in Table 10. The value of the actual resistor  $R_a$  ( $a = 1, 2, 3$ ) in Fig. 47 (left) is given by  $R_a = R_n/\omega_c C$ , where  $C$  is the value of the capacitance selected in advance.

**Butterworth filter** has plane amplitude characteristics in the pass frequency band and is used for shaping of the spectra. The delay characteristics of this filter, however, are not flat even in the pass band. These have a peak of delay around the cut off frequency. This means that the shape of the signal around the cut off frequency in the time domain is considerably distorted by filtering operation.

**Chebyshev filter** : If we allow the ripple in the pass band, we can perform a sharp cut off. *Chebyshev* filter belongs to this category. The transfer function is given for low pass filter

$$T(s) = \frac{1}{\sqrt{10^{R_p/10}}} \prod \frac{\Omega_0^2}{(s/\omega_c)^2 + (\Omega_0/Q)(s/\omega_c) + \Omega_0^2}, \text{ for even } n,$$

$$T(s) = \frac{\Omega_a}{(s/\omega_c) + \Omega_a} \prod \frac{\Omega_0^2}{(s/\omega_c)^2 + (\Omega_0/Q)(s/\omega_c) + \Omega_0^2}, \text{ for odd } n,$$

and for a high pass filter

$$T(s) = \frac{1}{\sqrt{10^{R_p/10}}} \prod \frac{\Omega_0^2}{(\omega_c/s)^2 + (\Omega_0/Q)(\omega_c/s) + \Omega_0^2}, \text{ for even } n,$$

$$T(s) = \frac{\Omega_a}{(\omega_c/s) + \Omega_a} \prod \frac{\Omega_0^2}{(\omega_c/s)^2 + (\Omega_0/Q)(\omega_c/s) + \Omega_0^2}, \text{ for odd } n,$$

where  $R_p$  is the ripple amplitude in a pass band measured in (db).

The coefficients are given in the following table.

Table 11 Coefficients for Chebyshev filter with ripple of 0.25 db in the pass band

$n$	$Q$	$\Omega_0$	$\Omega_a$	$C_1$	$C_2$	$C_3$
2	0.809254	1.453972	---	1.1132	0.4249	---
3	1.508026	1.156992	0.767223	1.6110	6.8272	0.0885
4	0.657249	0.674422	---	1.9491	1.1280	---
	2.536110	1.077939	---	4.7055	0.1829	---
5	1.035932	0.732405	0.436951	2.6625	5.0919	0.3147
	3.875683	1.046630	---	7.4060	0.1233	---

The delay characteristics of this filter are not flat in the pass band. In such a case, the filtering operation causes considerable distortion of waveform in the time domain.

**Bessel filter:** It can produce flat delay characteristics in the pass band. However, the cutoff is not sharp. The amplitude characteristics are also flat in the pass band. The transfer function is given by the same formulas as those of a **Butterworth** filter. The coefficients are given in the following table.

Table 12 Coefficients for Bessel filter

$n$	$Q$	$\Omega_0$	$\Omega_a$	$C_1$	$C_2$	$C_3$
2	0.577350	1.732051	---	0.6667	0.5000	---
3	0.691047	2.541547	2.322165	0.5647	0.8136	0.1451
4	0.521935	3.023265	---	0.3453	0.3169	---
	0.805538	3.389366	---	0.4753	0.1831	---
5	0.563536	3.777894	3.646739	0.3601	0.4171	0.1280
	0.916479	4.261031	---	0.4302	0.1280	---

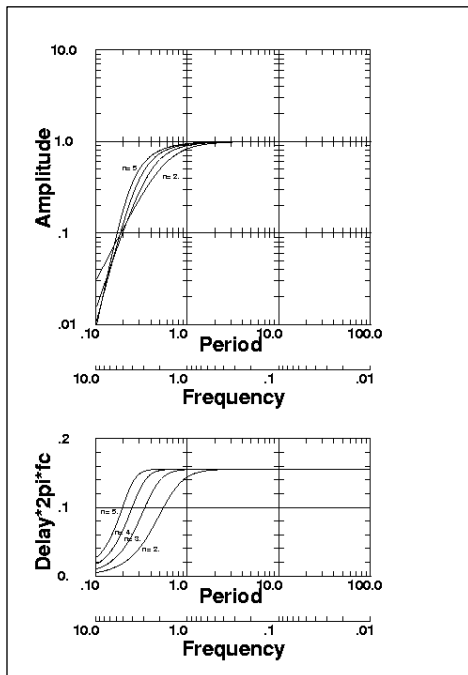


Fig. 48 (continued) (c) Amplitude and delay of the Bessel filter in the frequency domain plotted for the case of a low pass filter.

**Example: Recursive filter equivalent to analog filter circuit**

**Low Pass Filter:** The transfer function of the second-order analog low pass Butterworth filter circuit is

$$T_{2L}(s) = \frac{\Omega_0^2}{(s/\omega_c)^2 + (\Omega_0/Q)(s/\omega_c) + \Omega_0^2}.$$

The warped cutoff angular frequency and bi-linear transform are applied. Then,

$$s/\omega_c' = \frac{1-z^{-1}}{\alpha(1+z^{-1})}, \quad \alpha = \tan\left(\frac{\omega_c \Delta t}{2}\right).$$

Therefore,

$$T_{2L}(z) = \frac{a_0 + a_1 z^{-1} + a_2 z^{-2}}{b_0 + b_1 z^{-1} + b_2 z^{-2}}, \quad \begin{aligned} a_0 &= 1, a_1 = 2, a_2 = 1, b_0 = 1 + (1/\Omega_0 \alpha)/Q + (1/\Omega_0 \alpha)^2, \\ b_1 &= 2\{1 - (1/\Omega_0 \alpha)^2\}, b_2 = 1 - (1/\Omega_0 \alpha)/Q + (1/\Omega_0 \alpha)^2. \end{aligned}$$

The third-order filters can be written as

$$T_{3L}(s) = T_{1L}(s) \cdot T_{2L}(s), \quad T_{1L}(s) = \frac{\Omega_a}{(s/\omega_c) + \Omega_a}$$

The warped cutoff angular frequency and bi-linear transform give

$$T_{1L}(z) = \frac{c_0 + c_1 z^{-1}}{d_0 + d_1 z^{-1}}, \quad c_0 = 1, c_1 = 1, d_0 = 1 + 1/(\alpha \Omega_a), d_1 = 1 - 1/(\alpha \Omega_a).$$

Third-order filtering can be achieved by applying  $T_{1L}(z)$  and  $T_{2L}(z)$  sequentially.

It is also possible to apply it immediately by using the following.

$$T_{3L}(z) = T_{1L}(z) \cdot T_{2L}(z) = \frac{a_0 c_0 + (a_1 c_0 + a_0 c_1)z^{-1} + (a_2 c_0 + a_1 c_1)z^{-2} + a_2 c_1 z^{-3}}{b_0 d_0 + (b_1 d_0 + b_0 d_1)z^{-1} + (b_2 d_0 + b_1 d_1)z^{-2} + b_2 d_1 z^{-3}}.$$

**High Pass Filter:** The transfer function of an analog second-order high pass Butterworth filter is

$$T_{2H}(s) = \frac{\Omega_0^2}{(\omega_c/s)^2 + (\Omega_0/Q)(\omega_c/s) + \Omega_0^2}.$$

The equivalent recursive filter is

$$T_{2H}(z) = \frac{a_0 + a_1 z^{-1} + a_2 z^{-2}}{b_0 + b_1 z^{-1} + b_2 z^{-2}}, \quad \begin{aligned} a_0 &= 1, a_1 = -2, a_2 = 1, b_0 = 1 + (\alpha/\Omega_0)/Q + (\alpha/\Omega_0)^2, \\ b_1 &= 2\{(\alpha/\Omega_0)^2 - 1\}, b_2 = 1 - (\alpha/\Omega_0)/Q + (\alpha/\Omega_0)^2. \end{aligned}$$

The third-order filters can be written as

$$T_{3H}(s) = T_{1H}(s) \cdot T_{2H}(s), \quad T_{1H}(s) = \frac{\Omega_a}{(\omega_c/s) + \Omega_a}$$

The warped cutoff angular frequency and bi-linear transform give

$$T_{1H}(z) = \frac{c_0 + c_1 z^{-1}}{d_0 + d_1 z^{-1}}, \quad c_0 = 1, c_1 = -1, d_0 = 1 + \alpha/\Omega_a, d_1 = -1 + \alpha/\Omega_a.$$

These formulas are valid also for Chebyshev and Bessel filters because their transfer functions are

composed of  $T_{1L}(s), T_{2L}(s), T_{3L}(s)$  or  $T_{1H}(s), T_{2H}(s), T_{3H}(s)$ .

### 3.4.6. Exercise for Digital Filtering

The transfer function of a given recursive filter is easily obtained by using Eqs. (20) and (25). However, it is not easy to design the filter's coefficients; to achieve this, its amplitude and phase characteristics must be arranged in the desired manner. Saito (1978) has published a subroutine package written in Fortran, which gives the coefficients of a recursive filter whose frequency characteristics coincide with one of the four filters popularly used in electronics, i.e., Butterworth, Chebyshev-I (constant ripple in pass band), Chebyshev-II (constant ripple in stop band), and Elliptic filters. Each of these four can be arranged as high pass, low pass, band pass, and band stop filters, respectively. These programs have been provided as a free software.

However, it is requested that users of the software acknowledge this by stating that **“the subroutines for digital filtering published by Saito (1978) are used for processing”** with the reference **“ Saito, M. (1978):An automatic Design Algorithm for Band Selective Recursive Digital Filters, BUTURI-TANSA, Vol. 31, No. 4, pp112-135 (in Japanese).”** The distributed programs have been compiled and linked and provided as the library for *g77* on *Cygwin*. The filtering operation can be performed by just calling the subroutine *BANDPI* in the main routine as follows:

*CALL BANDPI (X,N,DT,FL,FH,FS,AP,AS, ntype, nchara,n causal)*

where *X* : Input time series / Filtered time series,  
*N* : Number of data included in *X*,  
*DT* : Sampling interval,  
*FL* : Lower limit frequency of the pass-band,  
*FH* : Upper limit frequency of the pass-band,  
*FS* : Limit frequency of the stop-band,  
*AP* :  $1/(1 + A_p^2)$  denotes the ripple in the pass-band,  
*AS* :  $1/(1 + A_s^2)$  denoted the ripple in the stop-band.

$A_p$  and  $A_s$  are schematically indicated in Fig. 49.

*NTYPE* : Flag that indicates the type of filter

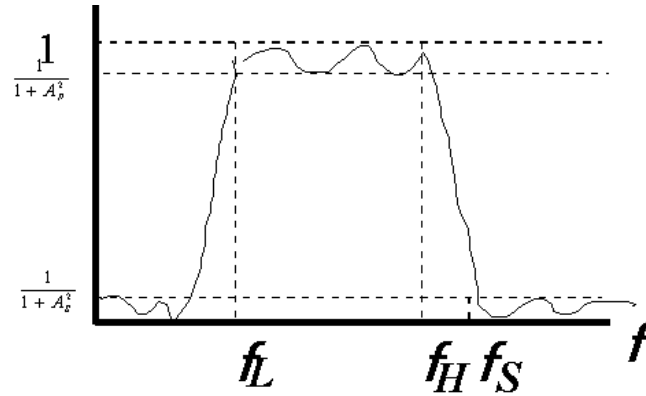
1	Butterworth filter
2	Chevyshev-I
3	Chevyshev-II
4	Elliptic

*NCHARA*: Flag that indicates the frequency characteristics of filter

1	Low cut (high pass)
2	High cut (low pass)
3	Band pass
4	Band stop

*NCAUSAL*: Flag that indicates the causality of filter

1	Causal
2	Zero phase



Zero-phase filtering always breaks the causality. For example, a small ringing caused by the filtering, appears before the initial break of the P-wave and makes it difficult to judge the arrival time. Causal filtering can be performed by these identical subroutines. The flag *NCAUSAL* controls the subroutine *GNF* for both causal and non-causal filtering.

**Exercise:**

- (1) Construct a time series that has an impulse of unit amplitude at  $t = 1.28$  s with  $\Delta t = 0.01$ ,  $N = 256$  and store the data in the file *UT* by using *TESTSIG*.
- (2) Draw the time series *UT* and check it by using *GSTOOLS*.
- (3) Apply FFT to the time series *UT* by using *FFT*.
- (4) Draw the Fourier spectra from *UF* by using *PSPEC*.
- (5) Apply a digital filter by using *TRFILT*. The output file name is *UT1*; A Butterworth type low cut causal filter with  $FS = 15.0$  Hz,  $FL = 30.0$  Hz,  $AP = 10.0$ ,  $AS = 0.1$  is selected for the following example.

```
a:\> TRFILT UT UT1
??? Filter Type ???
    Butterworth type      ==> 1
    Chebyshev-I type      ==> 2
    Chebyshev-II type     ==> 3
    Elliptic type         ==> 4

1
??? CHARACTERISTICS OF FILTER ???
    Low Cut (High Pass) Filter ==> 1
    High Cut (Low Pass) Filter ==> 2
    Band Pass Filter ==> 3
    Band Stop Filter ==> 4

1
      FL
      /-----
-----/
      FS
??? FS,FL,AP,AS ???
    AP,AS: Parameter defining the ripple
           in pass band and stop band.
           Use AP = 0.1, AS = 10.0, if you do
           not like to think.

15. 30. 0.1 10.0
??? Causal or Zero-Phase ???
    Causal Filter ==> 1
    Zero-Phase Filter ==> 2

1
```

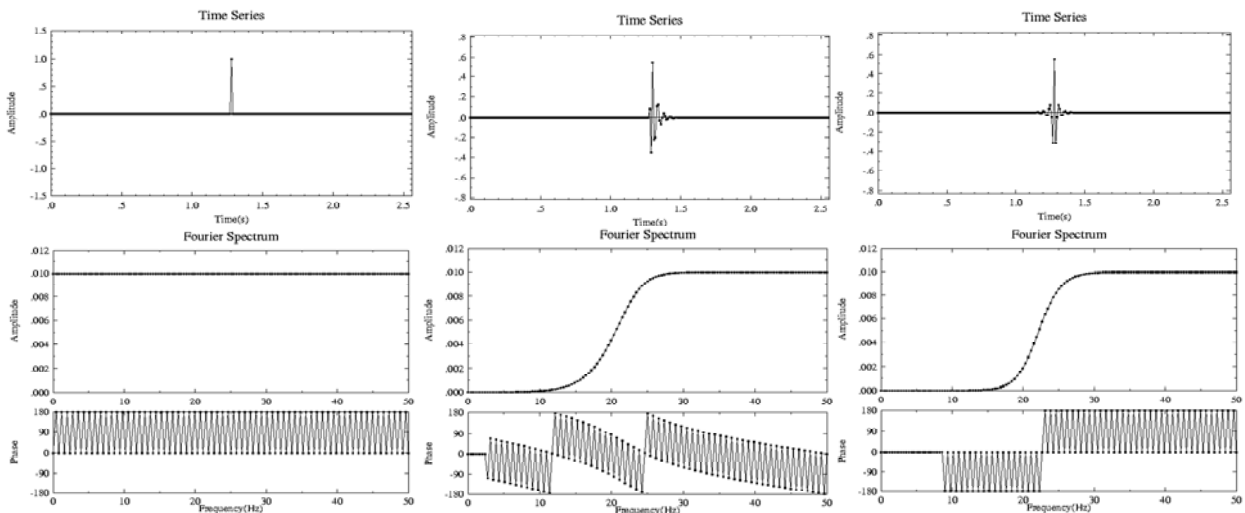


Fig. 50 Examples of the performance of the recursive filter designed by *TRFILT.EXE*. *Left*: the input signal that is an impulse of unit amplitude located at  $t = 1.28$  s. *Center*: the output from the causal Butterworth low pass filter. *Right*: the output from the zero phase Butterworth low pass filter. The parameters used are  $\Delta t = 0.01$ ,  $N = 256$ ,  $F_S = 15.0$  Hz,  $F_L = 30.0$  Hz,  $A_p = 0.1$ ,  $A_s = 10.0$ .

- (6) Draw the filtered time series *UTI* by using *PTIME*.
- (7) Apply *FFT* to the filtered time series *UTI* and store the result in the file *UF1*.
- (8) Draw the Fourier spectra stored in *UF1* by using *PSPEC* and compare it with the drawn *UF*.
- (9) Check the performance for filters of other types by repeating the above procedure with different values of control parameters.



### 3.5. Deconvolution or Inverse Filtering

Every observed seismic waveform data is an output from some filters. For example, the seismometer itself is a high pass filter. Amplifiers, sometimes, have frequency characteristics that are not uniform for all frequencies. Moreover, we use other high cut filter to suppress AC noise that has energy at 50 or 60 Hz since these distort seismic signals. Moreover, an anti-aliasing filter is used in digital data acquisition systems. Since seismometry is conducted not only for obtaining travel time data but also for obtaining waveform data, i. e., ground motion, we must compensate this distortion in order to obtain true ground motion. Remember that recorded signals are obtained by the convolution of the instrumental characteristics to the ground motion in the time domain and are given by their product in the frequency domain. The instrumental characteristics are identical to the response of instruments to an impulsive signal (Fig. 51 *top*).

$$\begin{aligned} r(t) &= f(t) * g(t), \\ R(\omega) &= F(\omega)G(\omega). \end{aligned} \tag{36}$$

where  $r(t)$ ,  $f(t)$ , and  $g(t)$  denote the recorded signals, the instrumental response, and the ground motion in the time domain, respectively;  $R(\omega)$ ,  $F(\omega)$ , and  $G(\omega)$ , respectively, are the same parameters in the frequency domain.

*Deconvolution* in seismology is usually the process for eliminating the effect of the instrumental characteristics from the observed data and to recover the true ground motion. Mathematically, this is the reverse process of convolution. Deconvolution in the time domain corresponds to the quotient in the frequency domain. In comparison with the complexity of deconvolution in the time domain, the frequency domain operation is composed of only three steps. These are for applying the FFT to the instrument response and recorded signal, to divide the recorded signal spectra by the instrumental response spectra, and to apply the inverse FFT to the quotient (Fig. 51 *middle*).

$$G(\omega) = \{F(\omega)\}^{-1} \cdot R(\omega). \tag{37}$$

However, we must consider the following. The information once lost in the observation or processing can never be recovered by any technique, even by extremely sophisticated and efficient ones. This is because we cannot avoid noises that are recorded simultaneously with the signal or those that invade into the record during the processing. Once the signals weaken and become smaller than the noise level, any recovery process only amplifies such noises. Such amplified noises can be dominant in the recovered ground motion. The signals weaken a little either only during the recording or the processing can be strengthened or recovered by the deconvolution technique.

Typically, true ground motion has band limited feature at the far field and low pass feature at the near field, whereas the ground noise is present at every frequency. Let us handle only far field ground motion in order to have a simple demonstration. ,  $G(\omega)$  is band limited. The instrumental response  $F(\omega)$  is also band limited, because a seismometer is a low cut filter and used with a high cut anti-alias filter. The recovering operator in the frequency domain  $\{F(\omega)\}^{-1}$  has a large amplitude at a frequency outside this limited frequency band. The recorded signal  $R(\omega)$  is almost band limited but it has little energy outside the frequency band of  $G(\omega)$ , i.e., the contribution of the noise. By applying the inverse filter, i.e., the recovering operator  $\{F(\omega)\}^{-1}$ , this small contribution of the noise will be considerably amplified and

contaminates the recovered ground motion transformed into the time domain by the inverse FFT (Fig. 51 bottom).

This shows that we have to select a frequency range such that the signal is sufficiently larger than the noise in order to prevent instability due to the application of the inverse filter. We cannot recover the ground motion outside this frequency band.

However, the recommended method is to handle the data only within the pass band of the instrumental response. Usually, this is given by the natural frequency of the seismometer and the cutoff frequency of the anti-alias filter. It is possible but not easy to use the information outside of this range. Even within this frequency range, the above mentioned instability problem can take place due to the smaller frequency band of  $G(\omega)$ . Thus, it is also recommended that the shape of  $R(\omega)$  be observed and the useful frequency band be selected before beginning the data processing.

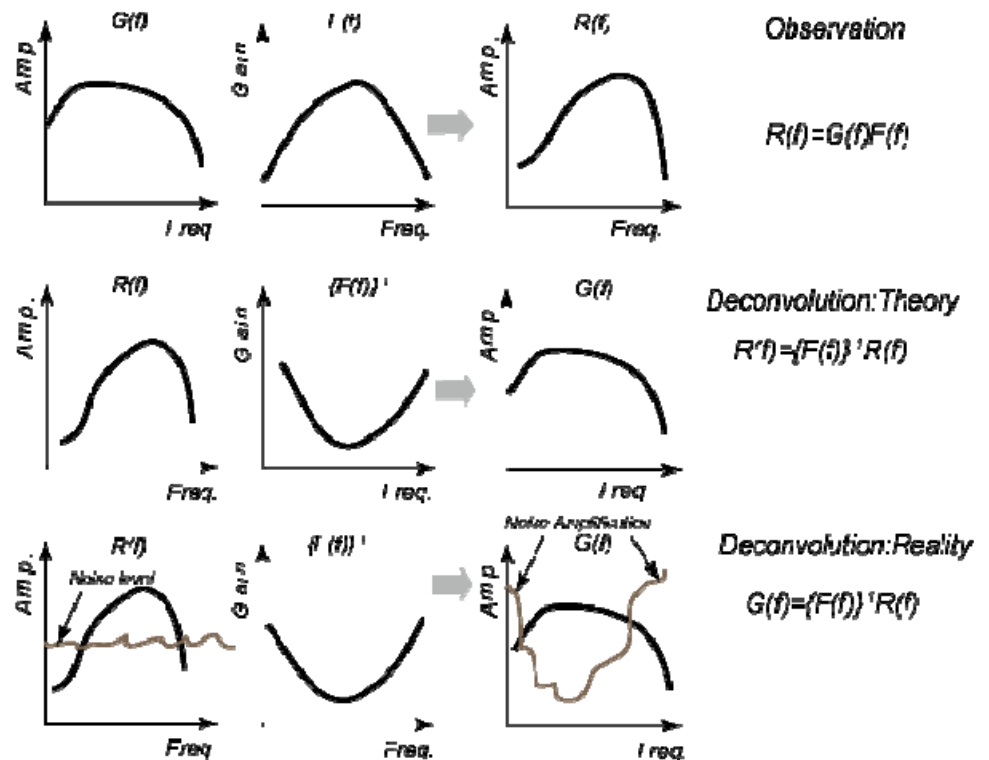


Fig. 51 Schematic illustration of Eq. (36) (top), theoretical inverse filtering process given by Eq. (37) (middle) and actual inverse filtering process with noise (bottom). Reproduced based on Sherbaum (1996).

### Example: Inverse Filter for a Simple Moving Coil Type Seismometer

The transfer function that gives the inverse response to the previous examples is given as follows:

$$T(s) = G^{-1} \cdot \frac{s^2 + 2h\omega_0 s + \omega_0^2}{s^3}, \quad G = \frac{G_0 R_s}{R_0 + R_s}. \quad (36)$$

Fig. 52 shows the amplitude response of this transfer function.

The solutions of the equation, i.e., the numerator is equal to zero, are  $s = \omega_0(-h \pm \sqrt{h^2 - 1})$ . This gives the zero position at

$$\left(-\omega_0 h, \omega_0 \sqrt{1 - h^2}\right), \left(-\omega_0 h, -\omega_0 \sqrt{1 - h^2}\right), \text{ for } h < 1.0, \text{ the under-damped case,}$$

$$\left(-\omega_0, 0.\right) \quad \text{doubled for } h = 1.0, \text{ critically damped case,}$$

$$\left(-\omega_0(h - \sqrt{h^2 - 1}), 0.\right), \left(-\omega_0(h + \sqrt{h^2 - 1}), 0.\right) \text{ for } h > 1.0, \text{ over-damped case.}$$

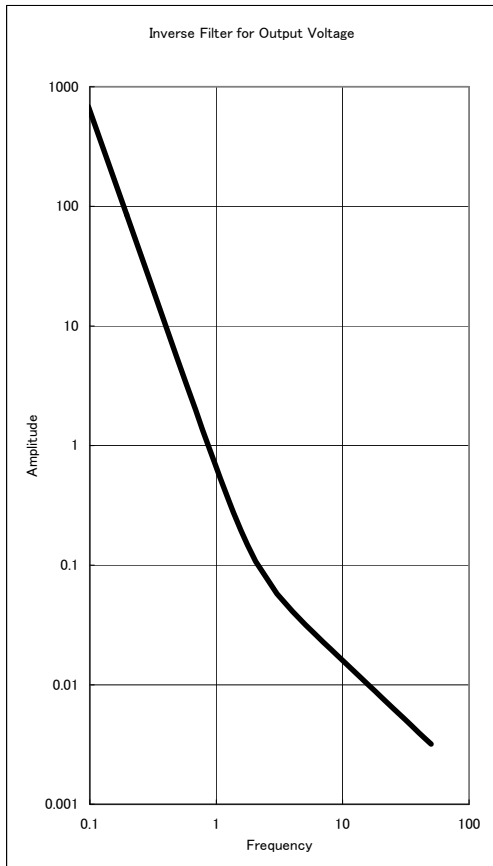
The denominator, however, gives a tripled pole at (0, 0).

By using the warped frequency, the transfer function is given approximately as follows.

$$T(s) = \frac{a_0 + a_1 z^{-1} + a_2 z^{-2} + a_3 z^{-3}}{b_0 + b_1 z^{-1} + b_2 z^{-2} + b_3 z^{-3}},$$

where

$$\begin{aligned} a_0 &= (\Delta t/2) \cdot \{1 + 2h \tan(\omega_0 \Delta t/2) + \tan^2(\omega_0 \Delta t/2)\}, \\ a_1 &= (\Delta t/2) \cdot \{-1 + 2h \tan(\omega_0 \Delta t/2) + 3 \tan^2(\omega_0 \Delta t/2)\}, \\ a_2 &= (\Delta t/2) \cdot \{-1 - 2h \tan(\omega_0 \Delta t/2) + 3 \tan^2(\omega_0 \Delta t/2)\}, \\ a_3 &= (\Delta t/2) \cdot \{1 - 2h \tan(\omega_0 \Delta t/2) + \tan^2(\omega_0 \Delta t/2)\} \\ b_0 &= G, b_1 = -3G, b_2 = 3G, b_3 = -G. \end{aligned}$$



Since a simple moving coil type seismometer is a low cut filter, its inverse filter makes the component grow in a frequency range that is lower than the natural frequency of the seismometer. This causes an instability in the output from the filter, because the signal to noise ratio is small at this frequency range.

Fig. 52 Amplitude response of the inverse filter for the voltage change between two output terminals of a simple moving coil type seismometer that is calculated by Eq. (36). The parameters used are  $T_0 = 0.5$ ,  $h = 0.71$ . Theoretically, this filter can convert the output voltage variation with the ground displacement. However, this amplifies the low frequency components considerably as shown clearly in this figure.

**Exercise: Design of a recursive filter equivalent to a seismometer (simulated seismogram).**

The program *ISEISM.EXE* calculates the coefficients of the abovementioned inverse filter and applies the recursive filter to an input time series. Fig. 53.1 shows the following procedure.

(1) Prepare an input time series by using *TESTSIG.EXE*. For example, an impulse of the unit amplitude located at  $t = 0.0$  s, with  $\Delta t = 0.05$  s,  $N = 128$  (Output: UT0). Then, apply a band pass filter by using *TRFILT.EXE* with this impulse. For example, for the Butterworth type,  $F_L = 0.3$ Hz,  $F_H = 8.0$  Hz,  $F_S = 9.5$  Hz,  $A_P = 0.1$ ,  $A_S = 10.0$  (Input: UT0, Output: UT1). This band pass impulse is used for artificial ground motion.

Calculate the Fourier transform (Input: UT1, Output: UF1)

(2) Run *VSEISM.EXE* with the natural period  $T_0 = 0.25$  s, damping factor  $h = 0.71$ ,  $\Delta t = 0.05$ , and gain  $G_0 = 1.0$  (Input: UT1, Output: UT2). This output is the simulated seismogram for this exercise. Calculate the Fourier transform by using *FFT.EXE* (Input: UT2, Output: UF2).

(4) Run *ISEISM.EXE* with the natural period  $T_0 = 0.25$  sec, the damping factor  $h = 0.71$ ,  $\Delta t = 0.05$ , and gain  $G_0 = 1.0$  (Input: UT2, Output: UT3). Calculate the Fourier transform by using *FFT.EXE* (Input: UT3, Output: UF3).

(5) Draw the time series UT1 and UT3 by using *PTIME.EXE* and its Fourier spectra UF1 and UF3 by using *FFT.EXE* and *PSPEC.EXE* and compare them. Consider the reason why the band pass filter is applied.

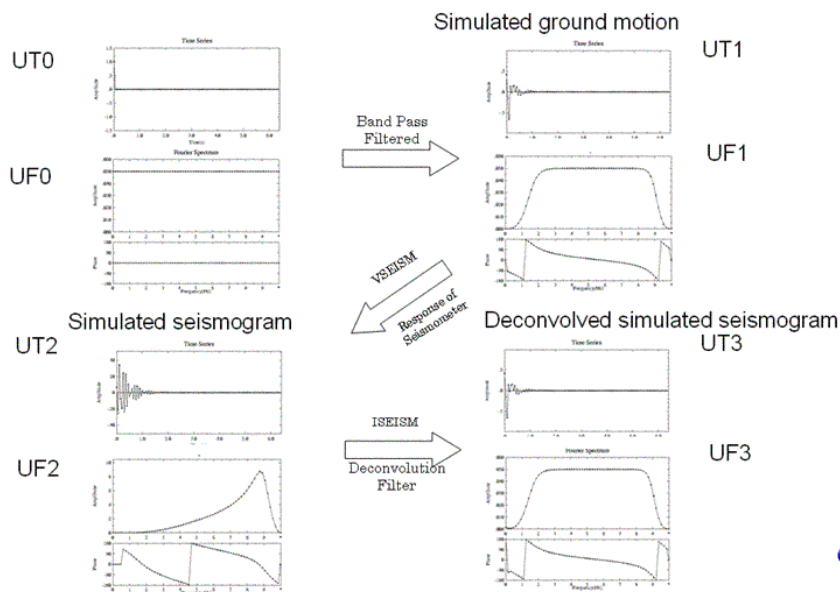


Fig. 53.1 Waveforms and its Fourier transform of each step of the example explained above. *Upper-Left*: Impulse. This is the beginning of the exercise, *Upper-Right*: Band pass filtered impulse, *Bottom Left*: Recursive filter equivalent to a seismometer is applied to the waveform shown in *Upper-Right*, *Bottom-Right*: Inverse filter is applied to the waveform shown in *Bottom-Left* panel. Compare *Upper-Right* and *Bottom-Right* panels to recognize the effect of inverse filter. Note that the band pass filtering at the second step is necessary to prevent the instability of output due to the big amplification at a high frequency by *VSEISM.EXE* and at a low frequency by *ISEISM.EXE*.

**Exercise: Inverse Filter for a Simple Moving Coil Type Seismometer—II (Real Seismogram)**

Ch1.dat is a seismogram (NS-component) observed by the L4C seismometer, the parameters of which are  $\Delta t = 0.05$ ,  $N = 8192$ ,  $T_0 = 1.0$  s,  $h = 0.64$ .

- (1) Draw wave forms of ch1.dat using *PTIME.EXE*.
- (2) Run *ISEISM.EXE* with  $T_0 = 1.0$  s,  $h = 0.64$ ,  $\Delta t = 0.02$  and gain  $G_0 = 2.7$  (V/cm/s) (Input: ch1.dat, Output: UT0).
- (3) Apply a band pass filter by using *TRFILT.EXE* to this impulse. For example, the Butterworth type,  $F_L = 0.1$  Hz,  $F_H = 18.0$  Hz,  $F_S = 20.0$  Hz,  $A_P = 0.1$ ,  $A_S = 10.0$  (Input: CH1.dat, Output: UT).
- (4) Run *ISEISM.EXE* with  $T_0 = 1.0$  s,  $h = 0.64$ ,  $\Delta t = 0.02$  and gain  $G_0 = 2.7$  (V/cm/s) (Input: UT, Output: UT1).
- (5) Apply a band pass filter by using *TRFILT.EXE* with the Butterworth type,  $F_L = 0.1$  Hz,  $F_H = 18.0$  Hz,  $F_S = 20.0$  Hz,  $A_P = 0.1$ ,  $A_S = 10.0$  (Input: UT1, Output: UT2).

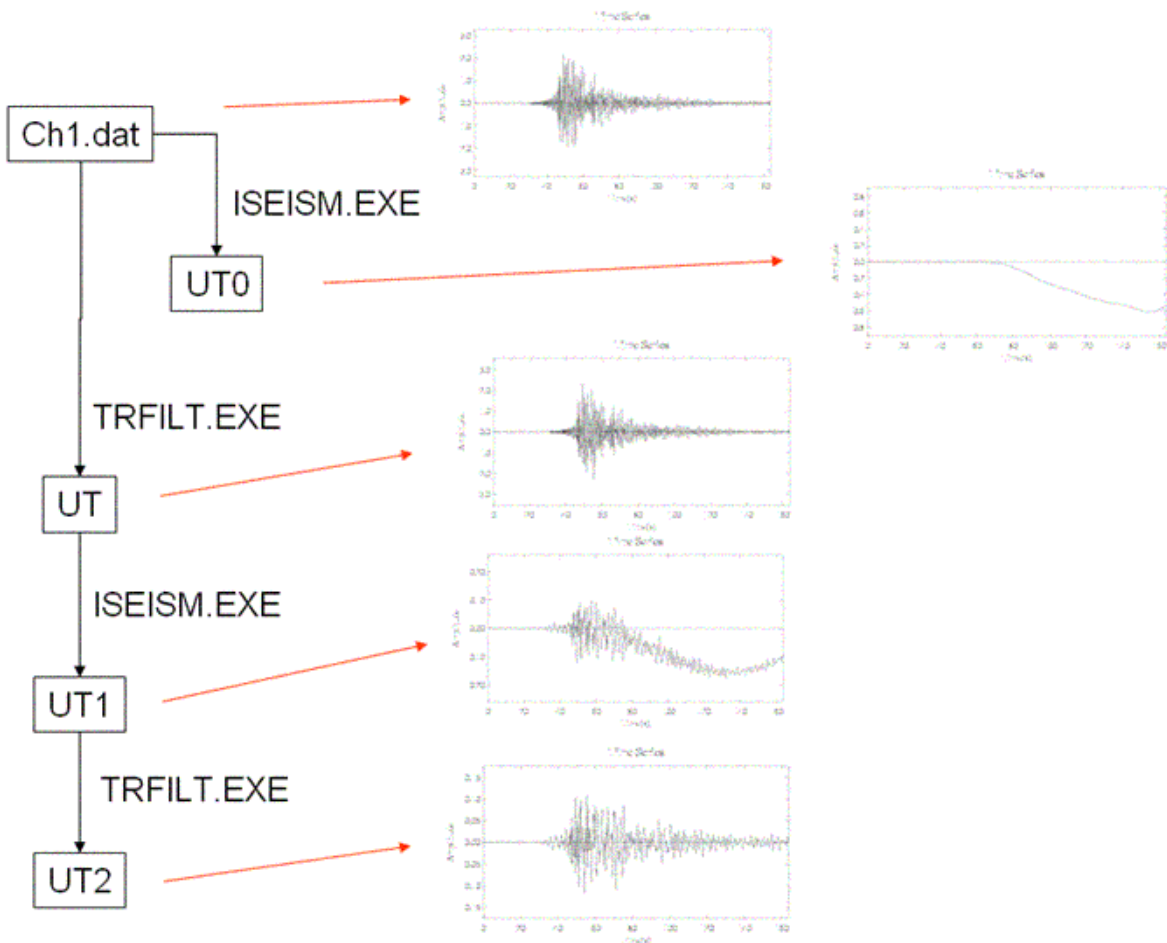


Fig. 53.2 Procedure and output of the example for a real seismogram.

**Example: Filter for Conversion of Seismogram obtained by a Simple Moving Coil Type Seismometer to that of Longer Period Seismometer.**

A simple way to avoid the abovementioned problem is to convert the system response with a longer-period seismometer.

Suppose that the damping constant and natural angular frequency of the seismometer used for the observations are  $h_0, \omega_0$ , then its response is

$$\frac{e_m}{y_m} = G_0 \cdot \frac{(i\omega)}{1 - 2ih_0(\omega_0/\omega) - (\omega_0/\omega)^2}.$$

Let us convert this with the response from a seismometer, in which the damping constant and natural angular frequency are  $h_1, \omega_1$ .

$$\frac{e_m}{y_m} = G_1 \cdot \frac{(i\omega)}{1 - 2ih_1(\omega_1/\omega) - (\omega_1/\omega)^2}.$$

The inverse function in the frequency domain is

$$G^{-1} \cdot \frac{1 - 2ih_0(\omega_0/\omega) - (\omega_0/\omega)^2}{1 - 2ih_1(\omega_1/\omega) - (\omega_1/\omega)^2}, \quad G = G_0/G_1.$$

The transfer function in the  $s$ -domain is given as follows:

$$T(s) = G^{-1} \cdot \frac{s^2 + 2h_0\omega_0s + \omega_0^2}{s^2 + 2h_1\omega_1s + \omega_1^2}. \quad (37)$$

Then, using the warped frequency gives

$$\begin{aligned} T(s) &= G^{-1} \cdot \frac{(1 - z^{-1})^2 + 2h_0 \tan\left(\frac{\omega_0\Delta t}{2}\right)(1 - z^{-1})(1 + z^{-1}) + \tan^2\left(\frac{\omega_0\Delta t}{2}\right)(1 + z^{-1})^2}{(1 - z^{-1})^2 + 2h_1 \tan\left(\frac{\omega_1\Delta t}{2}\right)(1 - z^{-1})(1 + z^{-1}) + \tan^2\left(\frac{\omega_1\Delta t}{2}\right)(1 + z^{-1})^2} \\ &= \frac{a_0 + a_1z^{-1} + a_2z^{-2}}{b_0 + b_1z^{-1} + b_2z^{-2}}, \end{aligned}$$

where

$$\begin{aligned} a_0 &= 1 + 2h_0 \tan(\omega_0\Delta t/2) + \tan^2(\omega_0\Delta t/2), \\ a_1 &= -2 + 2 \tan^2(\omega_0\Delta t/2), \\ a_2 &= 1 - 2h_0 \tan(\omega_0\Delta t/2) + \tan^2(\omega_0\Delta t/2), \\ b_0 &= G \{1 + 2h_1 \tan(\omega_1\Delta t/2) + \tan^2(\omega_1\Delta t/2)\}, \\ b_1 &= G \{-2 + 2 \tan^2(\omega_1\Delta t/2)\}, \\ b_2 &= G \{1 + 2h_1 \tan(\omega_1\Delta t/2) + \tan^2(\omega_1\Delta t/2)\}. \end{aligned}$$

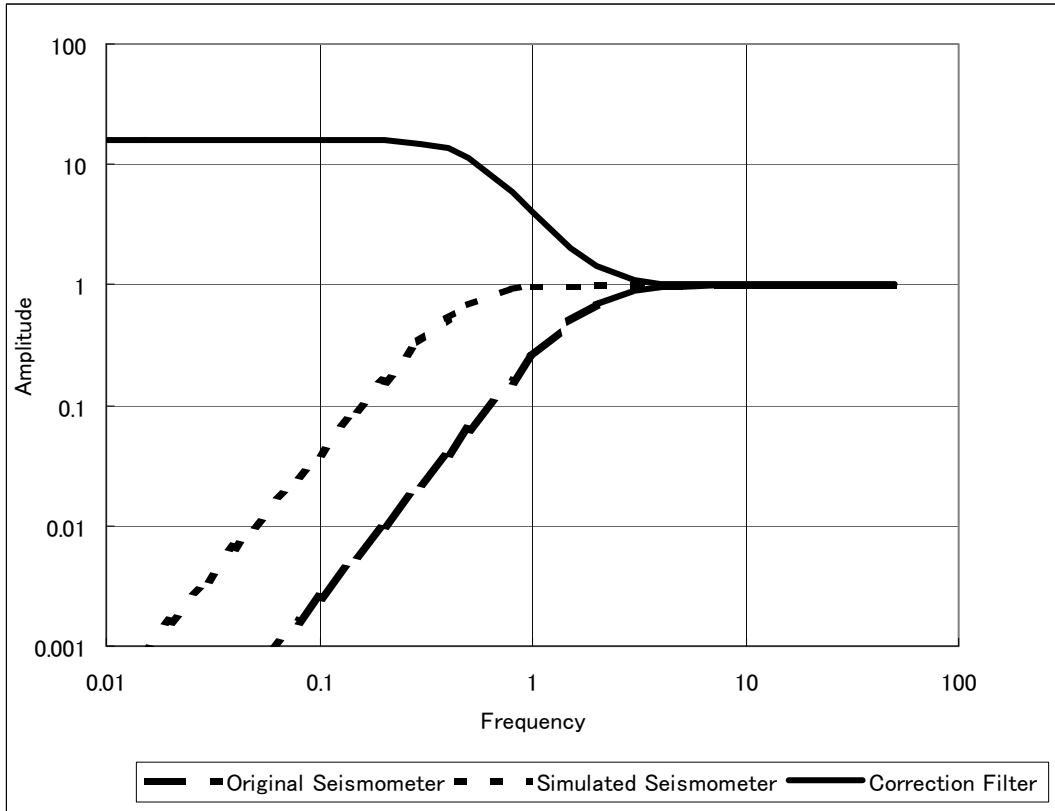


Fig. 54 Amplitude spectra of the correction filter explained above.

**Exercise: Filter for Conversion of Seismogram obtained by a Simple Moving Coil Type Seismometer to that of Longer Period Seismometer (An Impulse) .**

- (1) The program *CSEISM.EXE* calculates the coefficients of the recursive filter that can convert the characteristics of the simple moving coil type seismometer used originally for (1) observing those of the simulated seismometer with a longer period.
- (2) Prepare an input time series by using *TESTSIG.EXE*. For example, an impulse of unit amplitude located at  $t = 0.0$  s, with  $\Delta t = 0.05$  s,  $N = 128$  (Output: UT0)
- (3) Calculate Fourier transform by using *FFT.EXE* (Input: UT0, Output: UF0).
- (4) Run *DSEISM.EXE* with the natural period  $T_0 = 0.25$  s, the damping factor  $h_0 = 0.71$ ,  $\Delta t = 0.05$  and the gain  $G_0 = 1.0$  (Input: UT0, Output: UT1).  
Calculate Fourier transform by using *FFT.EXE* (Input: UT1, Output: UF1).
- (5) Run *CSEISM.EXE* with  $T_0 = 0.25$  s,  $h_0 = 0.71$ ,  $T_1 = 1.0$  s,  $h_1 = 0.71$ ,  $\Delta t = 0.05$ , and gain  $G_0 = 1.0$  (Input: UT1, Output: UT2).  
Calculate Fourier transform by using *FFT.EXE* (Input: UT2, Output: UF2).

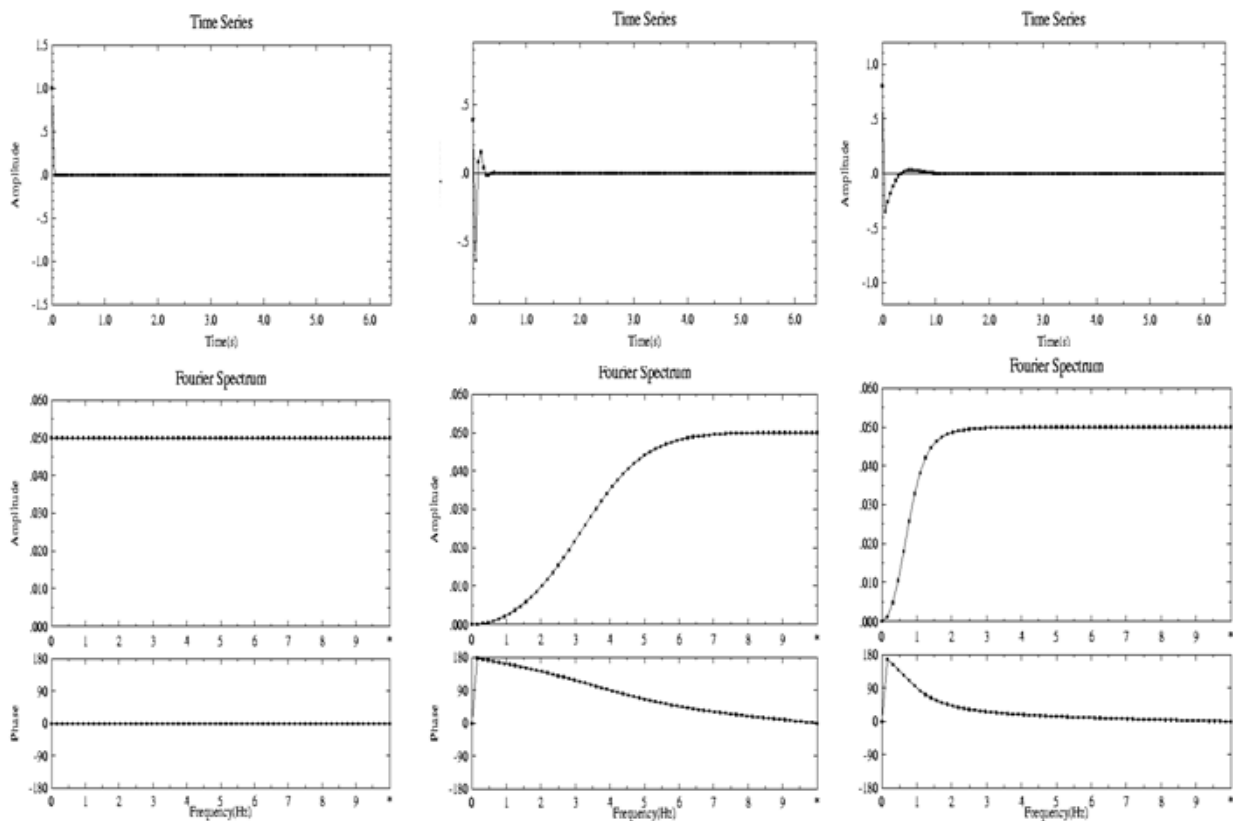


Fig. 55.1 Waveforms and spectra of the example mentioned above. *Left*: Impulse used as the input. *Center*: Simulated output for a seismometer  $T_0 = 0.25$  s,  $h_0 = 0.71$  by *DSEISM.EXE*. *Right*: Converted output from  $T_0 = 0.25$ ,  $h_0 = 0.71$  s to  $T_0 = 1.0$ ,  $h_0 = 0.71$  s by *CSEISMO.EXE*.



**Exercise: Filter for Conversion of Seismogram obtained by a Simple Moving Coil Type Seismometer to that of Longer Period Seismometer (Real Seismogram)**

Ch1.dat is a seismogram (NS-component) observed by L4C seismometer, the parameters of which are  $\Delta t = 0.05$ ,  $N = 8192$ ,  $T_0 = 1.0$  s,  $h = 0.64$ .

(1) Draw wave forms of ch1.dat using *PTIME.EXE*.

(2) Run *CSEISM.EXE* with

Original Seismometer

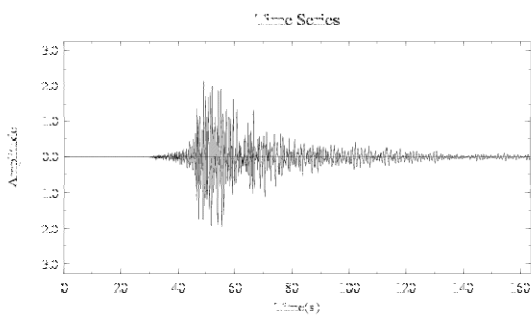
$T_0, h, dt, g_0$ : 1.0, 0.64, 0.02, and 2.7 (V/cm/s)

Simulated Seismometer

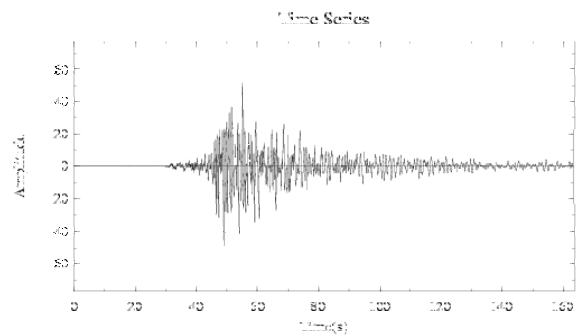
$T_0, h, dt, g_0$ : 5.0, 0.71, 0.02, and 50.0 (V/cm/s)

(Input: ch1.dat, Output: output.dat).

Calculate Fourier transform by using *FFT.EXE* (Input: output.dat, Output:output.spc).



Input.dat



Output.dat

Fig. 55.2 Input signal (*left*): Ch1.dat and output signal (*right*) of *CSEISMO.EXE*

## 3.6. Integration of accelerogram or Base Line Correction

### 3.6.1 Origin of Low-Frequency Errors

Chiu (1997) explains the origins of low-frequency errors in the following manner. We are interested mainly in low-frequency noises because the effects of high-frequency noises are considerably reduced by double integration since their frequency band is lower than Nyquist frequency.

- (1) Deviations of the instrument response from flat amplitude response and from linear phase shift. These are large in a high-frequency range. Careful calibration by the manufacturer is essential.
- (2) Limited resolution of analog-to-digital Converter. This causes the approximations and the drift.
- (3) Insufficient sampling rate that can be prevented by considering the Nyquist frequency. Although the frequency band of interest to engineering seismologists is lower than 10 Hz, sampling with at least 100 Hz is recommended in order to obtain the value of the peak acceleration correctly.
- (4) Electronic Noise: besides the AC noise at 50 or 60 Hz, low-frequency electric noise caused by fluctuation of temperature is present.
- (5) Ambient noises that include a long-period microtremor and slow deformation of observatory building, *e.g.*, due to sunshine affects the integrated displacement seismogram significantly.

Their effects appear superposed over each other and the seismic signals in the seismograms and in the output of digital data processing. It is almost impossible to separate the signal from these noises in the time domain.

These cause the base line drift of accelerograms and absurd waveforms of velocity and displacement seismograms obtained by numerical integration. An example is shown below.

Fig. 56.1 shows the original accelerograms obtained at K-NET maintained by NIED, Japan; its station is KGS005 in the northwest part of Kagoshima pref. and the seismic event ( $M_{JMA} = 3.6$ ) took place 21:26 on March 27, 1997. This is one of the aftershocks of Kagoshima-Ken Hokusei-bu Eq. ( $M = 6.3$ ). These include the noises mentioned above, although there is no visible noise in the shown scale.

The acceleration amplitude spectra of the seismic event (time window; [12.0 s, 24.0 s]) shown in the top panels of Fig. 56.2 has a wide peak that corresponds to the seismic signals and a portion of noise that is more clearly shown in the low frequency band. The latter is emphasized in the displacement amplitude spectra shown in the bottom panels. Fig. 56.3 shows the amplitude spectra of the noise (time window; [0.0 s, 10.0 s]) recorded before the P-wave arrivals. A comparison of Fig. 56.3 with Fig. 56.2 shows the effects of the low-frequency noises in the frequency domain.

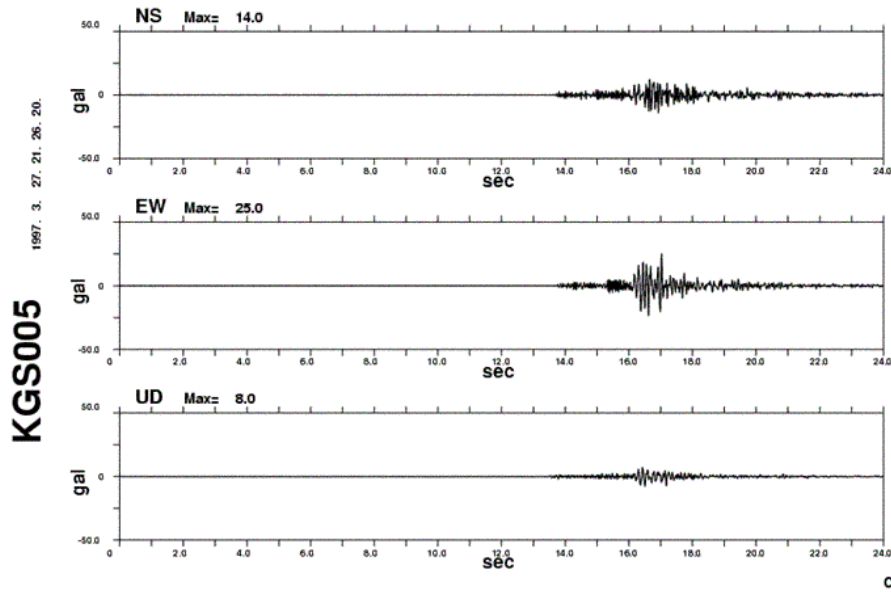


Fig. 56.1 Original Accelerograms. Plotted by *pltacc.exe*.

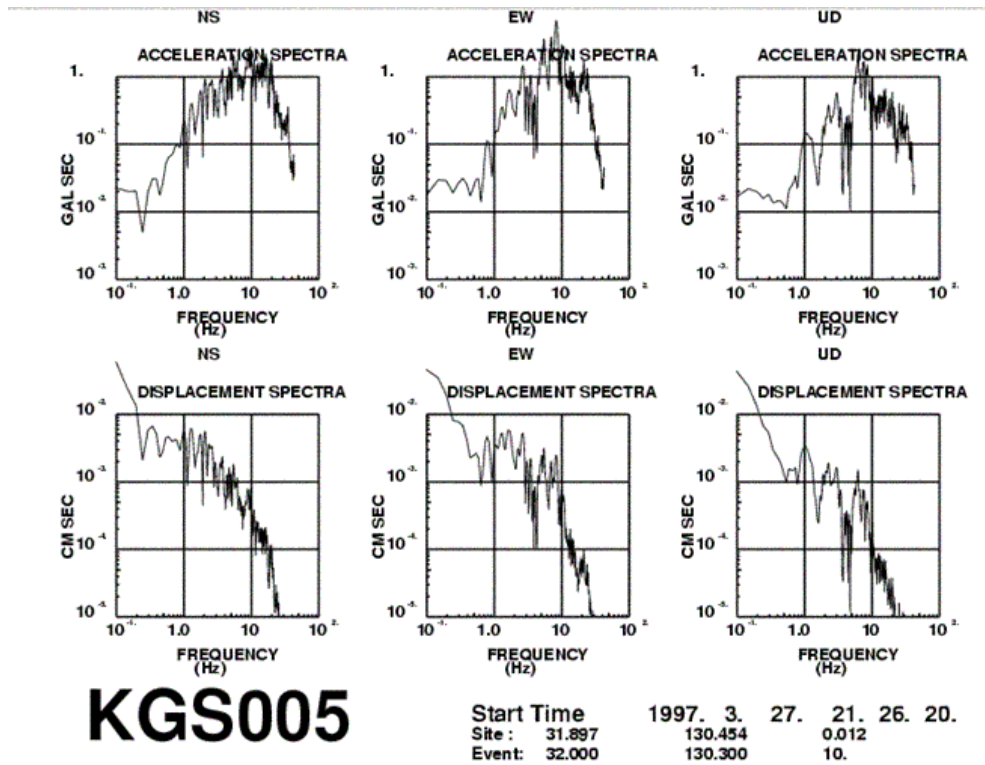


Fig. 56.2 Amplitude spectra of accelerograms shown in Fig. 56.1 for the time window [12.0 s, 24.0 s] that include the seismic event. The increasing peak at the frequency band lower than 0.4Hz is emphasized in the displacement amplitude spectra obtained by the division of acceleration spectra by  $(2\pi f)^2$ . Plotted by *splot.exe*.

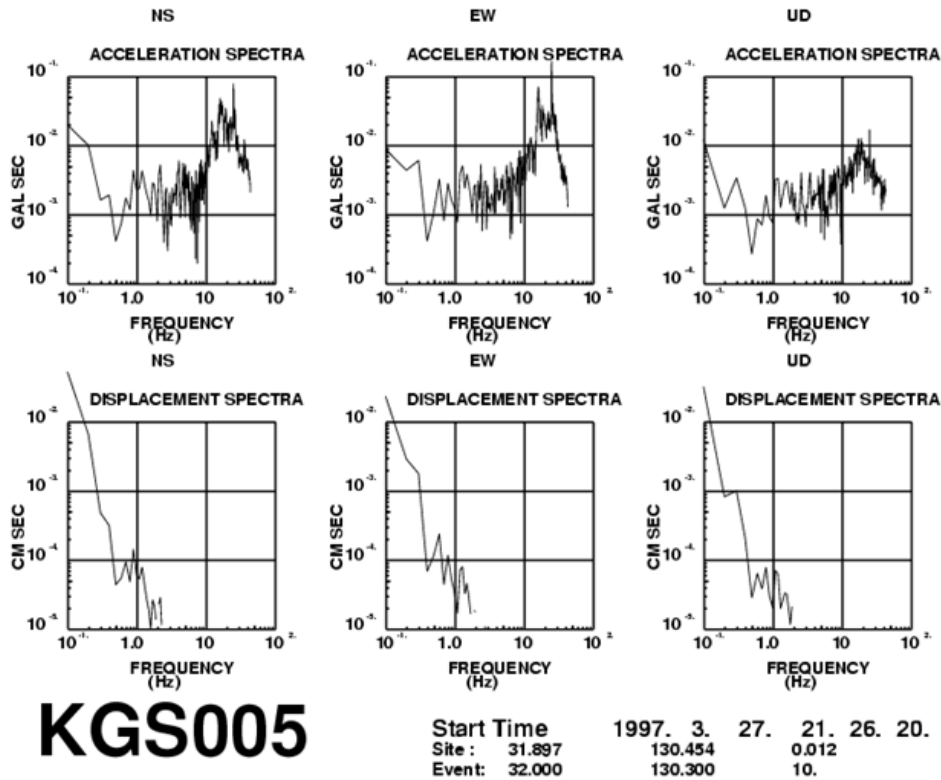


Fig. 56.3 Amplitude spectra of accelerograms shown in Fig. 56.1 for the time window [0.0 s, 10.0 s] that is the noise part. The increasing peak at the frequency band lower than 0.4 Hz is emphasized in the displacement amplitude spectra obtained by the division of acceleration spectra by  $(2\pi f)^2$ . Plotted by *splot.exe*.

The integration of the time series can be performed by recursive filtering as given by the following.

$$y_m = y_{m-1} + \Delta t \cdot x_m.$$

The velocity seismograms obtained by a direct integration of the accelerograms shown in Fig. 56.1 are shown in Fig. 56.4. These almost linearly increasing velocity seismograms show the effect of the constant offset included in the original accelerograms. The slightly observable curvature is due to the linear trend included in them.

The correction for the constant offset and linear trend in accelerograms is performed by the least square fitting as shown by Eq. (8). Then, the numerical integration is performed. The result is shown in Fig. 56.5. This is considerably improved in comparison with Fig. 56.4. The slight base line drift of the long period is supposed to be due to the influence of the low-frequency noise remaining in the corrected accelerograms.

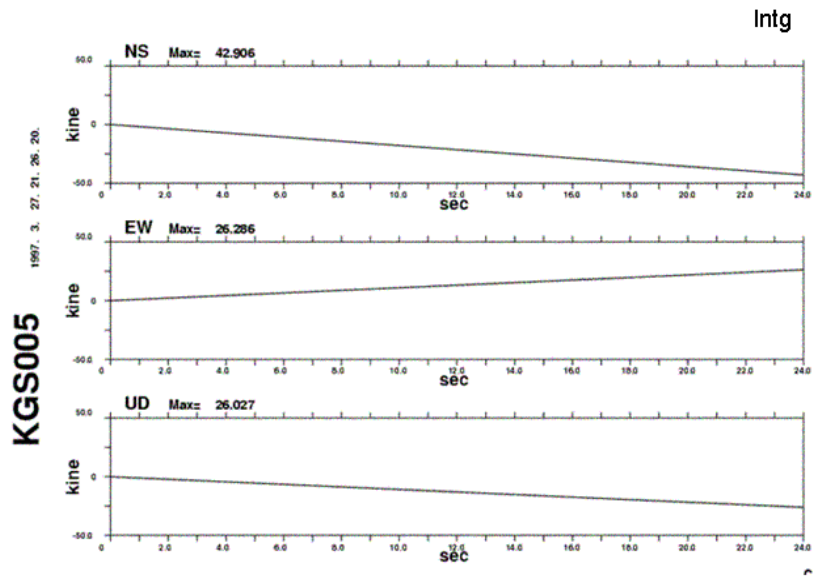


Fig. 56.4 Velocity seismograms obtained by a direct integration of the accelerograms shown in Fig. 56.1. The apparent linear trends are due to the influence of the constant drift in the accelerograms. Calculated by *acc2vel.exe* with comments for the offset-trend correction and band pass filtering, and then plotted by *pltvel.exe*.

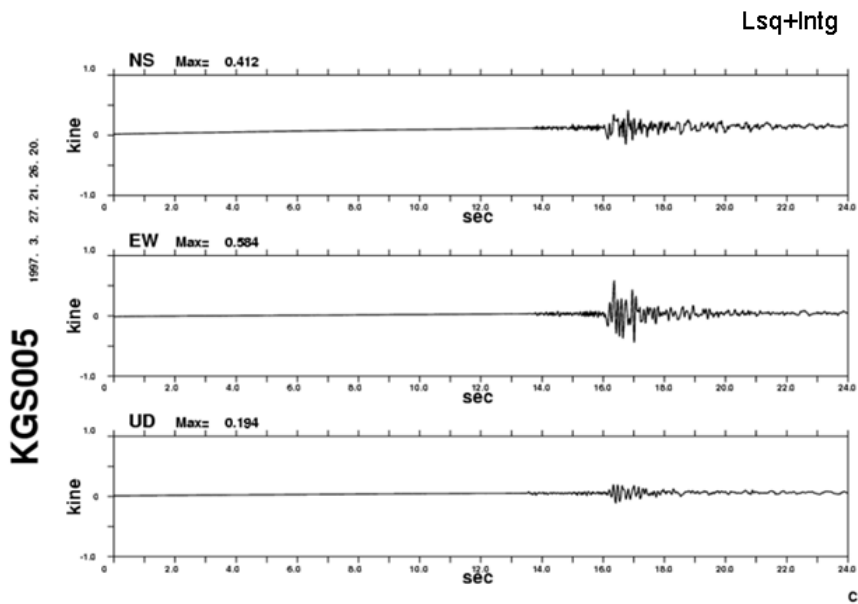


Fig. 56.5 The accelerograms shown in Fig. 56.1 are corrected by the least square fitting method and then integrated numerically. The slight drift in the base line is due to the influence of the low-frequency noise remaining in the corrected accelerograms. Calculated by *acc2vel.exe* with comments due to band pass filtering, and then plotted by *pltvel.exe*.

The remaining low-frequency noises can be removed by applying the high pass or band pass filter aimed to cut the low frequency component.

The example is shown in Fig. 56.6. The used values for the parameters of the band pass filter are as follows:.

fl, fh, fs/0.7, 25.0, 30.0/ ntype,nchara,ncausal/1,3,2/

$f_1 = 0.7$  Hz corresponds to the limit frequency in which the seismic signals become dominant in comparison with noises (Fig. 56.2 *top panels*).  $f_h = 25.0$  Hz and  $f_s = 30.0$  Hz are selected in order to ensure the suppression of high-frequency noises. Ntype = 1 means Butterworth filter and nchara = 3 means band pass feature. ncause= 2 that means that a zero-phase filter is selected in order to maintain the waveform of seismic signals.

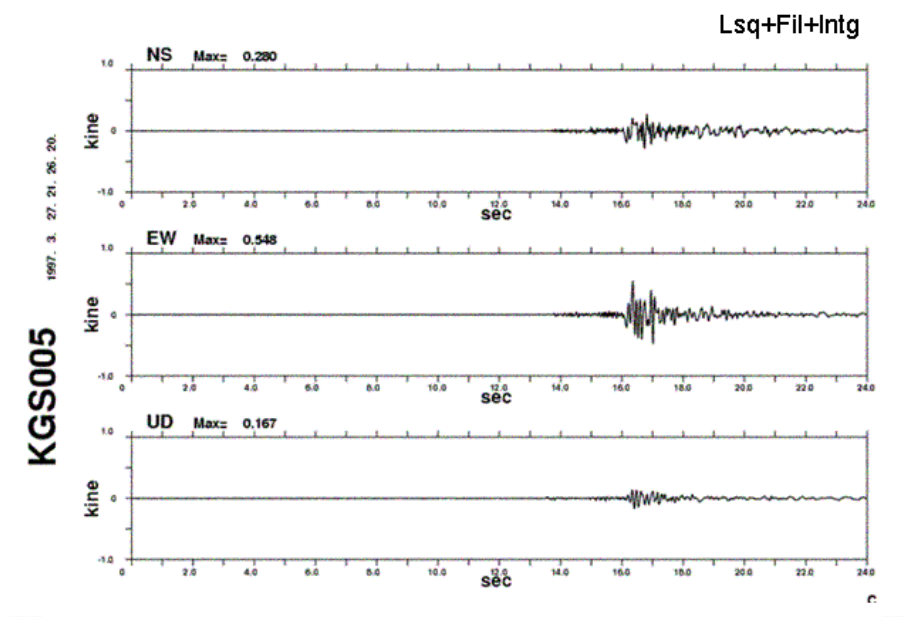


Fig. 56.6 The accelerograms shown in Fig. 56.1 are corrected by the least square fitting method, band pass filtered, and then integrated numerically. They were calculated by *acc2vel.exe* and plotted by *pltvel.exe*.

The conversion of the accelerogram to the displacement seismogram is tougher than that in the case of a velocity seismogram because of the double integrations. Fig. 56.7 shows the result of a direct integration of the velocity seismograms shown in Fig. 56.4, *i. e.*, those of the double direct integration of the accelerograms shown in Fig. 56.1. The parabolic feature is due to the integration of the almost linear velocity seismograms (Fig. 56.4). Note the absurdly big values of the maximum value.

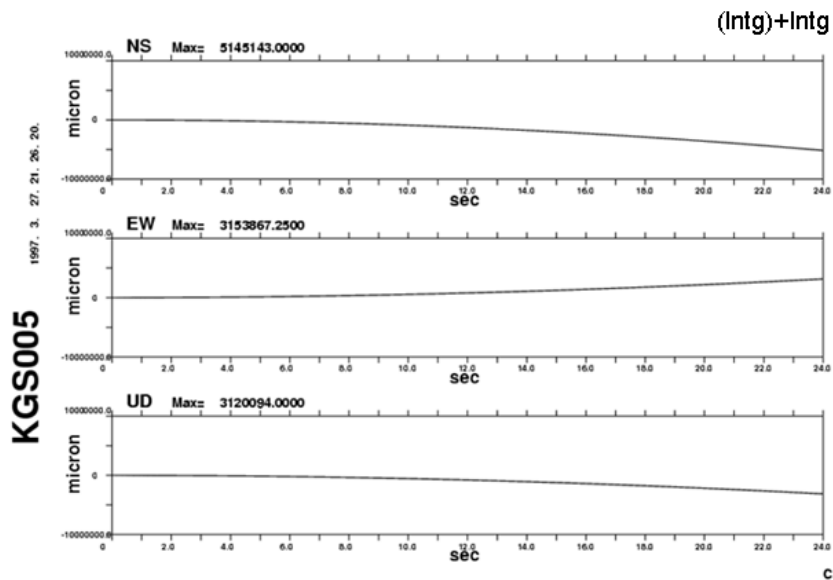


Fig. 56.7 Displacement seismogram obtained by a direct integration of the velocity seismograms shown in Fig. 56.4. Calculated by *acc2dis.exe* with comments on the offset-drift correction and band pass filtering for accelerograms and velocity seismograms; subsequently plotted by *pltdis.exe*.

Fig. 56.8 shows the displacement seismogram obtained by the least square correction and the integration of the velocity seismograms shown in Fig. 56.5. These show parabolic waveforms that are integrated low-frequency noises as well.

Fig. 56.9 Displacement seismogram obtained by a direct integration of the velocity seismograms shown in Fig. 56.6.

Fig. 56.10 Displacement seismogram obtained by the least square correction of the velocity seismograms shown in Fig. 56.6. This give a better performance than the others, but not sufficient. A slight trend is obvious in the NS component.

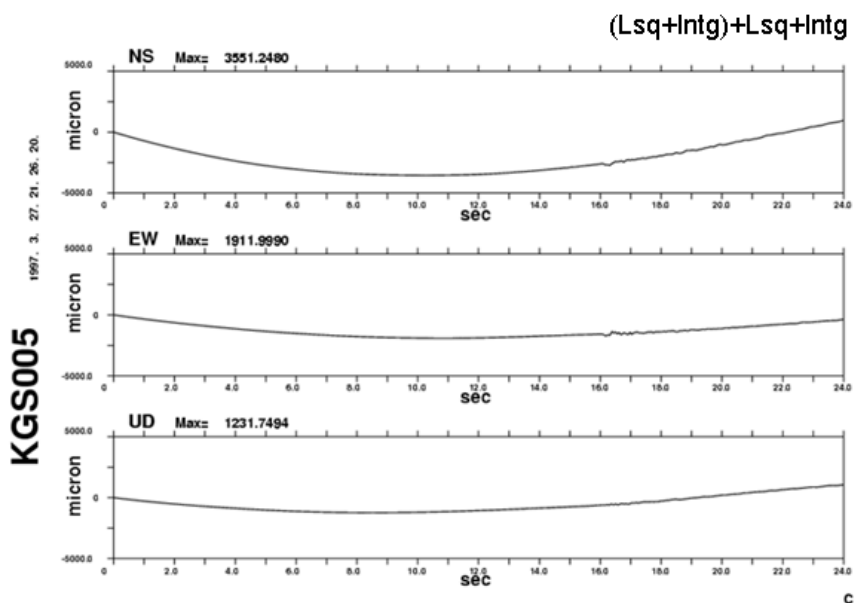


Fig. 56.8 Displacement seismogram obtained by the least square correction and integration of the velocity seismograms shown in Fig. 56.5. Calculated by *acc2dis.exe* with comments from the band pass filtering for accelerograms and velocity seismograms; subsequently plotted by *pltdis.exe*.



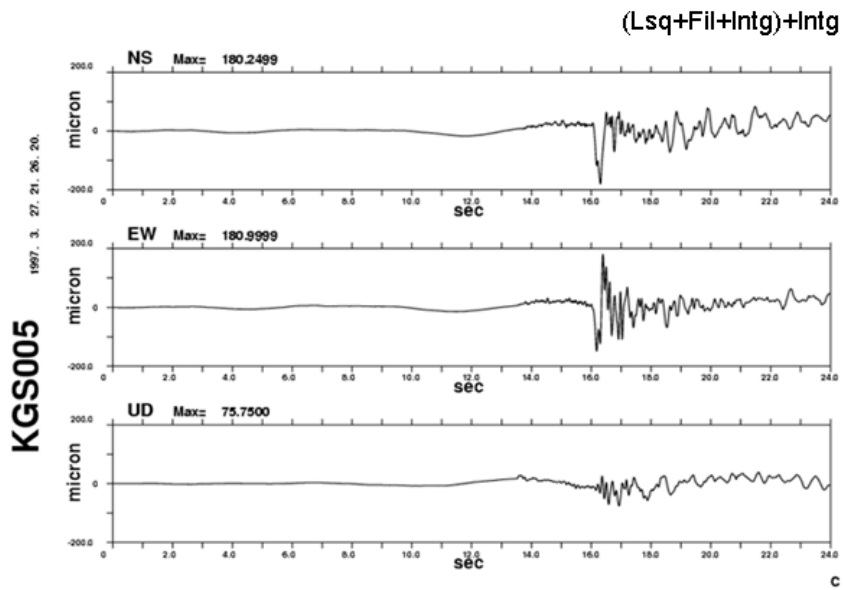


Fig. 56.9 Displacement seismogram obtained by a direct integration of the velocity seismograms shown in Fig. 56.6. Calculated by *acc2dis.exe* with comments from the offset-drift correction and band pass filtering for velocity seismograms; subsequently plotted by *pltdis.exe*.

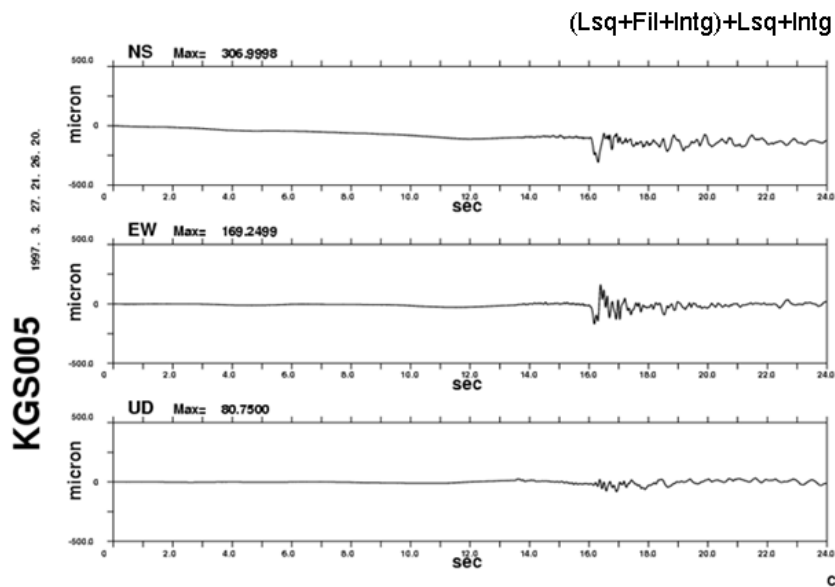


Fig. 56.10 Displacement seismogram obtained by the least square correction of the velocity seismograms shown in Fig. 56.6. Calculated by *acc2dis.exe* with comments from band pass filtering for velocity seismograms; subsequently plotted by *pltdis.exe*.

The attempts described above imply the necessity of an offset-drift correction and band pass filtering for accelerograms and velocity seismograms. Fig. 57.1 and Fig. 57.2 show stable results, i. e., displacement seismograms without a base line drift.

The difference between Fig. 57.1 and Fig. 57.2 is the value of the parameter  $f_i$ . In Fig. 57.1,  $f_i = 0.1$  Hz, and in Fig. 57.2  $f_i = 0.7$  Hz. A long-period ripple is observed before P-arrival in Fig. 57.1 due to the low-frequency noises that can be suppressed by selecting  $f_i = 0.7$  Hz

The base line correction of the velocity and displacement seismograms integrated from observed an accelerogram is a very basic task for strong motion observation. The techniques of digital data processing described in this lecture note are necessary in totality.

The recommended way of integration is as follows.

- (1) Calculate and plot the Fourier spectra of the observed accelerogram for both noise part and seismic event part.
- (2) Determine the value of  $f_i$  that is the low frequency limit of the band such that the seismic signal is dominant in comparison with the noise level.
- (3) Select the values of  $f_h$  and  $f_s$ , if the band pass filter is used. If a high pass filter is used, they are not necessary.
- (4) Apply the offset-trend correction by the least square method and filtering to the observed accelerograms.
- (5) Integrate the processed accelerograms; then the velocity seismograms are then obtained.
- (6) Apply the offset-trend correction by the least square method and filtering to the obtained velocity seismograms.
- (7) Integrate the processed velocity seismograms; then the displacement seismograms are obtained.
- (8) If the obtained displacement seismograms have a long-period deviation before P-arrival time, return to (2) and change the value of  $f_i$ .

The result can be affected by the types of filter applied. Try to check the performance of various filters on your own.

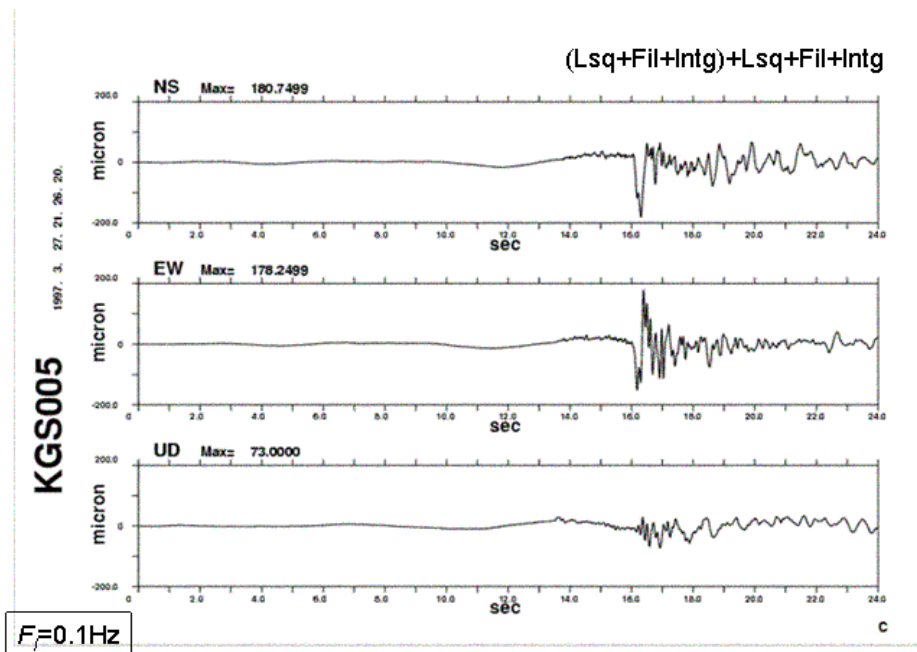


Fig. 57.1 Displacement seismogram obtained by the least square correction and band pass filtering of the velocity seismograms shown in Fig. 56.6. Long-period ripple observed before P-arrival is due to the low-frequency noises.

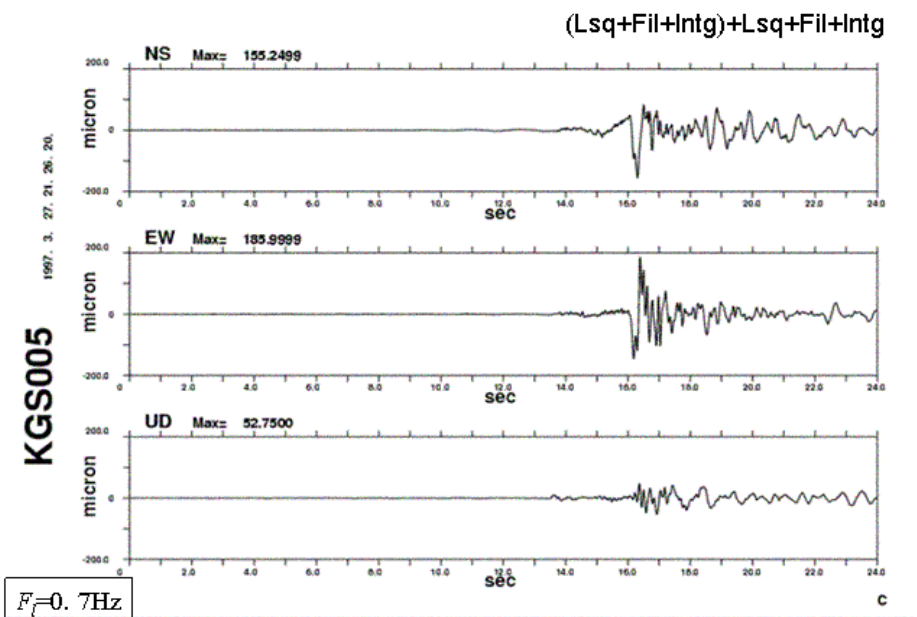


Fig. 57.2 Displacement seismogram obtained by the least square correction and band pass filtering of the velocity seismograms shown in Fig. 56.6. Long-period ripple observed before P-arrival in Fig. 57.1 is suppressed sufficiently.

## Fortran Programs

### CFFT.FOR

```
c*****
SUBROUTINE CFFT(X,N,IND)
c.....
c Complex Fast Fourier Transform
c Coded by Y. OHSAKI in "Introduction to spectral analysis for
c                               seismic waves"
c Typed by D.Suetsugu, IISEE
c Modified by H.Inoue, IISEE
c Usage.....   CALL FFT(N,X,IND)
c where
c X(N) In/Out Equi-spaced N complex data/Transformed values
c N   In   Number of data, MUST be a power of 2.
c IND In   -1=Fourier transform,  1=Inverse Fourier transform
c
c Note..... You must divide the output X() by N to get Fourier
c           transform when IND=-1.
c.....
complex x(n),temp,theta
c
  j=1
  do 140 i=1,n
    if(i.ge.j) goto 110
    temp=x(j)
    x(j)=x(i)
    x(i)=temp
110   m=n/2
120   if(j.le.m) goto 130
        j=j-m
        m=m/2
        if(m.ge.2) goto 120
130   j=j+m
140 continue
c
  kmax=1
  do 150 kkk=1,9999
    if(kmax.ge.n) return
    istep=kmax*2
    do 170 k=1,kmax
      theta=cmlpx(0.0,3.141593*float(ind*(k-1))/float(kmax))
      do 160 i=k,n,istep
        j=i+kmax
        temp=x(j)*cexp(theta)
        x(j)=x(i)-temp
        x(i)=x(i)+temp
160      continue
170      continue
      kmax=istep
150 continue
end
```

### TESTSIG.FOR

```
c*****
c Make a test signal u(t)
c*****
program main
parameter (nmax=8192,pi=3.1415926536)
real t(nmax),u(nmax)
character yn*1,ofile*80
c
if(nargs().le.1) then
  write(*,*)
  write(*,*) 'Usage: TESTSIG outfile'
  stop ''
end if
call getarg(1,ofile,istatus)
c ..... Input parameters
write(*,*)
write(*,*) 'Parameters for a test signal.'
write(*,*(a$)) ' Sampling interval dT (sec) >>'
read(*,*) dt
write(*,*(a$)) ' Number of data points N >>'
read(*,*) ndata
write(*,*(1x,a,f10.3,a)) 'Total time is ,dt*ndata, sec.'
write(*,*(1x,72a)) ('-',i=1,72)
do i=1,ndata
  t(i)=(i-1)*dt
  u(i)=0.0
end do
c ..... Superposition loop
do isum=1,999
  write(*,*(a$)) ' Cos/Sin=1, Const=2, Impulse=3,/'
  & ' Step=4, Triangle=5, Square=6 >>'
  read(*,*) ichoice
  if(ichoice.eq.1) then
    write(*,*(a$)) ' Period of the cosine wave (sec)>>'
```

```
read(*,*) period
write(*,*(a$)) ' Amplitude >>'
read(*,*) amp
write(*,*(a$)) ' Phase (0-360 in degree) >>'
read(*,*) phase
write(*,*(a$)) ' Damping factor >>'
read(*,*) damp
do i=1,ndata
  arg=phase/180.*pi+2*pi/period*t(i)
  u(i)=u(i)+amp*cos(arg)*exp(-damp*t(i))
end do
else if(ichoice.eq.2) then
  write(*,*(a$)) ' The constant value >>'
  read(*,*) const
  do i=1,ndata
    u(i)=u(i)+const
  end do
else if(ichoice.eq.3) then
  write(*,*(a$)) ' Location of the impulse (sec) >>'
  read(*,*) timpulse
  write(*,*(a$)) ' Amplitude of the impulse >>'
  read(*,*) aimpulse
  do i=1,ndata
    if(t(i).ge.timpulse) then
      u(i)=u(i)+aimpulse
      exit
    end if
  end do
else if(ichoice.eq.4) then
  write(*,*(a$)) ' Location of the step(sec) >>'
  read(*,*) tstep
  write(*,*(a$)) ' Amplitude of the step >>'
  read(*,*) astep
  do i=1,ndata
    if(t(i).ge.tstep) then
      u(i)=u(i)+astep
    end if
  end do
else if(ichoice.eq.5) then
  write(*,*(a$)) ' Period of the triangular wave(sec)>>'
  read(*,*) period
  write(*,*(a$)) ' Amplitude >>'
  read(*,*) amp
  write(*,*(a$)) ' Phase (0-360 in degree) >>'
  read(*,*) phase
  do i=1,ndata
    half=period/2.
    tt=t(i)+phase/360.*period
    tmod=mod(tt,half)
    if(mod(int(tt/half),2).eq.0) then
      u(i)=u(i)+amp-amp*2*tmod/half
    else
      u(i)=u(i)-amp+amp*2*tmod/half
    end if
  end do
else if(ichoice.eq.6) then
  write(*,*(a$)) ' Period of the square wave(sec) >>'
  read(*,*) period
  write(*,*(a$)) ' Amplitude >>'
  read(*,*) amp
  write(*,*(a$)) ' Phase (0-360 in degree) >>'
  read(*,*) phase
  do i=1,ndata
    half=period/2
    tt=t(i)+phase/360*period+period/4.0
    if(mod(int(tt/half),2).eq.0) then
      u(i)=u(i)+amp
    else
      u(i)=u(i)-amp
    end if
  end do
end if
write(*,*(a$)) ' Add more? (y/n) >>'
read(*,*(a)) yn
if(yn.ne.'y'.and.yn.ne.'Y') exit
end do
write(*,*(1x,72a)) ('-',i=1,72)
c ..... Output Time Series
open(10,file=ofile)
write(10,*) ndata,'=N'
write(10,*) dt,'=dT (sec)'
do i=1,ndata
  write(10,*) u(i)
end do
close(10)
write(*,*) 'Done'
c
end
```

```

PTIME.FOR
c*****
c      Plot time series in a file
c*****
      program main
      parameter (nmax=8192)
      real t(nmax),u(nmax)
      character ifile*80
c ..... Input Time Series
      if(nargs().le.1) then
         write(*,*)
         write(*,*) 'Usage: PTIME infile'
         stop ''
      end if
      call getarg(1,ifile,istatus)
      open(10,file=ifile,status='old')
      read(10,*) ndata
      read(10,*) dt
      do i=1,ndata
         read(10,*) u(i)
         t(i)=dt*(i-1)
      end do
      close(10)
c ..... Plot Time Series
      umax=u(1)
      umin=u(1)
      do i=2,ndata
         if(u(i).gt.umax) umax=u(i)
         if(u(i).lt.umin) umin=u(i)
      end do
      ymax=max(abs(umax),abs(umin))*1.5
c
      call opengraphics
      call setwindow(1,0,-ymax,dt*ndata,ymax)
      call setviewpt(1,0.3,0.3,1.1,0.7)
      call drawtitle('Time Series')
cc      call setline('width=2')
      call setttext('size=0.8')
      call drawaxis('labelx=Time(s),labely=Amplitude')
      call drawline(0,0,dt*ndata,0)
      call setmark('type=solid-square,size=0.5')
      call drawlines(ndata,t,u)
      call drawmarks(ndata,t,u)
      call closegraphics
      end

FFT.FOR
      include 'cfft.for'
c*****
c      Calculate Fourier Coefficients by FFT
c*****
      parameter(nmax=8192)
      real u(nmax)
      complex cf(nmax)
      character ifile*80,ofile*80
c
      if(nargs().le.2) then
         write(*,*)
         write(*,*) 'Usage: FFT infile outfile'
         stop ''
      end if
      call getarg(1,ifile,istatus)
      call getarg(2,ofile,istatus)
c ..... Read time series from file
      open(10,file=ifile,status='old')
      read(10,*) ndata
      read(10,*) dt
      do i=1,ndata
         read(10,*) u(i)
      end do
      close(10)
c ..... Check if ndata=2**n
      npow=nint(log(float(ndata))/log(2.0))
      nn=2**npow
      if(ndata.ne.nn) then
         write(*,*) 'N=',ndata,', Error: N must be a power of 2.'
stop
      end if
c ..... FFT
      do i=1,nn
         cf(i)=cmplx(u(i),0.0)
      end do
      call cfft(cf,nn,-1)
      do i=1,nn
         cf(i)=cf(i)/nn
      end do
c ..... Write Fourier Coeff. to file
      df=1.0/(dt*nn)
      open(20,file=ofile)
      write(20,*) nn,'=N'

```

```

      write(20,*) df,'=df'
      do i=1,nn
         write(20,*) cf(i)
      end do
      close(20)
      write(*,*) 'Done'
c .....
      end

PCFFT.FOR
c*****
c      Plot FFT Coefficients
c*****
      character*20 ifile,ofile
      parameter(nmax=8192,pi=3.1415926535,d2r=pi/180.)
      complex cf(nmax)
      real amp(nmax/2),phase(nmax/2),freq(nmax/2)
c ..... Input from file
      if(nargs().le.1) then
         write(*,*)
         write(*,*) 'Usage: PCFFT infile'
         stop ''
      end if
      call getarg(1,ifile,istatus)
      open(10,file=ifile,status='old')
      read(10,*) nn
      read(10,*) df
      do i=1,nn
         read(10,*) cf(i)
      end do
      close(10)
      fmax=nn/2*df
c ..... Plot spectrum
      i=0
      do k=nn/2+1,nn-1
         i=i+1
         freq(i)=(k-nn)*df
         amp(i)=cabs(cf(k+1))
         if(amp(i).ne.0.0) then
            phase(i)=atan2(imag(cf(k+1)),real(cf(k+1)))/d2r
         end if
      end do
      do k=0,nn/2
         i=i+1
         freq(i)=df*k
         amp(i)=cabs(cf(k+1))
         if(amp(i).ne.0.0) then
            phase(i)=atan2(imag(cf(k+1)),real(cf(k+1)))/d2r
         end if
      end do
c
      ampmax=amp(1)
      do i=2,nn
         if(amp(i).gt.ampmax) ampmax=amp(i)
      end do
      do i=1,nn
         if(amp(i)/ampmax.lt.1.e-5) then
            phase(i)=0.0
         end if
      end do
c ..... Graphics
      call opengraphics
      call setwindow(1,-fmax,0.0,fmax,ampmax*1.2)
      call setviewpt(1,0.3,0.45,1.1,0.8)
      call drawtitle('FFT Coefficients')
      call setttext('size=0.8')
      call drawaxis('labely=Amplitude')
      call setwindow(2,-fmax,-180.,fmax,180.)
      call setviewpt(2,0.3,0.2,1.1,0.40)
      call drawaxis('labelx=Frequency(Hz),labely=Phase,/'
& 'dy=90,ddy=90')
c ..... Plot amplitude and phase spectrum
      call setmark('type=solid-square,size=0.7')
      call selectwindow(1)
      call drawlines(nn,freq,amp)
      call drawmarks(nn,freq,amp)
      call selectwindow(2)
      call drawlines(nn,freq,phase)
      call drawmarks(nn,freq,phase)
      call closegraphics
c .....
      end

PSPEC.FOR
c*****
c      Plot Fourier Spectrum
c*****
      character*20 ifile,ofile

```



```

do i=nn/2+1,nn
  cf(i)=conjg(cf(nn-i+1))
enddo
c..... Write Fourier Coeff. to file
  open(20,file=ofile)
  write(20,*) nn, ' =N'
  write(20,*) df, ' =df'
  do i=1,nn
    write(20,*) cf(i)
  end do
  close(20)
  write(6,*) 'It is assumed that the phase spectra are zero.'
  write(6,*) 'Done'
c.....
end

```

### FPRDCT.FOR

```

c*****
c Product of filter spectra and input time series spectra
c*****
character*20 ifile,ofile
parameter(nmax=8192)
complex cf1(nmax),cf2(nmax)
c.....Input from file
10 write(6,*)'?? INPUT FILE NAME (Filter Spectra) ???'
  read(5,'(a)') ifile
  open(10,file=ifile,status='old')
  read(10,*) nn
  read(10,*) df
  do i=1,nn
    read(10,*) cf1(i)
  end do
  close(10)
  write(6,*)'?? INPUT FILE NAME (Input Signal Spectra) ???'
  read(5,'(a)') ifile
  open(10,file=ifile,status='old')
  read(10,*) nn1
  read(10,*) df1
  do i=1,nn
    read(10,*) cf2(i)
  end do
  close(10)
  if(nn1.ne.nn.or.df1.ne.df) then
    write(6,*)'Data number or frequency interval unmatched.'
    goto 10
  endif
  fmax=nn/2*df
  td=1/df
c..... Product
  do i=1,nn/2
    cf1(i)=cf1(i)*cf2(i)*td
  enddo
  do i=nn/2+1,nn
    cf1(i)=conjg(cf1(nn-i+1))
  enddo
c..... Write Fourier Coeff. to file
  write(6,*)'?? OUTPUT FILE NAME ???'
  read(5,'(a)') ofile
  open(20,file=ofile)
  write(20,*) nn, ' =N'
  write(20,*) df, ' =df'
  do i=1,nn
    write(20,*) cf1(i)
  end do
  close(20)
  write(6,*)'Done'
end

```

### FWVLET.FOR

```

c*****
c Make filter wavelet of a filter from the inverse
c Fourier Transform of given amplitude spectra
c*****
program main
parameter (nmax=8192)
real u(nmax),wt(nmax)
character ifile*80,ofile*80
c..... Input Time Series
write(6,*)'?? INPUT FILE NAME'
read(5,'(a)') ifile
write(6,*)'?? OUTPUT FILE NAME'
read(5,'(a)') ofile
open(10,file=ifile,status='old')
read(10,*) ndata
read(10,*) dt
do i=1,ndata
  read(10,*) u(i)
end do
close(10)

```

```

c..... make filter wavelet
11 write(6,*)'Causal (1) or Zero phase (2) ???'
  read(5,*) iflg
  if(iflg.eq.1) then
    write(6,*)
    *How many points will you use for the weighted moving
    average?
    read(5,*) np
    do i=1,np
      wt(i)=u(i)
    enddo
  else if(iflg.eq.2) then
    write(6,*)
    *How many points will you use for the weighted moving
    average?
    *The number must be odd.
    read(5,*) np
    nph=(np-1)/2
    do i=1,nph
      wt(i)=u(ndata-nph+i)
    enddo
    do i=nph+1,np
      wt(i)=u(i-nph)
    enddo
  else
    goto 11
  endif
  do i=1,np
    write(6,*) i,wt(i)
  enddo
  open(10,file=ofile,status='unknown')
  write(10,*) np,iflg
  write(10,*) dt, ' =dT (sec)'
  do i=1,np
    write(10,*) wt(i)
  end do
  close(10)
end

```

### FCONV.FOR

```

c*****
c Filtering by convolution in the time domain
c*****
program main
parameter (nmax=8192)
real wt(nmax),v(nmax),u(nmax)
character*80 ifile,ofile
c..... Input Time Series
write(6,*)'?? INPUT FILE NAME (Filter Wavelet)???'
read(5,'(a)') ifile
open(10,file=ifile,status='old')
read(10,*) np,iflg
read(10,*) dt
do i=1,np
  read(10,*) wt(i)
end do
close(10)
write(6,*)'?? INPUT FILE NAME (Input Signal)???'
read(5,'(a)') ifile
open(10,file=ifile,status='old')
read(10,*) ndata
read(10,*) dt1
if(dt1.ne.dt) then
  write(6,*)'Time interval unmatched!'
  stop
endif
do i=1,ndata
  read(10,*) u(i)
end do
close(10)
c..... Convolution
c Reverse filter time series
c ----> Coefficient of weighted moving average
if(iflg.eq.1) then
  write(6,*)'Causal Filtering'
  do n=1,np
    v(n)=0.0
    do m=1,n
      v(n)=v(n)+wt(n-m)*u(m)*dt
    enddo
  enddo
  do n=np+1,ndata
    v(n)=0.0
    do m=n-np,n
      v(n)=v(n)+wt(n-m)*u(m)*dt
    enddo
  enddo
else if(iflg.eq.2) then
  nph=(np-1)/2
  write(6,*)'Zero phase filtering'

```

```

do n=1,nph
v(n)=0.0
do m=1,n+nph
v(n)=v(n)+wt(n-m+nph+1)*u(m)*dt
enddo
enddo
do n=nph+1,ndata
v(n)=0.0
do m=n-nph,n+nph
v(n)=v(n)+wt(n-m+nph+1)*u(m)*dt
enddo
enddo
endif
c ..... Output Time Series
write(6,*)'???' OUTPUT FILE NAME ???'
read(5,'(a)') ofile
open(10,file=ofile,status='unknown')
write(10,*) ndata,'=N'
write(10,*) dt,'=dT (sec)'
do i=1,ndata
write(10,*) v(i)
end do
close(10)
write(6,*) 'Done'
stop
end

```

## BANDP1

```

subroutine bandp1(x,n,dt,fl,fh,fs,ap,as,ntype,nchara,ncausal)
real*4 x(n),y(65536),f(128)
c write(6,*)' Filter coefficients calculation started'
fl=fl*dt
fh=fh*dt
fs=fs*dt
c ..... Calculation of filter coefficients
if(ntype.eq.1) then
c Butterworth Filter
if(nchara.eq.1) then
c Low-Cut (High-Pass) Filter
call buthip(f,mf,gn,nchb,flu,fhu,fsu,ap,as)
else if(nchara.eq.2) then
c High-Cut (Low-Pass) Filter
write(6,*)' high cut'
call butlop(f,mf,gn,nchb,fhu,fsu,ap,as)
else if(nchara.eq.3) then
c Band-Pass Filter
call butpas(f,mf,gn,nchb,flu,fhu,fsu,ap,as)
else
c Band-Stop Filter
call butstp(f,mf,gn,nchb,flu,fhu,fsu,ap,as)
endif
else if(ntype.eq.2) then
c Chebyshev-I Filter
if(nchara.eq.1) then
c Low-Cut (High-Pass) Filter
call chbhip(f,mf,gn,nchb,eps,fhu,fsu,ap,as)
else if(nchara.eq.2) then
c High-Cut (Low-Pass) Filter
call chblop(f,mf,gn,nchb,eps,flu,fsu,ap,as)
else if(nchara.eq.3) then
c Band-Pass Filter
call chbpas(f,mf,gn,nchb,eps,flu,fhu,fsu,ap,as)
else
c Band-Stop Filter
call chbstp(f,mf,gn,nchb,eps,flu,fhu,fsu,ap,as)
endif
else if(ntype.eq.3) then
c Chebyshev-II Filter
if(nchara.eq.1) then
c Low-Cut (High-Pass) Filter
call tchhip(f,mf,gn,nchb,eps,fhu,fsu,ap,as)
else if(nchara.eq.2) then
c High-Cut (Low-Pass) Filter
call tchlop(f,mf,gn,nchb,eps,flu,fsu,ap,as)
else if(nchara.eq.3) then
c Band-Pass Filter
call tchpas(f,mf,gn,nchb,eps,flu,fhu,fsu,ap,as)
else

```

```

c Band-Stop Filter
call tchstp(f,mf,gn,nchb,eps,flu,fhu,fsu,ap,as)
endif
else
c Elliptic Filter
if(nchara.eq.1) then
c Low-Cut (High-Pass) Filter
call equhip(f,mf,gn,nchb,eps,ek,er,fhu,fsu,ap,as)
else if(nchara.eq.2) then
c High-Cut (Low-Pass) Filter
call equlop(f,mf,gn,nchb,eps,ek,er,flu,fsu,ap,as)
else if(nchara.eq.3) then
c Band-Pass Filter
call equipas(f,mf,gn,nchb,eps,ek,er,flu,fhu,fsu,ap,as)
else
c Band-Stop Filter
call equstp(f,mf,gn,nchb,eps,ek,er,flu,fhu,fsu,ap,as)
endif
endif
c write(6,*)' Filter coefficients are calculated.'
c write(6,(2x,a3,i3,2x,a3,f10.6))'mf=',mf,'gn=',gn
c write(6,(2x,i2,f10.6)) (i,f(i),i=1,mf)
c ..... Filtering subroutine
call gnf(x,y,n,f,mf,gn,ncausal)
return
end

```

## TRFIL.FOR

```

c*****
c Filtering in the time domain (Recursive filter)
c*****
include 'rfilter.for'
C
program main
parameter (nmax=8192)
real t(nmax),u(nmax)
character ifile*80,ofile*80
c ..... Input Time Series
write(6,*)'???' INPUT FILE NAME ???'
read(5,'(a)') ifile
write(6,*)'???' OUTPUT FILE NAME ???'
read(5,'(a)') ofile
open(10,file=ifile,status='old')
read(10,*) ndata
read(10,*) dt
do i=1,ndata
read(10,*) u(i)
t(i)=dt*real(i-1)
end do
close(10)
C ..... Filtering
10 write(6,*)' ??? Filter Type ???'
write(6,*)' Butterworth type ==> 1'
write(6,*)' Chebyshev-I type ==> 2'
write(6,*)' Chebyshev-II type ==> 3'
write(6,*)' Elliptic type ==> 4'
read(5,*) ntype
if(ntype.GT.4.OR.ntype.LT.1) GOTO 10
20 write(6,*)' ??? CHARACTERISTICS OF FILTER ???'
write(6,*)' Low-Cut (High-Pass) Filter ==> 1'
write(6,*)' High-cut (Low-Pass) Filter ==> 2'
write(6,*)' Band-Pass Filter ==> 3'
write(6,*)' Band-Stop Filter ==> 4'
read(5,*) nchara
if(nchara.GT.4.OR.nchara.LT.1) GOTO 20
if(nchara.eq.1) then
c Low-Cut (High-Pass) Filter
write(6,*)' FL
write(6,*)' /-----'
write(6,*)' /
write(6,*)' /
write(6,*)' /
write(6,*)' /-----/
write(6,*)' FS
write(6,*)' ??? FS,FL,AP,AS ???'
write(6,*)' AP,AS: Parameter defining the ripple '
write(6,*)' in pass band and stop band. '
write(6,*)' Use AP=0.1, AS=10.0, if you do '
write(6,*)' not like to think. '
read(5,*) fs,fl,ap,as
else if(nchara.eq.2) then
c High-Cut (Low-Pass) Filter

```



```

write(6,*) '      FH
write(6,*) '-----\
write(6,*) '
write(6,*) '
write(6,*) '
write(6,*) '
write(6,*) '
write(6,*) '      FS
write(6,*) ' ??? FH,FS,AP,AS ???'
write(6,*) '      AP,AS: Parameter defining the ripple '
write(6,*) '                  in pass band and stop band.
write(6,*) '                  Use AP=0.1, AS=10.0, if you do'
write(6,*) '                  not like to think.
read(5,*) fh,fs,ap,as
c
else if(nchara.eq.3) then
Band Pass Filter
write(6,*) '
write(6,*) '      FL      FH
write(6,*) '-----\
write(6,*) '
write(6,*) '
write(6,*) '
write(6,*) '
write(6,*) '      FS
write(6,*) ' ??? FL,FH,FS,AP,AS ???'
write(6,*) '      AP,AS: Parameter defining the ripple '
write(6,*) '                  in pass band and stop band.
write(6,*) '                  Use AP=0.1, AS=10.0, if you do'
write(6,*) '                  not like to think.
read(5,*) fl,fh,fs,ap,as
c
else if(nchara.eq.4) then
Band Stop Filter
write(6,*) '
write(6,*) '      FL      FH
write(6,*) '-----\
write(6,*) '
write(6,*) '
write(6,*) '
write(6,*) '
write(6,*) '      FS
write(6,*) ' ??? FL,FS,FH,AP,AS ???'
write(6,*) '      AP,AS: Parameter defining the ripple '
write(6,*) '                  in pass band and stop band.
write(6,*) '                  Use AP=0.1, AS=10.0, if you do'
write(6,*) '                  not like to think.
read(5,*) fl,fs,fh,ap,as
endif
c
write(6,*) fl,fh,fs,ap,as
write(6,*) ' ??? Causal or Zero-Phase ???'
write(6,*) '      Causal      Filter ==> 1'
write(6,*) '      Zero-Phase Filter ==> 2'
read(5,*) ncausal
call bandp1(u,ndata,dt,fl,fh,fs,ap,as,ntype,nchara,ncausal)
c.....Output Time Series
open(10,file=ofile)
write(10,*) ndata,'=N'
write(10,*) dt,'=dT (sec)'
do i=1,ndata
write(10,*) u(i)
end do
close(10)
write(*,*) 'Done'
stop
end

```

## RFILTER.FOR

```

c*****
c slave routine for trfilt.for
c*****
include 'buthip.f'
include 'butlop.f'
include 'butpas.f'
include 'butstp.f'
include 'chbhip.f'
include 'chblop.f'
include 'chbpas.f'
include 'chbstp.f'
include 'tchhip.f'
include 'tchlop.f'
include 'tchpas.f'
include 'tchstp.f'
include 'equhip.f'
include 'equlop.f'
include 'equpas.f'
include 'equstp.f'
include 'equpol.f'
include 'celin1.f'
include 'snthet.f'
include 'tandem.f'

```

```

include 'recfil.f'
c-----
subroutine bandp(x,n,dt,fl,fh,fs,ap,as,ntype,nchara,ncausal)
real*4 x(n),y(8192),f(128)
c write(6,*) 'Filter coefficients calculation started'
flu=fl*dt
fhu=fh*dt
fsu=fs*dt
c.....Calculation of filter coefficients
if(ntype.eq.1) then
c Butterworth Filter
if(nchara.eq.1) then
c Low-Cut (High-Pass) Filter
call buthip(f,mf,gn,nchb,flu,fsu,ap,as)
else if(nchara.eq.2) then
c High-Cut (Low-Pass) Filter
write(6,*) 'high cut'
call butlop(f,mf,gn,nchb,fhu,fsu,ap,as)
else if(nchara.eq.3) then
c Band-Pass Filter
call butpas(f,mf,gn,nchb,flu,fhu,fsu,ap,as)
else
c Band-Stop Filter
call butstp(f,mf,gn,nchb,flu,fhu,fsu,ap,as)
endif
else if(ntype.eq.2) then
c Chevyshev-I Filter
if(nchara.eq.1) then
c Low-Cut (High-Pass) Filter
call chbhip(f,mf,gn,nchb,eps,fhu,fsu,ap,as)
else if(nchara.eq.2) then
c High-Cut (Low-Pass) Filter
call chblop(f,mf,gn,nchb,eps,flu,fsu,ap,as)
else if(nchara.eq.3) then
c Band-Pass Filter
call chbpas(f,mf,gn,nchb,eps,flu,fhu,fsu,ap,as)
else
c Band-Stop Filter
call chbstp(f,mf,gn,nchb,eps,flu,fhu,fsu,ap,as)
endif
else if(ntype.eq.3) then
c Chevyshev-II Filter
if(nchara.eq.1) then
c Low-Cut (High-Pass) Filter
call tchhip(f,mf,gn,nchb,eps,fhu,fsu,ap,as)
else if(nchara.eq.2) then
c High-Cut (Low-Pass) Filter
call tchlop(f,mf,gn,nchb,eps,flu,fsu,ap,as)
else if(nchara.eq.3) then
c Band-Pass Filter
call tchpas(f,mf,gn,nchb,eps,flu,fhu,fsu,ap,as)
else
c Band-Stop Filter
call tchstp(f,mf,gn,nchb,eps,flu,fhu,fsu,ap,as)
endif
endif
else
c Elliptic Filter
if(nchara.eq.1) then
c Low-Cut (High-Pass) Filter
call equhip(f,mf,gn,nchb,eps,ek,er,fhu,fsu,ap,as)
else if(nchara.eq.2) then
c High-Cut (Low-Pass) Filter
call equlop(f,mf,gn,nchb,eps,ek,er,flu,fsu,ap,as)
else if(nchara.eq.3) then
c Band-Pass Filter
call equipas(f,mf,gn,nchb,eps,ek,er,flu,fhu,fsu,ap,as)
else
c Band-Stop Filter
call equstp(f,mf,gn,nchb,eps,ek,er,flu,fhu,fsu,ap,as)
endif
endif
endif
c write(6,*) 'Filter coefficients are calculated.'
write(6, '(2x,a3,i3,2x,a3,f10.6)') mf=',mf,gn=',gn
write(6, '(2x,i2,f10.6)') (i,f(i),i=1,mf)
c.....Filtering subroutine
call gnf(x,y,n,f,mf,gn,ncausal)
return

```

```

end
C-----
subroutine gnfx(x,y,nn,f,mf,gn,ncausal)
real*4 x(nn),y(8192),f(4)
if(ncausal.eq.2) then
  call tandem(x,y,nn,f,mf,-1)
  write(6,*) 'TANDEM-1'
  do 813 i=1,nn
813   x(i)=y(i)*gn
  endif
  call tandem(x,y,nn,f,mf, 1)
  write(6,*) 'TANDEM 1'
  do 815 i=1,nn
815   x(i)=y(i)*gn
  return
end

```

## DSEISM.FOR

```

program dseism
c*****
c digital filter equivalent to displacement seismometer
c*****
parameter (nmax=8192)
real v(nmax),u(nmax)
character*80 ifile,ofile
data pai/3.1415926535/
write(6,*) 'Natural Period:T0, Damping Factor:h'
write(6,*) 'Sampling interval:dt, Gain:G0'
read(5,*) t0,h,dt,g0
omg0=2.*pai/t0
write(6,*) omg0
omg0=(2./dt)*tan(omg0*dt/2.)
qq=tan(omg0*dt/2.)
qq2=qq*qq
a0= 1.0
a1=-2.0
a2= 1.0
b0= 1.0+2.*h*qq+qq2
b1=-2.0+2.*qq2
b2= 1.0-2.*h*qq+qq2
write(6,*) 'Numerator =',a0,a1,a2
write(6,*) 'Denominator=',b0,b1,b2
c-----input
write(6,*)'???' INPUT FILE NAME (Input Signal)???'
read(5, '(a)') ifile
open(10,file=ifile,status='old')
read(10,*) ndata
read(10,*) dt1
if(dt1.ne.dt) then
  write(6,*)'Time interval unmatched!'
  stop
endif
do i=1,ndata
  read(10,*) u(i)
end do
close(10)
c-----Recursive Filtering
write(6,*)'Causal Filtering'
v(1)= a0*u(1)/b0
v(2)=(a0*u(2)+a1*u(1)-b1*v(1))/b0
do i=3,ndata
  v(i)=(a0*u(i)+a1*u(i-1)+a2*u(i-2)
* -b1*v(i-1)-b2*v(i-2))/b0
  enddo
c----- Output Time Series
write(6,*)'???' OUTPUT FILE NAME ???'
read(5, '(a)') ofile
open(10,file=ofile,status='unknown')
write(10,*) ndata, '=N'
write(10,*) dt, '=dT (sec)'
do i=1,ndata
  write(10,*) v(i)
end do
close(10)
write(6,*) 'Done'
stop
end

```

## FWARP.FOR

```

c*****
c Calculate and Plot warped angular frequency
c*****
program main
parameter (nmax=8192)
real x(nmax),y(nmax)
data pai/3.1415926535/
c----- Input

```

```

write(6,*)'???' Dt, Ndata'
read(5,*) dt, ndata
td=ndata*dt
domg=2.*pai/td
omgnyq=pai/td
do i=1,ndata
  x(i)=real(i-1)*domg
  y(i)=(2./dt)*tan(x(i)*dt/2.)
end do
c----- Plot Time Series
umax=y(1)
umin=y(1)
do i=2,ndata
  if(y(i).gt.umax) umax=y(i)
  if(y(i).lt.umin) umin=y(i)
end do
cc ymax=max(abs(umax),abs(umin))*1.5
ymax=pai/td
c-----
call opengraphics
call setwindow(1,0.,0.,omgnyq*2,ymax)
call setviewpt(1,0.3,0.3,1,0.5)
call drawtitle('Warped Angular Frequency')
call setline('width=2')
call settxt('size=0.8')
call drawaxis('labelx=omg,labely=omg(k)')
c-----
call drawline(omgnyq,0.,omgnyq,ymax)
call setmark('type=solid-square,size=0.5')
call drawlines(ndata,x,y)
c-----
call drawmarks(ndata,x,y)
call drawtext(omgnyq*2*1.01,ymax,'pai/td')
call settxt('alignment=center-bottom')
call drawtext(omgnyq,0.0,'pai/td')
call closegraphics
end
DSEISMO.FOR
program dseism
c digital filter equivalent to displacement seismometer
parameter (nmax=8192)
real v(nmax),u(nmax)
character*80 ifile,ofile
data pai/3.1415926535/
write(6,*) 'Natural Period:T0, Damping Factor:h'
write(6,*) 'Sampling interval:dt, Gain:G0'
read(5,*) t0,h,dt,g0
omg0=2.*pai/t0
write(6,*) omg0
qq=tan(omg0*dt/2.)
qq2=qq*qq
a0= 1.0
a1=-2.0
a2= 1.0
b0= 1.0+2.*h*qq+qq2
b1=-2.0+2.*qq2
b2= 1.0-2.*h*qq+qq2
write(6,*) 'Numerator =',a0,a1,a2
write(6,*) 'Denominator=',b0,b1,b2
c-----input
write(6,*)'???' INPUT FILE NAME (Input Signal)???'
read(5, '(a)') ifile
open(10,file=ifile,status='old')
read(10,*) ndata
read(10,*) dt1
if(dt1.ne.dt) then
  write(6,*)'Time interval unmatched!'
  stop
endif
do i=1,ndata
  read(10,*) u(i)
end do
close(10)
c-----Recursive Filtering
write(6,*)'Causal Filtering'
v(1)= a0*u(1)/b0
v(2)=(a0*u(2)+a1*u(1)-b1*v(1))/b0
do i=3,ndata
  v(i)=(a0*u(i)+a1*u(i-1)+a2*u(i-2)
* -b1*v(i-1)-b2*v(i-2))/b0
  enddo
c----- Output Time Series
write(6,*)'???' OUTPUT FILE NAME ???'
read(5, '(a)') ofile
open(10,file=ofile,status='unknown')
write(10,*) ndata, '=N'
write(10,*) dt, '=dT (sec)'
do i=1,ndata
  write(10,*) v(i)
end do
close(10)
write(6,*) 'Done'
stop
end

```

```

open(10,file=ofile,status='unknown')
write(10,*) ndata,'=N'
write(10,*) dt,'=dT (sec)'
do i=1,ndata
  write(10,*) v(i)
end do
close(10)
write(6,*) 'Done'
stop
end

```

### VSEISMO.FOR

```

program vseism
c digital filter equivalent to displacement seismometer
parameter (nmax=8192)
real v(nmax),u(nmax)
character*80 ifile,ofile
data pai/3.1415926535/
write(6,*) 'Natural Period:T0, Damping Factor:h'
write(6,*) 'Sampling interval:dt, Gain:G0'
read(5,*) t0,h,dt,g0
omg0=2.*pai/t0
write(6,*) omg0
qq=tan(omg0*dt/2.)
qq2=qq*qq
a0= 1.0*g0
a1=-3.0*g0
a2= 3.0*g0
a3=-1.0*g0
b0=( 1.0+2.*h*qq+ qq2)*dt/2.
b1=(-1.0+2.*h*qq+3.*qq2)*dt/2.
b2=(-1.0-2.*h*qq+3.*qq2)*dt/2.
b3=( 1.0-2.*h*qq+ qq2)*dt/2.
write(6,*) 'Numerator =',a0,a1,a2,a3
write(6,*) 'Denominator=',b0,b1,b2,b3

```

```

c.....input
write(6,*)'???' INPUT FILE NAME (Input Signal)???'
read(5,'(a)') ifile
open(10,file=ifile,status='old')
read(10,*) ndata
read(10,*) dt1
if(dt1.ne.dt) then
  write(6,*)'Time interval unmatched!'
  stop
endif
do i=1,ndata
  read(10,*) u(i)
end do
close(10)

```

```

c.....Recursive Filtering
write(6,*)'Causal Filtering'
v(1)= a0*u(1)/b0
v(2)=(a0*u(2)+a1*u(1)-b1*v(1))/b0
v(3)=(a0*u(3)+a1*u(2)+a2*u(1)-b1*v(2)-b2*v(1))/b0
do i=4,ndata
  v(i)=(a0*u(i)+a1*u(i-1)+a2*u(i-2)+a3*u(i-3)
  * -b1*v(i-1)-b2*v(i-2)-b3*v(i-3))/b0
enddo

```

```

c ..... Output Time Series
100 write(6,*)'???' OUTPUT FILE NAME ???'
read(5,'(a)') ofile
open(10,file=ofile,status='unknown')
write(10,*) ndata,'=N'
write(10,*) dt,'=dT (sec)'
do i=1,ndata
  write(10,*) v(i)
end do
close(10)
write(6,*) 'Done'
stop
end

```

### ISEISMO.FOR

```

program iseism
c digital inverse filter to obtain ground displacement from velocity
seismograms
parameter (nmax=8192)
real v(nmax),u(nmax)
character*80 ifile,ofile
data pai/3.1415926535/

```

```

write(6,*) 'Natural Period:T0, Damping Factor:h'
write(6,*) 'Sampling interval:dt, Gain:G0'
read(5,*) t0,h,dt,g0
omg0=2.*pai/t0
write(6,*) omg0
qq=tan(omg0*dt/2.)
qq2=qq*qq
b0= 1.0*g0
b1=-3.0*g0
b2= 3.0*g0
b3=-1.0*g0
a0=( 1.0+2.*h*qq+ qq2)*dt/2.
a1=(-1.0+2.*h*qq+3.*qq2)*dt/2.
a2=(-1.0-2.*h*qq+3.*qq2)*dt/2.
a3=( 1.0-2.*h*qq+ qq2)*dt/2.
write(6,*) 'Numerator =',a0,a1,a2,a3
write(6,*) 'Denominator=',b0,b1,b2,b3

```

```

c.....input
write(6,*)'???' INPUT FILE NAME (Input Signal)???'
read(5,'(a)') ifile
open(10,file=ifile,status='old')
read(10,*) ndata
read(10,*) dt1
if(dt1.ne.dt) then
  write(6,*)'Time interval unmatched!'
  stop
endif
do i=1,ndata
  read(10,*) u(i)
end do
close(10)

```

```

c.....Recursive Filtering
write(6,*)'Causal Filtering'
v(1)= a0*u(1)/b0
v(2)=(a0*u(2)+a1*u(1)-b1*v(1))/b0
v(3)=(a0*u(3)+a1*u(2)+a2*u(1)-b1*v(2)-b2*v(1))/b0
do i=4,ndata
  v(i)=(a0*u(i)+a1*u(i-1)+a2*u(i-2)+a3*u(i-3)
  * -b1*v(i-1)-b2*v(i-2)-b3*v(i-3))/b0
enddo

```

```

c ..... Output Time Series
100 write(6,*)'???' OUTPUT FILE NAME ???'
read(5,'(a)') ofile
open(10,file=ofile,status='unknown')
write(10,*) ndata,'=N'
write(10,*) dt,'=dT (sec)'
do i=1,ndata
  write(10,*) v(i)
end do
close(10)
write(6,*) 'Done'
stop
end

```

### CSEISMO.FOR

```

program cseism
c digital filter to convert the longer natural period seismometer
parameter (nmax=8192)
real v(nmax),u(nmax)
character*80 ifile,ofile
data pai/3.1415926535/
write(6,*) 'Original Seismometer'
write(6,*) 'Natural Period:T0, Damping Factor:h0'
write(6,*) 'Sampling interval:dt, Gain:G0'
read(5,*) t0,h0,dt0,g0
omg0=2.*pai/t0
write(6,*) omg0
qq=tan(omg0*dt0/2.)
qq2=qq*qq
a0= 1.0+2.*h0*qq+qq2
a1=-2.0+2.*qq2
a2= 1.0-2.*h0*qq+qq2
write(6,*) 'Simulated Seismometer'
write(6,*) 'Natural Period:T1, Damping Factor:h1'
write(6,*) 'Sampling interval:dt, Gain:G1'
read(5,*) t1,h1,dt1,g1
omg1=2.*pai/t1
write(6,*) omg1

```

```

qq=tan(omg1*dt1/2.)
qq2=qq*qq
b0=( 1.0+2.*h1*qq+qq2)*g0/g1
b1=(-2.0      +2.*qq2)*g0/g1
b2=( 1.0-2.*h1*qq+qq2)*g0/g1
write(6,*)' Numerator   =',a0,a1,a2
write(6,*)' Denominator=',b0,b1,b2
c.....input
write(6,*)'?? INPUT FILE NAME (Input Signal)???'
read(5,'(a)') ifile
open(10,file=ifile,status='old')
read(10,*) ndata
read(10,*) dt
if(dt1.ne.dt.or.dt0.ne.dt) then
  write(6,*)'Time interval unmatched!'
  stop
endif
do i=1,ndata
  read(10,*) u(i)
end do
close(10)
c.....Recursive Filtering
write(6,*)'Causal Filtering'
v(1)= a0*u(1)/b0
v(2)=(a0*u(2)+a1*u(1)-b1*v(1))/b0
do i=3,ndata
  v(i)=(a0*u(i)+a1*u(i-1)+a2*u(i-2)
*      -b1*v(i-1)-b2*v(i-2))/b0
enddo
c ..... Output Time Series
write(6,*)'?? OUTPUT FILE NAME ???'
read(5,'(a)') ofile
open(10,file=ofile,status='unknown')
write(10,*) ndata,'=N'
write(10,*) dt,'=dT (sec)'
do i=1,ndata
  write(10,*) v(i)
end do
close(10)
write(6,*) 'Done'
stop
end

```

## PLTACC.FOR

```

include 'pltwv2.for'
include 'timesft.for'
include 'timecrt.for'
program pltacc
c k-net data viewer (raw acceleration data)
parameter(nmax=50000,nch=3)
integer ix(nmax)
real*4 x(nch,nmax),xx(nmax),amaxacc(nch)
character cid*6,filen*47,dummy*18,cdir*3,cmemo(nch)*80,
*      cstcode*6,cdrv*9,ccomp(3)*2
data aym,x0,y1,DTL/16.0,2.0,2.0,1.0/ fac/1./
data ccomp/'NS','EW','UD'/ tst,ted/0.01,500./
open(10,file='stlist.txt',status='old')
read(10,*) cdrv,nsite
do 2000 isite=1,nsite
  read(10,*) cid,tst
c write(6,'(A)') cdrv//cid//'.**'
  filen=cdrv//cid//'.ns'
  open(1,file=filen,status='old')
  write(6,'(a)') filen
  read(1,'(a18,i4,1x,i2,1x,i2,1x,i2,1x,i2)')
*      dummy,irgyear,irgmonth,irgday,irghour,irgmin
  read(1,'(a18,f5.1)')dummy,orglat
  read(1,'(a18,f5.1)')dummy,orglong
  read(1,'(a18,i4)') dummy,idepth
  orgdepth=real(idepth)
  read(1,'(a18,f4.1)')dummy,amag
  read(1,'(a18,a6)')dummy,cstcode
  read(1,'(a18,f8.4)')dummy,stlat
  read(1,'(a18,f8.4)')dummy,stlong
  read(1,'(a18,i4)') dummy,idepth
  stheight=real(idepth)
  read(1,'(a18,i4,1x,i2,1x,i2,1x,i2,1x,i2,1x,i2)')

```

```

*      dummy,ircyear,ircmonth,ircday,irchour,ircmin,ircsec
call timecrt(ircyear,ircmonth,ircday,irchour,ircmin,ircsec)
read(1,'(a18,i3)')dummy,iratio
sratio=real(iratio)
dt=1.0/sratio
read(1,'(a18,i3)')dummy,idur
ndata=idur*100
duration=real(idur)
read(1,'(a18,A3)')dummy,cdir
read(1,'(a18,10x,i7)')dummy,iscal
scale=2000./real(iscal)
read(1,'(a18,i4)')dummy,imaxacc
amaxacc(1)=real(imaxacc)
read(1,'(a18)')dummy
read(1,'(a80)')cmemo(1)
read(1,'(8(i8,1x))')(ix(i),i=1,ndata)
do 10 i=1,ndata
10  x(1,i)=real(ix(i))*scale
  close(1)
  filen=cdrv//cid//'.ew'
  open(1,file=filen,status='old')
  write(6,'(a)') filen
  do 20 kk=1,14
20  read(1,'(a18)')dummy
  read(1,'(a18,i4)')dummy,imaxacc
  amaxacc(2)=real(imaxacc)
  read(1,'(a18)')dummy
  read(1,'(a80)')cmemo(2)
  read(1,'(8(i8,1x))')(ix(i),i=1,ndata)
  do 30 i=1,ndata
30  x(2,i)=real(ix(i))*scale
  close(1)
  filen=cdrv//cid//'.ud'
  open(1,file=filen,status='old')
  write(6,'(a)') filen
  do 40 kk=1,14
40  read(1,'(a18)')dummy
  read(1,'(a18,i4)')dummy,imaxacc
  amaxacc(3)=real(imaxacc)
  read(1,'(a18)')dummy
  read(1,'(a80)')cmemo(3)
  read(1,'(8(i8,1x))')(ix(i),i=1,ndata)
  do 50 i=1,ndata
50  x(3,i)=real(ix(i))*scale
  close(1)
  ntst=int(tst)
  nst=int(tst/dt+0.5)
  ned=int(ted/dt+0.5)
  call timesft(ntst,ircyear,ircmonth,ircday,
*      irchour,ircmin,ircsec)
  filen=cdrv//cid//'.ac.ps'
  open(25,file=filen,status='unknown')
  write(6,'(a)') filen
  CALL PLOTS(25)
  call yoko(25)
  call symbol(1.5,5.0,0.8,cid,90.0,6)
  call number(1.5,11.00,0.25,real(ircyear),90.0,0)
  call number(1.5,12.25,0.25,real(ircmonth),90.0,0)
  call number(1.5,13.00,0.25,real(ircday),90.0,0)
  call number(1.5,13.75,0.25,real(irchour),90.0,0)
  call number(1.5,14.50,0.25,real(ircmin),90.0,0)
  call number(1.5,15.25,0.25,real(ircsec),90.0,0)
  CALL FACTOR(FAC)
  amaxx=max1(amaxacc(1),amaxacc(2),amaxacc(3))
  if(amaxx.lt.2000.) xxmax=2000.
  if(amaxx.lt.1000.) xxmax=1000.
  if(amaxx.lt. 500.) xxmax= 500.
  if(amaxx.lt. 200.) xxmax= 200.
  if(amaxx.lt. 100.) xxmax= 100.
  if(amaxx.lt. 50.) xxmax= 50.
  if(amaxx.lt. 20.) xxmax= 20.
  if(amaxx.lt. 10.) xxmax= 10.
  if(amaxx.lt. 5.) xxmax= 5.
  if(amaxx.lt. 2.) xxmax= 2.
  if(amaxx.lt. 1.) xxmax= 1.
  call plot(x0,0.,-3)
  CALL NEWPEN(1)
  ay=aym

```

```

ndatamax=min1(real(ned),real(ndata),24.*dtl/dt)+nst-1
if(ndatamax.lt.ndata) call symbol(26.,0.5,0.3,'C',0.0,1)
do 1000 kk=1,3
  call symbol(2.5,ay+y1+0.3,0.4,ccomp(kk),0.0,2)
  call symbol(4.0,ay+y1+0.3,0.3,'Max=',0.0,4)
  call number(5.5,ay+y1+0.3,0.3,amaxacc(kk),0.0,1)
  do 500 i=1,ndata
500    xx(i)=x(kk,i+nst-1)
        sum=xx(1)
        do 600 i=2,ndata
600    sum=sum+xx(i)
        ave=sum/real(ndata)
        do 610 i=1,ndata
610    xx(i)=xx(i)-ave
        CALL PLOT( 2., AY,-3)
        call pltwv2(xx,ndatamax-nst+1,dt,dtl,xxmax,y1,1)
        CALL PLOT(-2.,-AY,-3)
        ay=ay-3.*yl
1000 continue
    call plote(25)
    close(25)
2000 continue
    close(10)
    stop
    end
PLTWV2.FOR
  subroutine pltwv2(d,n,dt,dtl,YMAX,y1,ind)
  REAL*4 d(n)
  data htitle,hcha,htick/0.4,0.2,0.2/
  tlong=real(int(real(n)*dt/dtl)+1)
  ntlong=int(tlong)
  dsize=dt/dtl
  IF(YMAX.LT.1.0E-30) YMAX=1.0E-30
  call plot( 0.,-y1,3)
  call plot(tlong,-y1,2)
  call plot( 0., y1,3)
  call plot(tlong, y1,2)
  call symbol(0.5*tlong-htitle*1.5,
  *
-(y1+htick*2.+htitle+0.1),htitle,'sec',0.0,3)
  do 10 i=0,ntlong
    call plot(real(i),-y1,3)
10  call plot(real(i),-y1-htick,2)
  do 11 i=0,ntlong
    call plot(real(i),y1,3)
11  call plot(real(i),y1+htick,2)
  do 12 i=0,ntlong,2
12  call number(real(i)-hcha*2.,-(y1+htick+hcha+0.1),hcha,
  *      real(i)*dtl,0.0,1)
  DYL=YMAX/YL
  call plot(0.,-y1,3)
  call plot(0., y1,2)
  if(ind.eq.1)
  *call symbol(-(htick*2.+hcha*3.5),-htitle*1.5,
  *      htitle,'gal',90.0,3)
  if(ind.eq.2)
  *call symbol(-(htick*2.+hcha*3.5),-htitle*1.5,
  *
htitle,'kine',90.0,4)
  if(ind.eq.3)
  *call symbol(-(htick*2.+hcha*3.5),-htitle*1.5,
  *
htitle,'micron',90.0,6)
  do 110 i=-int(y1+0.5),int(y1+0.5)
    call plot(-htick,real(i),3)
    call plot( 0.,real(i),2)
110 continue
  if(ymax.ge. 1.) then
    ncfra=4
  if(ymax.ge. 10.) then
    ncfra=5
  if(ymax.ge.100.) then
    ncfra=6
  if(ymax.ge.1000.) ncfra=7.
  endif
  endif
  endif
  xx=-(htick+hcha*real(ncfra)+0.1)

```

```

yy=y1+hcha*0.5
call number(xx,-yy,hcha,-ymax,0.0,1)
call number(-hcha*3.25,-hcha*0.5,hcha,0.0,0.0,1)
call number(xx, yy,hcha,ymax,0.0,1)
call plot(tlong, y1,3)
call plot(tlong,-y1,2)
call plot(0.,d(1)/DYL,3)
do 130 i=2,n
  Y=D(I)/DYL
  call plot((i-1)*dsize,Y,2)
130 continue
  return
  end
TIMECRT.FOR
  subroutine timecrt(ircyear,ircmonth,ircday,
  *      irchour,ircmin,ircsec)
  if(ircsec.ge.15) then
    ircsec=ircsec-15
  else
    if(ircmin.ge.1) then
      ircmin=ircmin-1
      ircsec=ircsec+45
    else
      if(irchour.ge.1) then
        irchour=irchour-1
        ircmin =59
        ircsec=ircsec+45
      else
        ircday =ircday-1
        irchour=23
        ircmin =59
        ircsec=ircsec+45
      endif
    endif
  endif
  return
  end
TIMESFT.FOR
  subroutine timesft(ntst,ircyear,ircmonth,ircday,
  *      irchour,ircmin,ircsec)
  if(ircsec+ntst.lt.60) then
    ircsec=ircsec+ntst
  else
    if(ircmin.lt.59) then
      ircmin=ircmin+1
      ircsec=ircsec-60+ntst
    else
      if(irchour.lt.23) then
        irchour=irchour+1
        ircmin =0
        ircsec=ircsec-60+ntst
      else
        ircday =ircday+1
        irchour=0
        ircmin =0
        ircsec=ircsec-60+ntst
      endif
    endif
  endif
  return
  end
SPLOT.FOR
  include 'cfft.for'
  include 'timecrt.for'
  include 'timesft.for'
  parameter(nmax=50000,nch=3)
  complex c(nmax*2)
  integer ix(nmax)
  real*4
  x(nch,nmax),xx(nmax),amaxacc(nch),fr(nmax),yy(nmax),
  *      amaxdsp(nch),y(nch,nmax)
  character
  cid*6,flen*47,dummy*18,cdir*3,cmemo(nch)*80,cstcode*6,
  *      cdrv*9,ccomp(3)*2
  data yl/2.0/ fac/0.7/
  data ccomp/'NS','EW','UD'/ tst,ted/ 0.01,70./ asmfth/0.02/
  data pai,c0/3.141592,(0.0,0.0)/ ircdelay/0.0/

```

```

data NBX,NBY,ABX,ABY/3,3,2.5,2.5/ HCH,HTP/0.4,0.1/
FMIN/0.1/
open(10,file='stlist.txt',status='old')
read(10,*) cdrv,nsite
do 2000 isite=1,nsite
read(10,*) cid,tst,tcd
write(6,('a'))cdrv//cid//.*
filen=cdrv//cid//.ns'
open(1,file=filen,status="old")
read(1,('a18,i4,1x,i2,1x,i2,1x,i2,1x,i2'))
* dummy,irgyear,irgmonth,irgday,irghour,irgmin
read(1,('a18,f5.1'))dummy,orglat
read(1,('a18,f5.1'))dummy,orglong
read(1,('a18,i4')) dummy,idepth
orgdepth=real(idepth)
read(1,('a18,f4.1'))dummy,amag
read(1,('a18,a6'))dummy,cstcode
read(1,('a18,f8.4'))dummy,stlat
read(1,('a18,f8.4'))dummy,stlong
read(1,('a18,i4')) dummy,idepth
stheight=real(idepth)
read(1,('a18,i4,1x,i2,1x,i2,1x,i2,1x,i2'))
* dummy,ircyear,ircmonth,ircday,irchour,ircmin,ircsec
call timecrt(ircyear,ircmonth,ircday,irchour,ircmin,ircsec)
read(1,('a18,i3'))dummy,iratio
sratio=real(iratio)
dt=1.0/sratio
read(1,('a18,i3'))dummy,idur
ndata=idur*100
duration=real(idur)
read(1,('a18,A3'))dummy,cdir
read(1,('a18,10x,i7'))dummy,iscalc
scale=2000./real(iscalc)
read(1,('a18,i4'))dummy,imaxacc
amaxacc(1)=real(imaxacc)
read(1,('a18'))dummy
read(1,('a80'))cmemo(1)
read(1,('8(i8,1x)'))(ix(i),i=1,ndata)
do 10 i=1,ndata
10 x(1,i)=real(ix(i))*scale
close(1)
filen=cdrv//cid//.ew'
open(1,file=filen,status="old")
do 20 kk=1,14
20 read(1,('a18'))dummy
read(1,('a18,i4'))dummy,imaxacc
amaxacc(2)=real(imaxacc)
read(1,('a18'))dummy
read(1,('a80'))cmemo(2)
read(1,('8(i8,1x)'))(ix(i),i=1,ndata)
do 30 i=1,ndata
30 x(2,i)=real(ix(i))*scale
close(1)
filen=cdrv//cid//.ud'
open(1,file=filen,status="old")
do 40 kk=1,14
40 read(1,('a18'))dummy
read(1,('a18,i4'))dummy,imaxacc
amaxacc(3)=real(imaxacc)
read(1,('a18'))dummy
read(1,('a80'))cmemo(3)
read(1,('8(i8,1x)'))(ix(i),i=1,ndata)
do 50 i=1,ndata
50 x(3,i)=real(ix(i))*scale
close(1)
ncomp=nch
ntst=int(tst)
nst=int(tst/dt+0.5)
ned=int(ted/dt+0.5)
write(6,*)nst,ned
c ned=ndata
call timesft(ntst,ircyear,ircmonth,ircday,
* irghour,ircmin,ircsec)
nn=ned-nst+1
n2e=int(alog10(real(nn))/0.30103+0.5)
mm=2*n2e
mm2=mm*2
mwin=int(0.1*real(mm))

```

```

datal=dt*real(mm)
write(6,*)datal
c do 106 i=1,mm
106 fr(i)=real(i-1)/datal
C FMAX=FMIN*(10**REAL(NBX))
FMAX=45.
NMX=INT(FMAX*DATAL)+1
NMN=INT(FMIN*DATAL)+1
do 525 mcomp=1,ncomp
do 207 i=1,nn
207 xx(i)=x(mcomp,i+nst-1)
c linear trend fitting
b1=xx(1)
do 392 i=2,nn
392 b1=b1+real(i)*xx(i)
b2=xx(1)
do 394 i=1,nn
394 b2=b2+xx(i)
a0=real(nn*(nn+1)*(2*nn+1))/6.*dt
b0=real(nn*(nn+1))/2.
c0=b0*dt
d0=real(nn)
f0=a0*d0-b0*c0
aa=(d0*b1-b0*b2)/f0
bb=(-c0*b1+a0*b2)/f0
do 396 i=1,nn
396 xx(i)=xx(i)-aa*real(i)*dt-bb
do 409 i=1,mwin
409 xx(i)=xx(i)*SIN(PI/2.*REAL(I-1)/REAL(MWIN))
do 412 i=nn-mwin,nn
412 xx(i)=xx(i)*SIN(PI/2.*REAL(nn-i)/REAL(MWIN))
do 413 i=nn+1,mm
413 xx(i)=0.0
do 414 i=1,mm
414 c(i)=cplx(xx(i))
write(6,*) mm
CALL cfft(c,mm,-1)
do 518 i=2,mm
518 xx(i)=cabs(c(i))*dt/(2.*PI*FR(I))*2
do 519 i=2,mm
519 yy(i)=cabs(c(i))*dt
amaxdsp(mcomp)=abs(xx(nmn))
do 521 i=nmn,nmx
521 if(abs(xx(i)).GT.amaxdsp(mcomp))
amaxdsp(mcomp)=abs(xx(i))
do 522 i=2,mm
522 x(mcomp,i)=xx(i)
amaxacc(mcomp)=abs(yy(nmn))
do 523 i=nmn,nmx
523 if(abs(yy(i)).GT.amaxacc(mcomp))
amaxacc(mcomp)=abs(yy(i))
do 524 i=2,mm
524 y(mcomp,i)=yy(i)
525 continue
filen=cdrv//cid//sp.ps'
write(6,*) filen
open(25,file=filen,status='unknown')
CALL PLOTS(25)
call yoko(25)
smaxx=amax1(amaxacc(1),amaxacc(2),amaxacc(3))
if(smaxx.lt.1000. ) sxmax=1000.
if(smaxx.lt. 100. ) sxmax= 100.
if(smaxx.lt. 10. ) sxmax= 10.
if(smaxx.lt. 1. ) sxmax= 1.
if(smaxx.lt. 0.1 ) sxmax= 0.1
if(smaxx.lt. 0.01 ) sxmax= 0.01
if(smaxx.lt. 0.001 ) sxmax= 0.001
if(smaxx.lt. 0.0001 ) sxmax= 0.0001
if(smaxx.lt. 0.00001 ) sxmax= 0.00001
if(smaxx.lt. 0.000001 ) sxmax= 0.000001
sminacc=sxmax*10**real(-nby-1)
smaxx=amax1(amaxdsp(1),amaxdsp(2),amaxdsp(3))
if(smaxx.lt. 100. ) sxmax= 100.
if(smaxx.lt. 10. ) sxmax= 10.
if(smaxx.lt. 1. ) sxmax= 1.
if(smaxx.lt. 0.1 ) sxmax= 0.1
if(smaxx.lt. 0.01 ) sxmax= 0.01
if(smaxx.lt. 0.001 ) sxmax= 0.001

```

```

if(smaxx.lt. 0.0001 ) sxmax= 0.0001
if(smaxx.lt. 0.00001 ) sxmax= 0.00001
if(smaxx.lt. 0.000001) sxmax= 0.000001
smindsp=sxmax*10**real(-nby-1)
call plot(2.0,1.65,-3)
call symbol( 0.0,0.5,1.2,cid,0.0,6)
call symbol(10.0,0.0,0.3,'Event:',0.0,6)
call number(11.8,0.0,0.3, orglat,0.0,3)
call number(15.4,0.0,0.3,orglong,0.0,3)
call number(19.0,0.0,0.3,orgdepth,0.0,0)
call symbol(10.0,0.5,0.3,'Site :',0.0,6)
call number(11.8,0.5,0.3, stlat,0.0,3)
call number(15.4,0.5,0.3,stlong,0.0,3)
call number(19.0,0.5,0.3,sheight/1000.,0.0,3)
call symbol(10.0,1.0,0.4,'Start Time ',0.0,11)
call plot(14.4,0.0,-3)
call number( 0.0,1.0,0.4,real(ircyear) ,0.0,0)
call number( 2.0,1.0,0.4,real(ircmonth),0.0,0)
call number( 3.2,1.0,0.4,real(ircday) ,0.0,0)
call number( 4.8,1.0,0.4,real(irchour) ,0.0,0)
call number( 6.0,1.0,0.4,real(ircmin) ,0.0,0)
call number( 7.2,1.0,0.4,real(ircsec)+real(ircdelay),0.0,0)
call plot(-14.4,0.5,-3)
CALL FACTOR(FAC)
CALL NEWPEN(1)
call plot(2.,5.5,-3)
xl=abx*real(nbx)+4.0
yl=aby*real(nby)+2.0
do 800 kk=1,3
  call symbol(abx+0.5,y1,0.4,ccomp(kk),0.0,2)
  do 790 i=1,mm
    xx(i)=x(kk,i)
    CALL
  ESMTHF(XX(nmn),XX(nmn),FR(nmn),NMx,asmthf)
  CALL
  PLTLOG(NBX,NBY,ABX,ABY,HCH,HTP,FMIN,SMINDSP,1)
  CALL NEWPEN(1)
  CALL
  PTLINE(FR,XX,ABX,ABY,FMIN,SMINDSP,NMAX,NMN,NMX)
  CALL NEWPEN(1)
  CALL PLOT(xl,0.,-3)
  800 continue
  call plot(-xl*3.,12.5,-3)
  do 1000 kk=1,3
    call symbol(abx+0.5,y1,0.4,ccomp(kk),0.0,2)
    do 900 i=1,mm
      yy(i)=y(kk,i)
      CALL ESMTHF(yy(nmn),yy(nmn),FR(nmn),NMx,asmthf)
      CALL
    PLTLOG(NBX,NBY,ABX,ABY,HCH,HTP,FMIN,SMINACC,3)
    CALL NEWPEN(1)
    CALL
    PTLINE(FR,yy,ABX,ABY,FMIN,SMINACC,NMAX,NMN,NMX)
    CALL NEWPEN(1)
    CALL PLOT(xl,0.,-3)
  1000 continue
  Call plot(0.,-12.,-3)
  call plote(25)
  2000 continue
  stop
  end
C+++++
SUBROUTINE
PLTLOG(N,M,DX,DY,HEIT,TIP,XMIN,YMIN,IDX)
C
C N,M : BLOCK NUMBER OF X-AXIS & Y-AXIS
C DX,DY : BLOCK LENGTH OF X-AXIS & Y-AXIS
(CM)
C TIP : LENGTH OF SCALE TIP IN AXISES
(CM)
C HEIT : HEIGHT OF CHARACTER
(CM)
C XMIN,YMIN : (XMIN,YMIN) REFERENCE POINT
DATA
C
CHARACTER
XUNIT*7,YUNIT*13,XNAME*10,YNAME*20
CHARACTER YAX(5)*8,YNA(5)*20

```

```

DATA YAX/' CM SEC ','CM/S SEC','GAL SEC ','
'/
DATA YNA/'DISPLACEMENT SPECTRA',' VELOCITY
SPECTRA ',
* 'ACCELERATION SPECTRA','SPECTRUM
RATIO (M/A)',
* 'SPECTRUM RATIO (S/M)'/
XUNIT=' (Hz) '
YUNIT=' '//YAX(IDX)/' '
XNAME='FREQUENCY '
YNAME=YNA(IDX)
CALL NEWPEN(3)
DO 100 I=0,N
  XL=XMIN*10**I
  IX=NINT(ALOG10(XL)+0.1)
  IF(IX.GE.2.OR.IX.LE.-1) GO TO 400
  IF(IX.EQ.1) THEN
    CALL SYMBOL(I*DX-HEIT,-0.3-HEIT,HEIT,'10',0.0,2)
  ELSE
    CALL
    NUMBER(I*DX-2*HEIT,-HEIT-0.3,HEIT,10.**IX+0.05,0.0,1)
  ENDIF
  GO TO 100
400 CALL SYMBOL(I*DX-HEIT,-0.3-HEIT,HEIT,'10',0.0,2)
  CALL
  NUMBER(I*DX+HEIT,-0.3 ,HEIT*0.5,REAL(IX),0.0,0)
  100 CONTINUE
  DO 101 I=0,M
    YL=YMIN*10**I
    IY=NINT(ALOG10(YL)+0.1)
    IF(IY.GE.2.OR.IY.LE.-1) GO TO 401
    IF(IY.EQ.1) THEN
      CALL SYMBOL(-0.3-3*HEIT,I*DY,HEIT,'10',0.0,2)
    ELSE
      CALL
      NUMBER(-0.3-4*HEIT,I*DY,HEIT,10.**IY+0.05,0.0,0)
    ENDIF
    GO TO 101
401 CALL SYMBOL(-0.3-3*HEIT,I*DY,HEIT,'10',0.0,2)
  CALL
  NUMBER(-0.3-HEIT,I*DY+HEIT,HEIT*0.5,REAL(IY),0.0,0)
  101 CONTINUE
  CALL
  SYMBOL(N*DX/2-2.5*HEIT,-1.-2.7*HEIT,HEIT*1.,XUNIT,0., 7)
  CALL
  SYMBOL(N*DX/2-3.8*HEIT,-1.-1.5*HEIT,HEIT*1.,XNAME,0.,10)
  CALL
  SYMBOL(-4*HEIT,N*DX/2-6*HEIT,HEIT*1., YUNIT,90.,13)
  CALL
  SYMBOL(N*DX/2-7.5*HEIT,M*DY+1.5*HEIT,HEIT*1.,YNAME,0.,
20)
  CALL NEWPEN(2)
  DO 1 I=0,N
    CALL PLOT(I*DX, 0.0,3)
    CALL PLOT(I*DX,M*DY,2)
    IF(I.EQ.N) GO TO 1
    DO 11 J=2,9
      CALL PLOT(I*DX+ALOG10(REAL(J))*DX,0.0,3)
    11 CALL PLOT(I*DX+ALOG10(REAL(J))*DX,TIP,2)
    CALL PLOT(I*DX+ALOG10(5.)*DX,0.0,3)
    CALL PLOT(I*DX+ALOG10(5.)*DX,TIP*1.5,2)
  1 CONTINUE
  DO 2 I=0,M
    CALL PLOT(0.0,I*DY,3)
    CALL PLOT(N*DX,I*DY,2)
    IF(I.EQ.M) GO TO 2
    DO 12 J=2,9
      CALL PLOT(0.0,I*DY+ALOG10(REAL(J))*DY,3)
    12 CALL PLOT(TIP,I*DY+ALOG10(REAL(J))*DY,2)
    CALL PLOT(0.0,I*DY+ALOG10(5.)*DY,3)
    CALL PLOT(TIP*1.5,I*DY+ALOG10(5.)*DY,2)
  2 CONTINUE
  CALL NEWPEN(1)
  RETURN
  END
C+++++
SUBROUTINE
PTLINE(F,D,DX,DY,XMIN,YMIN,NN,NMIN,NMAX)

```

```

REAL*4 X(4096),Y(4096),D(NN),F(NN)
DO 80 I=NMIN,NMAX
  IF(F(I).LT.1.E-10.OR.D(I).LT.1.E-10) GO TO 80
  X(I)=DX*ALOG10(F(I)/XMIN)
  Y(I)=DY*ALOG10(D(I)/YMIN)
80 CONTINUE
CALL PLOT(X(NMIN),Y(NMIN),3)
DO 300 I=NMIN+1,NMAX
  IF(Y(I-1).GE.0.0.AND.Y(I).GT.0.0) THEN
    CALL PLOT(X(I),Y(I),2)
  ELSE
    CALL PLOT(X(I),Y(I),3)
  ENDIF
300 CONTINUE
301 RETURN
END
C+++++
SUBROUTINE ESMTHF(AI,AO,F,NF,SMTHF)
C * ALISAR ESPECTRO
REAL*4 AI(NF),AO(NF),F(NF),MADD
IF(SMTHF.LE.0.0) GO TO 100
AO(1)=AI(1)
DO 35 I=2,NF
  F1=(1.0-0.5*SMTHF)*F(I)
  F2=(1.0+0.5*SMTHF)*F(I)
  IF(F2.LE.F(NF)) GO TO 31
  MNF=I-1
  GO TO 36
31 SUM=0.0
  MADD=0.0
  DO 32 K=1,NF
    IF(F(K).LT.F1) GO TO 32
    IF(F(K).GT.F2) GO TO 33
    SUM=SUM+AI(K)
    MADD=MADD+1
32 CONTINUE
33 CONTINUE
  AO(I)=SUM/MADD
35 CONTINUE
36 CONTINUE
  NF=MNF
  RETURN
100 CONTINUE
  DO 101 I=1,NF
101 AO(I)=AI(I)
  RETURN
  END
ACC2VEL.FOR
program acc2vel
c k-net data integrater (convert to velocity data)
parameter(nmax=65536,nch=3)
integer ix(nmax),imaxacc(nch),imaxvel(nch)
real*4 x(nch,nmax),xx(nmax),amaxacc(nch),amaxvel(nch)
complex cx(nmax)
character cid*6,filen(nch)*20,dummy(16)*18,cdir(nch)*3,
*
cmemo(nch)*80,cstcode*6,cdrv*9,ccomp(nch)*2,filenm*18
data nst/1/ ntype,nchara,ncausal/1,3,2/
ccomp/'NS','EW','UD'/
data ap,as/0.1,10.0/
ned=nmax
open(10,file='stlist.txt',status='old')
read(10,*) cdrrv,nsite
do 5000 isite=1,nsite
read(10,*) cid,tst,ted,fl,fh,fs
write(6,'(a)cdrrv/cid//.*')
filen(1)=cdrrv/cid//'.ns'
open(1,file=filen(1),status="old")
read(1,'(a18,i4,1x,i2,1x,i2,1x,i2,1x,i2)')
* dummy(1),irgyear,irgmonth,irgday,irghour,irgmin
read(1,'(a18,f5.1)')dummy(2),orglat
read(1,'(a18,f5.1)')dummy(3),orglong
read(1,'(a18,i4)') dummy(4),idepth
orgdepth=real(idepth)
read(1,'(a18,f4.1)')dummy(5),amag
read(1,'(a18,a6)')dummy(6),cstcode
read(1,'(a18,f8.4)')dummy(7),stlat
read(1,'(a18,f8.4)')dummy(8),stlong
read(1,'(a18,i4)') dummy(9),istdepth
stheight=real(istdepth)
read(1,'(a18,i4,1x,i2,1x,i2,1x,i2,1x,i2,1x,i2)')
* dummy(10),ircyear,ircmonth,ircday,irchour,ircmin,ircsec
read(1,'(a18,i3)')dummy(11),iratio
sratio=real(iratio)
dt=1.0/sratio
read(1,'(a18,i3)')dummy(12),idur
ndata=idur*100-8
nst=int(tst/dt+0.5)*0.0
c ned=min(int(ted/dt+0.5),ndata)
ned=ndata
duration=real(idur)
read(1,'(a18,A3)')dummy(13),cdir(1)
read(1,'(a18,10x,i7)')dummy(14),iscale
scale=2000./real(iscale)
read(1,'(a18,i4)')dummy(15),imaxacc(1)
amaxacc(1)=real(imaxacc(1))
read(1,'(a18,i4,1x,i2,1x,i2,1x,i2,1x,i2,1x,i2)')
* dummy(16),icryear,icrmonth,icrday,icrhour,icrmin,icrsec
read(1,'(a80)')cmemo(1)
read(1,'(8(i8,1x))')(ix(i),i=1,ndata)
do 5 i=1,nmax
5 cx(i)=(0.0,0.0)
do 10 i=1,ndata
10 x(1,i)=real(ix(i))*scale
close(1)
filen(2)=cdrrv/cid//'.ew'
open(1,file=filen(2),status="old")
do 20 kk=1,14
20 read(1,'(a18)')dummy(kk)
cdir(2)='EW'
read(1,'(a18,i4)')dummy(15),imaxacc(2)
amaxacc(2)=real(imaxacc(2))
read(1,'(a18)')dummy(16)
read(1,'(a80)')cmemo(2)
read(1,'(8(i8,1x))')(ix(i),i=1,ndata)
do 30 i=1,ndata
30 x(2,i)=real(ix(i))*scale
close(1)
filen(3)=cdrrv/cid//'.ud'
open(1,file=filen(3),status="old")
do 40 kk=1,14
40 read(1,'(a18)')dummy(kk)
cdir(3)='UD'
read(1,'(a18,i4)')dummy(15),imaxacc(3)
amaxacc(3)=real(imaxacc(3))
read(1,'(a18)')dummy(16)
read(1,'(a80)')cmemo(3)
read(1,'(8(i8,1x))')(ix(i),i=1,ndata)
do 50 i=1,ndata
50 x(3,i)=real(ix(i))*scale
close(1)
nn=ned-nst+1
idur=int(real(nn)*dt)
td =real(idur)
nn=td/dt
ntst=int(tst)
do 500 kk=1,3
do 180 i=1,nn
180 xx(i)=x(kk,i+nst-1)
c linear trend fitting
b1=xx(1)
do 192 i=2,nn
192 b1=b1+real(i)*xx(i)
b2=xx(1)
do 194 i=1,nn
194 b2=b2+xx(i)
a=real(nn*(nn+1)*(2*nn+1))/6.*dt
b=real(nn*(nn+1))/2.
c=b*dt
d=real(nn)
f=a*d-b*c
aa=( d*b1-b*b2)/f
bb=(-c*b1+a*b2)/f
do 196 i=1,nn
196 xx(i)=xx(i)-aa*real(i)*dt-bb
c band pass filtering

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```

        call bandp1(xx,nn,dt,fl,fh,fs,ap,as,nstype,nchara,ncausal)
c integration
200  xx(1)=xx(1)*dt
    do 300 i=2,nn
300  xx(i)=xx(i-1)+xx(i)*dt
    amaxvel(kk)=abs(xx(1))
    do 400 i=2,nn
400  if(amaxvel(kk).lt.abs(xx(i))) amaxvel(kk)=abs(xx(i))
    imaxvel(kk)=int(amaxvel(kk)*1000.+0.5)
    do 410 i=1,nn
410  x(kk,i)=xx(i)
500 continue
    do 600 kk=1,nch
        filenm=filen(kk)(1:15)//'VL.'
        filen(kk)=filenm//ccomp(kk)
600  write(6,'(1x,a23)') filen(kk)
c 25 micro kine per digit
    scale=2.5e-5
    iscale=int(200/scale+0.5)
c
    do 1000 kk=1,nch
    do 1010 i=1,nn
1010  ix(i)=int(x(kk,i)/scale)
    open(1,file=filen(kk),status='unknown')
    write(1,'(a18,i4,1x,i2,1x,i2,1x,i2,1x,i2)')
    *          dummy(1),
    *
irgyear,irgmonth,irgday,irghour,irgmin
    write(1,'(a18,f5.1)')dummy(2),orglat
    write(1,'(a18,f5.1)')dummy(3),orglong
    write(1,'(a18,i4)') dummy(4),idepth
    write(1,'(a18,f4.1)')dummy(5),amag
    write(1,'(a18,a6)') dummy(6),cstcode
    write(1,'(a18,f8.4)')dummy(7),stlat
    write(1,'(a18,f8.4)')dummy(8),stlong
    write(1,'(a18,i4)') dummy(9),istdepth
    write(1,'(a18,i4,1x,i2,1x,i2,1x,i2,1x,i2,1x,i2)')
    *          dummy(10),
    *
ircyear,ircmonth,ircday,irchour,ircmin,ircsec
    write(1,'(a18,i3)') dummy(11),iratio
    write(1,'(a18,i3)') dummy(12),idur
    write(1,'(a18,A3)') dummy(13),cdir(kk)
    write(1,'(a18,a10,i7)')dummy(14),'200(kine)',iscale
    write(1,'(a18,i8)') 'Max. Vel.(mkine) ',imaxvel(kk)
    write(1,'(a18,i4,1x,i2,1x,i2,1x,i2,1x,i2,1x,i2)')
    *          dummy(16),
    *
icryear,icrmonth,icrday,icrhour,icrmin,ircsec
    write(1,'(a80)') cmemo(kk)
    write(1,'(8(i8,1x))')(ix(i),i=1,nn)
    close(1)
1000 continue
5000 continue
    close(10)
    stop
    end
PLTVEL.FOR
include 'pltv2.for'
include 'timesft.for'
include 'timecrt.for'
program pltvel
c k-net data viewer (integrated velocity data)
parameter(nmax=50000,nch=3)
integer ix(nmax)
real*4 x(nch,nmax),xx(nmax),amaxvel(nch)
character
cid*6,filen*47,dummy*18,cdir*3,cmemo(nch)*80,cstcode*6,
* cdrv*9,ccomp(nch)*2
data aym,x0,y1,DTL/16.0,2.0,2.0,1.0/ fac/1./
data ccomp/'NS','EW','UD'/ tst,ted/0.01,500./
open(10,file='stlist.txt',status='old')
read(10,*)cdrv,nsite
do 5000 isite=1,nsite
read(10,*) cid,tst
filen=cdrv//cid//v1.ns'
write(6,'(a)')cdrv//cid//v1.**'
open(1,file=filen,status='old')

```

```

read(1,'(a18,i4,1x,i2,1x,i2,1x,i2,1x,i2)')
* dummy,irgyear,irgmonth,irgday,irghour,irgmin
read(1,'(a18,f5.1)')dummy,orglat
read(1,'(a18,f5.1)')dummy,orglong
read(1,'(a18,i4)') dummy,idepth
orgdepth=real(idepth)
read(1,'(a18,f4.1)')dummy,amag
read(1,'(a18,a6)')dummy,cstcode
read(1,'(a18,f8.4)')dummy,stlat
read(1,'(a18,f8.4)')dummy,stlong
read(1,'(a18,i4)') dummy,idepth
stheight=real(idepth)
read(1,'(a18,i4,1x,i2,1x,i2,1x,i2,1x,i2,1x,i2)')
* dummy,ircyear,ircmonth,ircday,irchour,ircmin,ircsec
call timecrt(ircyear,ircmonth,ircday,irchour,ircmin,ircsec)
read(1,'(a18,i3)')dummy,iratio
sratio=real(iratio)
dt=1.0/sratio
read(1,'(a18,i3)')dummy,idur
ndata=idur*100
duration=real(idur)
read(1,'(a18,A3)')dummy,cdir
read(1,'(a18,10x,i7)')dummy,iscale
scale=200./real(iscale)
read(1,'(a18,i8)')dummy,imaxvel
amaxvel(1)=real(imaxvel)
read(1,'(a18)')dummy
read(1,'(a80)')cmemo(1)
read(1,'(8(i8,1x))')(ix(i),i=1,ndata)
do 10 i=1,ndata
10  x(1,i)=real(ix(i))*scale
    close(1)
    ntst=int(tst)
    nst=int(tst/dt+0.5)
    ned=int(ted/dt+0.5)
    call timesft(ntst,ircyear,ircmonth,ircday,
    *          irchour,ircmin,ircsec)
    filen=cdrv//cid//v1.ew'
    open(1,file=filen,status="old")
    do 20 kk=1,14
20  read(1,'(a18)')dummy
    read(1,'(a18,i8)')dummy,imaxvel
    amaxvel(2)=real(imaxvel)
    read(1,'(a18)')dummy
    read(1,'(a80)')cmemo(2)
    read(1,'(8(i8,1x))')(ix(i),i=1,ndata)
    do 30 i=1,ndata
30  x(2,i)=real(ix(i))*scale
    close(1)
    filen=cdrv//cid//v1.ud'
    open(1,file=filen,status="old")
    do 40 kk=1,14
40  read(1,'(a18)')dummy
    read(1,'(a18,i8)')dummy,imaxvel
    amaxvel(3)=real(imaxvel)
    read(1,'(a18)')dummy
    read(1,'(a80)')cmemo(3)
    read(1,'(8(i8,1x))')(ix(i),i=1,ndata)
    do 50 i=1,ndata
50  x(3,i)=real(ix(i))*scale
    close(1)
    filen=cdrv//cid//v1.ps'
    write(6,'(A)')filen
    open(25,file=filen,status='unknown')
    CALL PLOTS(25)
    call yoko(25)
    call symbol(1.5,5.0,0.8,cid,90.0,6)
    call number(1.5,11.0,0.25,real(ircyear),90.0,0)
    call number(1.5,12.25,0.25,real(ircmonth),90.0,0)
    call number(1.5,13.00,0.25,real(ircday),90.0,0)
    call number(1.5,13.75,0.25,real(irchour),90.0,0)
    call number(1.5,14.50,0.25,real(ircmin),90.0,0)
    call number(1.5,15.25,0.25,real(ircsec),90.0,0)
    CALL FACTOR(FAC)
    ndatamax=min(1,real(ndata-nst+1),24.*dt/dt)+nst-1
    do 60 kk=1,nch
        amaxvel(kk)=x(kk,nst)
        do 60 i=2,ndatamax-nst+1

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        if(amaxvel(kk).lt.abs(x(kk,nst+i-1))) then
            amaxvel(kk) =abs(x(kk,nst+i-1))
        endif
60 continue
    amaxx=amax1(amaxvel(1),amaxvel(2),amaxvel(3))
    if(amaxx.lt. 500.) xxmax= 500.
    if(amaxx.lt. 200.) xxmax= 200.
    if(amaxx.lt. 100.) xxmax= 100.
    if(amaxx.lt. 50.) xxmax= 50.
    if(amaxx.lt. 20.) xxmax= 20.
    if(amaxx.lt. 10.) xxmax= 10.
    if(amaxx.lt. 5.) xxmax= 5.
    if(amaxx.lt. 2.) xxmax= 2.
    if(amaxx.lt. 1.) xxmax= 1.
    if(amaxx.lt.0.5) xxmax= 0.5
    if(amaxx.lt.0.2) xxmax= 0.2
    if(amaxx.lt.0.1) xxmax= 0.1
    if(amaxx.lt.0.05) xxmax= 0.05
    if(amaxx.lt.0.02) xxmax= 0.02
    if(amaxx.lt.0.01) xxmax= 0.01
    call plot(x0,0,-3)
    CALL NEWPEN(1)
    ay=aym
    if(ndatamax.lt.ndata) call symbol(26.,0.5,0.3,'C',0.0,1)
    do 1000 kk=1,3
        call symbol(2.5,ay+y1+0.3,0.4,ccomp(kk),0.0,2)
        call symbol(4.0,ay+y1+0.3,0.3,'Max=',0.0,4)
        call number(5.5,ay+y1+0.3,0.3,amaxvel(kk),0.0,3)
    do 100 i=1,ndata
100    xx(i)=x(kk,i+nst-1)
        CALL PLOT( 2., AY,-3)
        call pltwv2(xx,ndatamax-nst+1,dt,dtl,xxmax,yl,2)
        CALL PLOT(-2.,-AY,-3)
        ay=ay-3.*yl
1000 continue
    call plote(25)
    close(25)
5000    continue
    close(10)
    stop
    end
ACC2DIS.FOR
    program acc2dis
c    k-net data integrater (convert to displacement data)
    parameter(nmax=65536,nch=3)
    integer ix(nmax),imaxacc(nch),imaxdsp(nch)
    real*4 x(nch,nmax),xx(nmax),amaxacc(nch),amaxdsp(nch)
    character
cid*6,filein(nch)*20,filein*15,dummy(16)*18,cdir(nch)*3,
*
cmemo(nch)*80,cstcode*6,cdrv*9,ccomp(nch)*2,filenm*18
data nst/1/ ntype,nchara,ncausal/1,3,2/
ccomp/'NS','EW','UD'/
data ap,as/0.1,10.0/
ned=nmax
open(10,file='stlist.txt',status='old')
read(10,*) cdrv,nsite
do 5000 isite=1,nsite
read(10,*) cid,tst,ted,fl,fs
filein=cdrv//cid
filein(1)=filein//'.ns'
write(6,(a)filein//'.**')
open(1,file=filein(1),status="old")
read(1,'(a18,i4,1x,i2,1x,i2,1x,i2,1x,i2)')
* dummy(1),irgyear,irgmonth,irgday,irghour,irgmin
read(1,'(a18,f5.1)')dummy(2),orglat
read(1,'(a18,f5.1)')dummy(3),orglong
read(1,'(a18,i4)') dummy(4),idepth
orgdepth=real(idepth)
read(1,'(a18,f4.1)')dummy(5),amag
read(1,'(a18,a6)')dummy(6),cstcode
read(1,'(a18,f8.4)')dummy(7),stlat
read(1,'(a18,f8.4)')dummy(8),stlong
read(1,'(a18,i4)') dummy(9),idepth
stheight=real(idepth)
read(1,'(a18,i4,1x,i2,1x,i2,1x,i2,1x,i2)')
* dummy(10),ircyear,ircmonth,ircday,irchour,ircmin,ircsec
read(1,'(a18,i3)')dummy(11),iratio

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sratio=real(iratio)
dt=1.0/sratio
read(1,'(a18,i3)')dummy(12),idur
ndata=idur*100-8
nst=int(tst/dt+0.5)*0.0
c ned=min(int((ted/dt+0.5),ndata)
ned=ndata
duration=real(idur)
read(1,'(a18,A3)')dummy(13),cdir(1)
read(1,'(a18,10x,i7)')dummy(14),iscale
scale=2000./real(iscale)
read(1,'(a18,i4)')dummy(15),imaxacc(1)
amaxacc(1)=real(imaxacc(1))
read(1,'(a18,i4,1x,i2,1x,i2,1x,i2,1x,i2,1x,i2)')
* dummy(16),icryear,icrmonth,icrday,icrhour,icrmin,ircsec
read(1,'(a80)')cmemo(1)
read(1,'(8(i8,1x))')(ix(i),i=1,ndata)
do 10 i=1,ndata
10 x(1,i)=real(ix(i))*scale
close(1)
filein(2)=filein//'.ew'
open(1,file=filein(2),status="old")
do 20 kk=1,14
20 read(1,'(a18)')dummy(kk)
cdir(2)='EW'
read(1,'(a18,i4)')dummy(15),imaxacc(2)
amaxacc(2)=real(imaxacc(2))
read(1,'(a18)')dummy(16)
read(1,'(a80)')cmemo(2)
read(1,'(8(i8,1x))')(ix(i),i=1,ndata)
do 30 i=1,ndata
30 x(2,i)=real(ix(i))*scale
close(1)
filein(3)=filein//'.ud'
open(1,file=filein(3),status="old")
do 40 kk=1,14
40 read(1,'(a18)')dummy(kk)
cdir(3)='UD'
read(1,'(a18,i4)')dummy(15),imaxacc(3)
amaxacc(3)=real(imaxacc(3))
read(1,'(a18)')dummy(16)
read(1,'(a80)')cmemo(3)
read(1,'(8(i8,1x))')(ix(i),i=1,ndata)
do 50 i=1,ndata
50 x(3,i)=real(ix(i))*scale
close(1)
nn=ned-nst+1
idur=int(real(nn)*dt)
td =real(idur)
nn=td/dt
ntst=int(tst)
do 500 kk=1,3
do 180 i=1,nn
180 xx(i)=x(kk,i+nst-1)
c linear trend fitting
b1=xx(1)
do 192 i=2,nn
192 b1=b1+real(i)*xx(i)
b2=xx(1)
do 194 i=1,nn
194 b2=b2+xx(i)
a=real(nn*(nn+1)*(2*nn+1))/6.*dt
b=real(nn*(nn+1))/2.
c=b*dt
d=real(nn)
f=a*d-b*c
aa=( d*b1-b*b2)/f
bb=(-c*b1+a*b2)/f
do 186 i=1,nn
186 xx(i)=xx(i)-aa*real(i)*dt-bb
c band pass filtering
call bandp1(xx,nn,dt,fl,fs,ap,as,ntype,nchara,ncausal)
c integration
888 xx(1)=xx(1)*dt
do 200 i=2,nn
200 xx(i)=xx(i-1)+xx(i)*dt
c linear trend fitting
b1=xx(1)

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```

do 292 i=2,nn
292   b1=b1+real(i)*xx(i)
      b2=xx(1)
do 294 i=1,nn
294   b2=b2+xx(i)
      a=real(nn*(nn+1)*(2*nn+1))/6.*dt
      b=real(nn*(nn+1))/2.
      c=b*dt
      d=real(nn)
      f=a*d-b*c
      aa=( d*b1-b*b2)/f
      bb=(-c*b1+a*b2)/f
do 286 i=1,nn
286   xx(i)=xx(i)-aa*real(i)*dt-bb
c band pass filtering
call bandp1(xx,nn,dt,fl,fh,fs,ap,as,ntype,nchara,ncausal)
c integration
xx(1)=xx(1)*dt
do 300 i=2,nn
300   xx(i)=xx(i-1)+xx(i)*dt
      amaxdsp(kk)=abs(xx(1))
do 400 i=2,nn
400   if(amaxdsp(kk).lt.abs(xx(i))) amaxdsp(kk)=abs(xx(i))
      imaxdsp(kk)=int(amaxdsp(kk)*10000.+0.5)
do 410 i=1,nn
410   x(kk,i)=xx(i)
500 continue
do 600 kk=1,nch
      filenm=filen(kk)(1:15)//'DS.'
      filen(kk)=filenm//ccomp(kk)
      write(6,'(1x,a23)') filen(kk)
600 continue
c 0.25 micron per digit
scale=2.5e-5
iscale=int(25./scale+0.5)
c
do 1000 kk=1,nch
do 1010 i=1,nn
1010  ix(i)=int(x(kk,i)/scale)
      open(1,file=filen(kk),status='unknown')
      write(1,'(a18,i4,1x,i2,1x,i2,1x,i2,1x,i2)')
      *
      *
      dummy(1),
irgyear,irgmonth,irgday,irghour,irgmin
write(1,'(a18,f5.1)')dummy(2),orglat
write(1,'(a18,f5.1)')dummy(3),orglong
write(1,'(a18,i4)') dummy(4),idepth
write(1,'(a18,f4.1)')dummy(5),amag
write(1,'(a18,a6)') dummy(6),cstcode
write(1,'(a18,f8.4)')dummy(7),stlat
write(1,'(a18,f8.4)')dummy(8),stlong
write(1,'(a18,i4)') dummy(9),idepth
write(1,'(a18,i4,1x,i2,1x,i2,1x,i2,1x,i2,1x,i2)')
*
*
      dummy(10),
ircyear,ircmonth,ircday,irchour,ircmin,ircsec
write(1,'(a18,i3)') dummy(11),iratio
write(1,'(a18,i3)') dummy(12),idur
write(1,'(a18,A3)') dummy(13),cdir(kk)
write(1,'(a18,a10,i7)')dummy(14),'25( cm )',iscale
write(1,'(a18,i8)') 'Max.Disp.(micron) ',imaxdsp(kk)
write(1,'(a18,i4,1x,i2,1x,i2,1x,i2,1x,i2,1x,i2)')
*
*
      dummy(16),
icryear,icrmonth,icrday,icrhour,icrmin,icrsec
write(1,'(a80)') cmemo(kk)
write(1,'(8(i8,1x))')(ix(i),i=1,nn)
close(1)
1000 continue
5000 continue
close(10)
stop
end
PLTDIS.FOR
include 'pltwv2.for'
include 'timesft.for'
include 'timecrt.for'
program pltdis
c
k-net data viewer (integrated displacement data)
parameter(nmax=50000,nch=3)
integer ix(nmax)
real*4 x(nch,nmax),xx(nmax),amaxvel(nch)
character
cid*6,filen*47,dummy*18,cdir*3,cmemo(nch)*80,cstcode*6,
*
*
      cdrv*9,ccomp(nch)*2
      data aym,x0,yl,DTL/16.0,2.0,2.0,1.0/ fac/1./
      data ccomp/'NS','EW','UD'/ tst,tcd/0.01,500./
      open(10,file='stlist.txt',status='old')
      read(10,*)cdrv,nsite
      do 5000 isite=1,nsite
      read(10,*) cid,tst
      filen=cdrv//cid//ds.ns'
      write(6,'(a)')cdrv//cid//ds.*'
      open(1,file=filen,status="old")
      read(1,'(a18,i4,1x,i2,1x,i2,1x,i2,1x,i2)')
      *
      dummy,irgyear,irgmonth,irgday,irghour,irgmin
      read(1,'(a18,f5.1)')dummy,orglat
      read(1,'(a18,f5.1)')dummy,orglong
      read(1,'(a18,i4)') dummy,idepth
      orgdepth=real(idepth)
      read(1,'(a18,f4.1)')dummy,amag
      read(1,'(a18,a6)')dummy,cstcode
      read(1,'(a18,f8.4)')dummy,stlat
      read(1,'(a18,f8.4)')dummy,stlong
      read(1,'(a18,i4)') dummy,idepth
      stheight=real(idepth)
      read(1,'(a18,i4,1x,i2,1x,i2,1x,i2,1x,i2,1x,i2)')
      *
      dummy,ircyear,ircmonth,ircday,irchour,ircmin,ircsec
      call timecrt(ircyear,ircmonth,ircday,irchour,ircmin,ircsec)
      read(1,'(a18,i3)')dummy,iratio
      sratio=real(iratio)
      dt=1.0/sratio
      read(1,'(a18,i3)')dummy,idur
      ndata=idur*100
      duration=real(idur)
      read(1,'(a18,A3)')dummy,cdir
      read(1,'(a18,10x,i7)')dummy,iscale
      scale=25./real(iscale+0.5)
      read(1,'(a18,i8)')dummy,imaxvel
      amaxvel(1)=real(imaxvel)
      read(1,'(a18)')dummy
      read(1,'(a80)')cmemo(1)
      read(1,'(8(i8,1x))')(ix(i),i=1,ndata)
do 10 i=1,ndata
10  x(1,i)=real(ix(i))*scale
      close(1)
      filen=cdrv//cid//ds.ew'
      open(1,file=filen,status="old")
      do 20 kk=1,14
20  read(1,'(a18)')dummy
      read(1,'(a18,i8)')dummy,imaxvel
      amaxvel(2)=real(imaxvel)
      read(1,'(a18)')dummy
      read(1,'(a80)')cmemo(2)
      read(1,'(8(i8,1x))')(ix(i),i=1,ndata)
do 30 i=1,ndata
30  x(2,i)=real(ix(i))*scale
      close(1)
      filen=cdrv//cid//ds.ud'
      open(1,file=filen,status="old")
      do 40 kk=1,14
40  read(1,'(a18)')dummy
      read(1,'(a18,i8)')dummy,imaxvel
      amaxvel(3)=real(imaxvel)
      read(1,'(a18)')dummy
      read(1,'(a80)')cmemo(3)
      read(1,'(8(i8,1x))')(ix(i),i=1,ndata)
do 50 i=1,ndata
50  x(3,i)=real(ix(i))*scale
      close(1)
      ntst=int(tst)
      nst=int(tst/dt+0.5)
      ned=int(ted/dt+0.5)
      call timesft(ntst,ircyear,ircmonth,ircday,
      *
      *
      irchour,ircmin,ircsec)
      filen=cdrv//cid//ds.ps'

```

```

write(6,*) filen
open(25,file=filen,status='unknown')
CALL PLOTS(25)
call yoko(25)
call symbol(1.5,5.0,0.8,cid,90.0,6)
call number(1.5,11.0,0.25,real(ircyear),90.0,0)
call number(1.5,12.25,0.25,real(ircmonth),90.0,0)
call number(1.5,13.00,0.25,real(ircday),90.0,0)
call number(1.5,13.75,0.25,real(irchour),90.0,0)
call number(1.5,14.50,0.25,real(ircmin),90.0,0)
call number(1.5,15.25,0.25,real(ircsec),90.0,0)
CALL FACTOR(FAC)
ndatamax=min(1,real(ndata-nst+1),24.*dtl/dt)+nst-1
do 60 kk=1,nch
  amaxvel(kk)=x(kk,nst)
  do 60 i=2,ndatamax-nst+1
    if(amaxvel(kk).lt.abs(x(kk,nst+i-1))) then
      amaxvel(kk) =abs(x(kk,nst+i-1))
    endif
  60 continue
  amaxx=amax1(amaxvel(1),amaxvel(2),amaxvel(3))
c write(6,*)amaxvel(1),amaxvel(2),amaxvel(3)
  if(amaxx.lt.500000.) xxmax=500000.
  if(amaxx.lt.200000.) xxmax=200000.
  if(amaxx.lt.100000.) xxmax=100000.
  if(amaxx.lt.50000.) xxmax=50000.
  if(amaxx.lt.20000.) xxmax=20000.
  if(amaxx.lt.10000.) xxmax=10000.
  if(amaxx.lt.5000.) xxmax=5000.
  if(amaxx.lt.2000.) xxmax=2000.
  if(amaxx.lt.1000.) xxmax=1000.
  if(amaxx.lt. 500.) xxmax= 500.
  if(amaxx.lt. 200.) xxmax= 200.
  if(amaxx.lt. 100.) xxmax= 100.
  if(amaxx.lt. 50.) xxmax= 50.
  if(amaxx.lt. 20.) xxmax= 20.

  if(amaxx.lt. 10.) xxmax= 10.
  if(amaxx.lt. 5.) xxmax= 5.
  if(amaxx.lt. 2.) xxmax= 2.
  if(amaxx.lt. 1.) xxmax= 1.
  if(amaxx.lt.0.5) xxmax= 0.5
  if(amaxx.lt.0.2) xxmax= 0.2
  if(amaxx.lt.0.1) xxmax= 0.1
  if(amaxx.lt.0.05) xxmax= 0.05
  if(amaxx.lt.0.02) xxmax= 0.02
  if(amaxx.lt.0.01) xxmax= 0.01
  if(amaxx.lt.0.005) xxmax= 0.005
  if(amaxx.lt.0.002) xxmax= 0.002
  if(amaxx.lt.0.001) xxmax= 0.001
  xxmax=xxmax*10000.
  call plot(x(0,.-3)
  CALL NEWPEN(1)
  ay=aym
  if(ndatamax.lt.ndata) call symbol(26.,0.5,0.3,'C',0.0,1)
  do 1000 kk=1,3
    call symbol(2.5,ay+y1+0.3,0.4,ccomp(kk),0.0,2)
    call symbol(4.0,ay+y1+0.3,0.3,'Max=',0.0,4)
    call number(5.5,ay+y1+0.3,0.3,amaxvel(kk)*10000.,0.0,4)
  100 xx(i)=x(kk,i+nst-1)*10000.
    CALL PLOT( 2., AY,-3)
    call pltwv2(xx,ndatamax-nst+1,dt,dtl,xxmax,y1,3)
    CALL PLOT(-2.,-AY,-3)
    ay=ay-3.*y1
  1000 continue
    call plote(25)
    close(25)
  5000 continue
    close(10)
    stop
    end

```

## Reference for further reading

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